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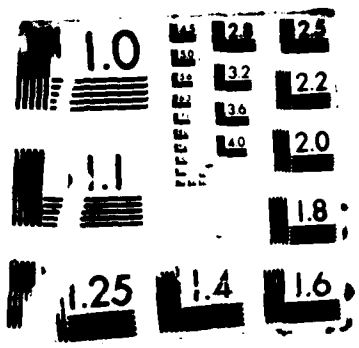
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The hydrogen spectra from the CH(4)/N(2)O and nitramine composite flames allow use of each of the concurrent capabilities of CARS: (1) rovibrational state resolution better than a wavenumber, (2) time resolution better than 10 nrs, and (3) spatial resolution better than 100 micrometers. The rovibrational state resolution results in direct observation of previously unobserved higher J transitions in the Q branch v" = 0 and v" = 1 and S branch v" = 0 transitions. These transitions are in accord with transitions calculated from ab initio resluts and with transitions indirectly obtained from electronic emission ~~(cont)~~

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spectra. The ab initio results give a compact set of constants that gives results within the experimental accuracy of the CARS experimental results.

The experimental transition energies were used to construct Boltzmann plots that allow discrimination among the various spectra used to determine temperature.

CARS provides the spectral resolution and intensity to allow direct observation of previously unobserved hydrogen higher J and v rotational transitions. These transitions were used to obtain temperature profiles and in the CH(4)/N(2) and nitramine composite flames that require both the high spatial and temporal resolution of CARS. These temperature profiles provided additional confirmation of previous kinetic mechanisms used to explain observations in the nitramine composite flame. The concurrent capabilities of CARS used here demonstrate its high potential for use as the probe beam in kinetic flash photolysis and direct in situ measurements of combustion flames.

Keywords

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INTRODUCTION

Coherent Antistokes Raman Spectroscopy (CARS) provides one of the most powerful new techniques for monitoring the progress of fast chemical reactions in uses such as the probe beam in flash photolysis. Electronic spectroscopy was first used as the probe beam in flash photolysis as originally developed by Norrish and Porter (ref 1). Similarly to electronic spectroscopy, CARS has also been extensively used for analysis of flames (refs 2 through 4) prior to use as the probe beam in flash photolysis, as for example, in refs 5 and 6. CARS allows direct monitoring of the rovibrational state distributions of reactant, transient, and product species. In addition, the nonlinear scattering processes from which CARS originates enables excellent time resolution, presently shorter than 10 ns (ref 4). CARS can also be easily optically configured to give a sampling extent less than 100 μm. These concurrent capabilities, rovibrational state resolution better than a wavenumber, time resolution better than 10 ns, and spatial resolution better than 100 μm, allow direct observation of the fast combustion reactions occuring in the reaction zone of flames where preceptible chemical change was spatially resolved on the scale of the sampling extent. This was first demonstrated experimentally in studies of the reaction zone of lean CH_4/N_2O flames (refs 2 and 3) where decay of N_2O and concomitant N_2 formation were simultaneously observed on a 100 μm spatial scale (which corresponds to millisecond time resolution for the burning velocity of the flame studied).

This capability was subsequently extended to the reaction zone of rich CH_4/N_2O flames and the flames of solids (the nitramine RDX, hexahydro-1,3,5,-trinitro-3-triazine in an organic ester matrix) (refs 7 through 9).

Spectra were obtained in the regions 4200-3900, 2400-2050, and 1900-1200 cm^{-1}. The reaction zone of the rich CH_4/N_2O flame was studied primarily to provide a stationary flame analog to the transient propellant flame. In the CH_4/N_2O flames, the decay of the initial products was observed through the Q branch of the ν_1 and ν_3 modes of N_2O and Q, O, and S branches of the ν_2 and $2\nu_2$ modes of CH_4. The formation of the products N_2, H_2[Q(v" = 0 and v" = 1) and S(5) - S(9)], CO, and CO_2 (ν_1) were also observed. In the nitramine propellant flame near the surface of the propellant, reactant RDX (1599 cm^{-1} tentatively assigned as asymmetric NO_2 stretch) and transients HCN (ν_1) and NO are observed at moderate concentration (>1%). The final product N_2 is observed at low concentration (~ 1%); H_2 (Q and S branches) and CO are observed at higher concentration ($\geq$10%). RDX and HCN decay within 2 mm of the propellant surface, while NO remains constant until 4 mm, where it decays with a concomitant rise in N_2 concentration and temperature. H_2 and CO also increase in temperature and concentration. The nitramine reaction zone was seen to consist of two regions characterized by the reactions of RDX and HCN near the surface, consistent with the high-temperature mechanism of RDX decomposition, and the conversion of NO to N_2 to generate the luminous flame further above the surface.

The observed spectra demonstrate many of the features previously discussed which make CARS useful for the probe beam in flash photolysis. This is clearly

seen in the rotationally well-resolved hydrogen spectra from the reaction zone of rich CH_4/N_2O and nitramine composite flames. These CARS hydrogen spectra are treated here in greater detail than that given in previous reports (refs 7 through 9). Higher rotational transitions in the Q and S branches of the ground and first excited vibrational state, many of which were directly observed for the first time (ref 10), are compared to transitions calculated from constants derived from ab initio calculations (ref 11) and transitions calculated from those observed in the rovibrational structure of electronic emission spectra (ref 12). These hydrogen transitions are used to obtain temperature profiles through the reaction zone of rich CH_4/N_2O flames and, for the first time, of nitramine composite flames. The interpretatins involved in the reduction of the hydrogen data remove some of the ambiguities of data interpretation previously reported (ref 9). These observed temperature profiles aid in the interpretation of the observed species in terms of elementary reactions.

Flash photolysis using CARS as the probe beam will be used in future work for determining the elementary reactions required to elucidate the kinetic mechanisms directly observed using CARS. Consequently, CARS will be useful not only for observing combustion flame phenomena directly in situ but also using the same or similar apparatus for elucidating these phenomena through elementary reactions determined using flash photolysis and related techniques.

EXPERIMENTAL

CARS spectra were generated using the folded BOXCARS apparatus (ref 9). Briefly, the output of the Quantra-Ray DCR-2A Nd/YAG laser at 1.06 μm (700 mJ) is doubled to generate the pump beam at 5320A (250 mJ) with a bandwidth of near 1 cm^{-1}. The pump beam is split to generate ω_{1s} and ω_{1p}. ω_{1s} is used to pump a dye laser to generate the Stokes beam ω_2. The dye laser consists of a flowing dye cell in a planar Fabry-Perot oscillator cavity pumped slightly off-axis by 20% of ω_{1s}, with the output amplified by an additional flowing dye cell pumped by the remainder of ω_{1s}. The dye laser is operated broadband with the laser dyes Exciton DCM and LDS in ethanol to generate a nonresonant spectra centered near 4210 cm^{-1} with a bandwidth of 300 cm^{-1} for H_2 v = 0 Q branch spectra. ω_1 is split such that ω_1, ω_1', and ω_2 are placed on a 12.5 mm circle. ω_3 was focused into a monochromator equipped with a PAR SIT detector interfaced to a PAR OMA2 system. The full-width-at-half-maximum (FWHM) of calibration lines near the center of the SIT detector is 3.0 cm^{-1}, giving approximately 1 cm^{-1} per channel over the spectral range investigated.

Stationary flame measurements were made on a premixed CH_4/N_2O flame maintained on a circular burner with a 2.0-cm-diameter head whose surface was constructed of a matrix of steel syringe needles of 1 mm outer diameter. A mixture of Matheson technical grade methane and chemically pure nitrous oxide were flowed through the burner at 13 cm/s to maintain a 3.2 equivalence ratio (ϕ)

flame, where ϕ is defined as the fuel/oxidant ratio divided by the stoichiometric fuel/oxidant ratio. To obtain CARS spectra in the reaction zone, the center of the burner surface was displaced vertically at intervals of 0.5 mm (0.25 mm in the vicinity of the reaction zone). The nitramine composite grains burned were 14 x 14 mm^2 cylinders of mass 3.2 g. The composite consisted of 76% RDX (hexahydro-1,3,5,-trinitro-s-triazine), 20% carbohydrates, and 4% nitrocellulose. The propellant grains were burned in air with spectra taken along the centerline above the burning propellant surface during the approximately 1-min burn time. The calculated approximate gas velocity from the burning cylinder is 50 cm/s. Thermochemical calculations were performed for both the CH_4/N_2O flame and nitramine composite flames. The calculated flame temperature for CH_4/N_2O was 1745 K with 23% CO, 1% CO_2, 42% H_2, 5% H_2O, and 29% N_2; and for the nitramine composite 2076 K with 27% H_2, 22% N_2, 10% H_2O, 3% CO_2, and 38% CO.

RESULTS

The observed CARS spectrum in the isolated line approximation (ref 13) is proportional to the square of the modulus of the third-order susceptibility ($\chi^{(3)}$) which is the sum of a resonant term χ_r related to nuclear displacement and χ_{nr} related to electronic displacement. χ_r is composed of a real component χ', which displays dispersive behavior and an imaginary component χ which displays resonant behavior such that:

$$\left|\chi^{(3)}\right|^2 = \left|\chi' + \chi'' + \chi_{nr}\right|^2 = \sum_j \left|\frac{kj\,\Gamma j}{2\,\Delta\omega j - i\Gamma_j} + \chi_{nr}\right|^2 \qquad (1)$$

$$kj = (N/M\omega o)\left|M_j\right|^2 (\Delta j)\,\Gamma_j^{-1} \qquad (2)$$

where M_j, Δj, and Γj are the polarizability matrix element, normalized population difference and line width, respectively, $\Delta\omega_j = \omega_1 - \omega_2 - \omega_j$, m is the reduced mass, ω_o is the resonant Raman frequency, and N is the molecular gas density. $M = \alpha^2(\nu + 1)$ and $7/45\ bj\ \nu^2\ (\nu + 1)$ for Q and O, S branches, respectively, where α, ν, and b_J are the derivatives of the mean isotropic and anisotropic molecular polarizability, and b_J are the Placzck-Teller coefficients, ν is the vibrational quantum number, and $(\nu + 1)$ is contributed by the vibrational matrix element.

The interpretation of hydrogen CARS spectra under the conditions in which it is observed here is particularly straightforward. The Raman cross section of hydrogen (from which Mj can be calculated) is sufficiently large (it is twice that of nitrogen) that at the hydrogen concentrations observed in most of these experiments, the effect of the χ_{nr} may be neglected. In this case, the peak CARS intensity, Ij^{max}, is given by the imaginary component of $\chi^{(3)}$

$$(Ij^{max})^{1/2} = \frac{(N \quad)}{m\omega_o} \left| Mj \right|^2 (\Delta j)\, \Gamma j^{-1}$$

Neglecting the aniostropic component of Mj and defining $\Delta j = (2J + 1)\, gj\, \bar{\Delta} j$ where gj is the nuclear spin statistical degeneracy (g = 3 for odd levels and g = 1 for even levels) and $\bar{\Delta} j$ is the population difference per rotational energy level

$$(\bar{I}j^{max})^{1/2} = \frac{(Ij^{max})^{1/2}}{(2J + 1)\, g_j\, (\nu + 1)} = \left(\frac{N}{m\,\omega_o}\right) \alpha \left| \right|^2 \Gamma_j^{-1}\, \bar{\Delta} j$$

Γj^{-1} as observed in these experiments is determined by the convolution of the instrument response function (FWFM = 3.0 cm^{-1}) over the doppler broadened transitions (where the linewidth is of the order of tenths of a wavenumber). Under these conditions the observed linewidth is effectively constant and equal to that of the instrument. Therefore, $(Ij^{max})^{1/2}$ is related by a constant to $\bar{\Delta} j$. If the excited state population is neglected, $\bar{\Delta} j$ can be assumed to be given by the Boltzmann relation $\bar{\Delta} j = \exp\left[\frac{-k\,\Delta Ej}{T}\right]$

The relation

$$\ln\,(\bar{I}j^{max}) = C\,\frac{k}{T}\,\Delta Ej$$

(where k is the Boltzmann constant) can be used to determine temperature from the peak heights. The separation of the rotational lines and the lack of modulation by χ_{nr} eliminate the complexities of data reduction recently considered (ref 13).

CH_4/N_2O Flames

A summary of the species observed for the ϕ = 3.2 rich CH_4/N_2O flame is given in table 1. Hydrogen Q-band structure as it increases through the reaction zone is shown in figure 1. Hydrogen is seen at a concentration less than 1% at 4 mm. The line positions as shown for v" = 0 (fig. 1) and v" = 0, J $\leq$ 11 and v" = 1, J $\leq$ 9 (fig. 2), and the S bands J = 5 - 9 (table 1) for v" = 0 are compared to experimentally determined transitions (ref 12) and to transitions calculated from spectroscopic constants derived from ab initio calculations (ref 11). The constants given in table 2 are the same as those given in ref 11 with the exception that typographical errors for H_1 and H_2 have been corrected. These constants (table 2) result in transition energies (tables 3 and 4) that agree with rovibrational transition energies obtained from analysis of the Lyman $B^1\,{}^+_u$ --> $X^1\,{}^+_g$ and Werner $C^1\,\Pi u$ --> $X^1\,{}^+_g$ electronic transitions. The directly observed rovibrational transitions agree within their 1 cm^{-1} experimental accuracy with both the previously observed experimental and calculated transitions. The advantage of the constants given in table 2 is that they agree with the CARS data within the experimental error and are in a compact form convenient for use in spectral analysis.

Higher than 8 mm above the burner surface there is an apparent bimodal distribution in which approximately half of the observed spectra have Boltzmann distributions consistent with the random experimental error, while the other spectra show deviations from a Boltzmann distribution which is much larger (greater then 2σ). An example of this behavior is illustrated by Boltzmann plots shown in figures 3 and 4. At the same position in the flame (18 mm above the burner surface) spectra taken under similar conditions give widely disparate temperatures. However, of the two distributions only the lower temperature result appears Boltzmann. In the non-Boltzmann distributions, the odd levels are preferentially populated over the even levels. While these apparently non-Boltzmann distributions may reflect the actual hydrogen rotational distribution, it seems more likely due to flame instability which increases with distance above the burner. The outer blue diffusion flame is noticeably floppy and of higher temperature than the adiabatic CH_4/N_2O flame. Spectra which result from a mixture of adiabatic and diffusion flames due to flame flicker would yield non-Boltzmann results perhaps similar to that in figure 4. Use of only Boltzmann results allows construction of the temperature profile shown in figure 5.

Nitramine Composite Flames

A summary of species observed in the nitramine composite flame is given in table 5. Average spectra (100 scans, 10 s) were taken as a function of distance from the propellant surface to 6 mm above the surface at intervals of 1 mm. Each spectrum was taken nominally 10 s after ignition. In addition, time sequences of ten-scans (1 s) spectra were taken approximately every 6 s from ignition to extinction. The intensity of the H_2 signal permitted acquisition of single-shot spectra. Time-resolved spectra of the H_2 Q branch at the propellant surface are given in figures 6 and 7. Single-shot spectra (fig. 6) are obtained at good signal-to-noise and are in substantial agreement with 10-shot averaged spectra (fig. 7). The reduction of the signal-to-noise ratio from bottom to top is a reflection of the increasing temperature as a function of distance above the propellant surface. Average spectra taken 12 and 18 seconds after ignition were used to construct the temperature profile shown in figure 8. The average surface temperature is 1100 K $\pm$ 200 K (averaging data taken within 2 mm of the surface). The dispersion reflects not only the noise in the individual spectra but also the variation of the distance of the surface with respect to the CARS sampling volume, since spectra are taken as close to the surface as possible. Both sets of spectra show the same trend, i.e., an increase in temperature from the surface to 4 mm where the temperature levels off at 2000 K which is close to the calculated adiabatic flame temperature. The spread of the data is such that the functional form of the variation of temperature with distance is difficult to determine. One of the simplest interpretations is to consider all points within 2 mm of the surface as equivalent with the lowest temperature measured (900 $\pm$ 100 K) being considered as the upperbound of the gas-surface interface temperature. Alternatively, both sets of data may be taken as indicating a decrease in temperature from the surface to about 2 mm. More precise data are needed to determine the precise functional form of the temperature gradient near the surface.

DISCUSSION

The H_2 spectra from the CH_4/N_2O and nitramine composite flame allow use of each of the concurrent capabilities of CARS: (1) rovibrational state resolution better than a wavenumber, (2) time resolution better than 10 ns, and (3) spatial resolution better than 100 μm. The revibrational state resolution results in direct observation of previously unobserved higher J transitions in the Q-branch v" = 0 and v" = 1 and S-branch v" = 0 transitions. These transitions are in accord with transitions calculated from ab initio results and with transitions indirectly obtained from electronic emission spectra (tables 3 and 4). The ab initio results (ref 10) give a compact set of constants given in table 2 that give results within the experimental accuracy of the CARS experimental results.

The experimental transition energies have been used to construct Boltzmann plots (examples of which figs. 3 and 4) that allow discrimination among the various spectra used to determine temperature. A bimodal temperature distribution is obtained above a certain height from the burner surface. However, only the lower temperatures, which are in accord with the calculated flame temperature, give a Blotzmann plot consistant with the experimental error. This is perhaps indicative of the influence of flame flicker; although, the occurrence of a certain fraction of nonequilibrated ortho/para hydrogen cannot be definitively discounted. Whatever the source of the non-Boltzmann hydrogen distributions, it is clear that these cannot be used to obtain temperature, which is only defined in the context of a Blotzmann distribution.

Use of this Boltzmann discrimination allows construction of temperature profiles in both CH_4/N_2O (fig. 5) and nitramine composite flames. Reaction occurs over a region extending from 2 mm to 14 mm with the steepest concentration gradients occurring between 10 mm and 14 mm with a post flame region above 14 mm. As given in table 1, only decay of the initial products was observed through the Q branch of the ν_1 NN and ν_3 NO stretching modes of N_2O and the Q, O, and S branches of the ν_2 and $2\nu_2$ modes of CH_4. No intermediate species were detected. Initial decomposition of the reactants was observed to occur near 500 K. Secondary reactions for N_2O given by Balaklnine et. al. (ref 14)

$$H + N_2O \text{ ---> } N_2 + OH \qquad k = 6 \times 10^{13}\exp(-13100/RT)$$

and CH_4 given by Tabayaski and Bauer (ref 15)

$$CH_4 + H \text{ ---> } CH_3 + H_2 \qquad k = 7.23 \times 10^{14}\exp(-15600/RT)$$

explain the reactivity at 500 K.

The conversion of methane to final products is thought to proceed by the following global mechanism (ref 16)

$$CH_4 \text{ ---> } CH_3 \text{ ---> } CH_2O,\ CHO \text{ ---> } H_2,\ CO.$$

The conversion of fuel-bound nitrogen is an area of active current research, thought to occur by the global mechanism (refs 17 and 18)

$$\text{Fuel N} \longrightarrow \text{HCN} \longrightarrow NH_i \longrightarrow NO \quad i - 0, 1, 2$$
$$NH_i \xrightarrow{NO} N_2$$

The CH_4/N_2O flame results provide a comparison for results obtained in nitramine composite flames.

The Boltzmann criteria discussed above was also used to obtain a temperature profile in the nitramine composite flame (fig. 8) as a function of time after ignition. Fewer non-Boltzmann distributions were encountered in the nitramine composite flame than in the CH_4/N_2O flame. The spectra taken at 12 and 18 seconds after ignition give similar results within the experimental error. If the data within 2 mm of the surface are averaged, a temperature of 1100 $\pm$ 200 K is obtained. The lower-bound of the averaged data is 900 $\pm$ 100 K which maybe identified with the upper-bound of the gas/surface interface temperature. However, the temperature profile may be interpreted alternately as showing a temperature decrease from the surface to 2 mm. Above 2 mm, both sets of data show a temperature increase to 4 mm where the temperature levels off near the calculated abiabatic flame temperature of 2076 K. This temperature profile is consistent with the previous observation of RDX and HCN at the surface and decaying within 2 mm of the surface, with NO remaining constant to 4 mm. Near 4 mm, NO decayed rapidly with a concomitant increase in N_2 concentration. H_2, CO, and CO_2 increased in concentration throughout this region. These observations are qualitatively consistent with the kinetic mechanism previously invoked to explain spectra obtained in the nitramine composite flame (ref 9).

The observation of HCN and RDX and lack of observation of N_2O (<0.1%) is consistent with the high temperature (T>600K) nitramine decomposition mechanism.

$$(CH_2NNO_2)_3 \longrightarrow 3HCN + 6NO_2 + 3H_2$$

(NO_2 is not observable by CARS under the condition of the experiment.) Therefore, RDX decomposition, as opposed to processes occuring in the CH_4/N_2O flame, allows a direct buildup of the intermediates HCN and NO near the propellant surface so that they are directly observable in CARS. At these pressures, it is the decomposition processes of the species at the surface, RDX and HCN, that supply the heat that determines the burning rate of the nitramine propellant. NO conversion to N_2, which provides the heat for the luminous flame occurs too far upstream to affect the surface.

The reaction zone of nitraine propellant as thus seen to consist of two characteristic areas: (1) an inner flame area near the solid gas interface, which is at a temperature above 900 $\pm$ 100 K and is characterized by the gas-phase reactions of RDX and HCN that provide the heat that determine the burning rate and (2) an outer flame area where NO is converted to N_2 to generate the luminous flame.

CARS provides the spectral resolution and intensity to allow direct observation of previously unobserved hydrogen higher J and v rotational transitions.

These transitions have been used to obtain temperature profiles in the CH_4/N_2O and nitramine composite flames which require both the high spatial and temporal resolution of CARS. These temperature profiles have provided additional confirmation of previous kinetic mechanisms used to explain observations in the nitramine composite flame. These concurrent capabilities of CARS used here demonstrate the high potential of CARS for use as the probe beam in kinetic flash photolysis and direct in situ measurements of combustion flames.

Table 1. Summary of species identified in a $\phi = 3.2$ CH_4/N_2O flame

Observed I_{max} (cm^{-1})	Species	Comment
4155-4075	H_2 Q branch	Intensity increases up the reaction zone
3240-3100	CH_4 ($2\ \nu_2$)	Rapid decrease in reaction zone
2325	N_2	Gradual increase in reaction zone
2222	N_2O (ν_3)	Rapid decrease in reaction zone
2136	CO	Increase in reaction zone
2129	H_2 S(9)	Rapid increase of intensity in reaction zone; large compared to CO
1813	H_2 S(7)	Intensity increases up the reaction zone
1636	H_2 S(6)	Weak signal seen in post flame
1531	CH_4 (ν_2)	Gradual decrease in reaction zone
1447	H_2 S(5)	Intensity increases up the reaction zone
1294	N_2O (ν_1)	Intense signal that decreased rapidly in reaction zone

Table 2. Hydrogen spectroscopic constants derived from ab initio calculations

W_e	4400.39	Be	60.7922
X_eW_e	120.814	α_e	-3.0320
Y_eW_e	0.7241	γ_e	0.0350
Z_eW_e	0.0	H_0	3.23×10^{-5}
D_e	0.0448	H_1	-8.00×10^{-9}
β_e	-0.0016	H_2	-5.00×10^{-8}
δ_e	4.5×10^{-5}		

$$E(v, J) = W_e (v + 0.5) - X_eW_e(v + 0.5)^2 + Y_eW_e(v + 0.5)^3$$

$$- Z_eW_e(v + 0.5)^4 + B_v J(J + 1) - D_vJ^2 (J + 1)^2$$

$$+ H_v J^3(J + 1)^3$$

$$B_v = B_e + \alpha_e (v + 0.5) + \gamma_e (v + 0.5)^2$$

$$D_v = D_e + \beta_e (v + 0.5) + \delta_e (v + 0.5)^2$$

$$H_v = H_o + H_1 (v + 0.5) + H_2 (v + 0.5)^2$$

Table 3. Raman frequencies (cm^{-1}) of H_2 Q-transitions

J	Calculations	v" = 0, Q(J)*	Observed	Calculations	v" = 1, Q(J)*	Observed
0	4161.11	4161.14		3926.00	3925.79	
1	4155.20	4155.25	4154	3920.23	3920.06	3919
2	4143.4	4143.47	4143	3908.70	3908.47	3907
3	4125.79	4125.87	4126	3891.50	3891.29	3892
4	4102.48	4102.58	4102	3868.73	3868.50	3868
5	4073.61	4073.74	4073	3840.52	3840.20	3840
6	4039.37	4039.52	4038	3807.03	3806.92	3806
7	3999.96	3999.87	4000	3768.47	3768.24	3768
8	3955.64	3956.04	3954	3725.06	3724.35	3725
9	3906.69	3906.31	3905	3677.08	3676.52	3676
10	3853.42	3852.98	3853	3568.52		
11	3796.19	3795.09	3794	3508.62		

* T. Haw, W. Y. Cheung, G. C. Chiu, and L. E. Harris, "A study of Flame Species Using CARS," 40th Symposium on Molecular Spectroscopy, Ohio Stante University Abstract WH10, ; 106 (1985).

Table 4. Raman frequencies (cm^{-1}) of H_2 S-transitions

J"	Calculated for v" = 0, S(J")*		Observed
0	354.13	354.35	
1	586.74	587.04	
2	814.22	814.43	
3	1034.67	1034.67	
4	1246.37	1245.98	
5	1447.80	1447.36	1447
6	1637.69	1636.97	1636
7	1815.08	1814.40	1813
8	1979.37	1979.08	
9	2130.32	2130.06	2129

* T. Haw, W. Y. Cheung, G. C. Chiu, and L. E. Harris, "A Study of Flame Species Using CARS," 40th Symposium on Molecular Spectroscopy, Ohio State University Abstract WH10, p 106 (1985).

Table 5. Summary of species identified in nitramine propellant flame

Observed I_{max} (cm^{-1})	Species	Command
4155-4075	H_2 Q branch	Temperature calculations indicate a temperature of 900 K at the surface of the propellant
2325	N_2	Slow increase until near the end of reaction zone where a large increase occurs
2136	CO	Intensity increases up reaction zone
2129	H_2 S(9)	Similar intensity to CO observed
2086	HCN (ν_1)*	Strong signal initially which diminishes rapidly
1872	NO	Low concentration modulation which remains constant throughout reaction zone; decreases rapidly at end of reaction zone
1814	H_2 S(7)	Signal intensity increases up the reaction zone
1599	RDX (NO_2 asymetric stretch)*	Modernate concentration early in the reaction zone
1447	H_2 S(5)	Intensity increases up the reaction zone
1387	CO_2 (ν_1)	Moderate intensity early in reaction

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* Tentative

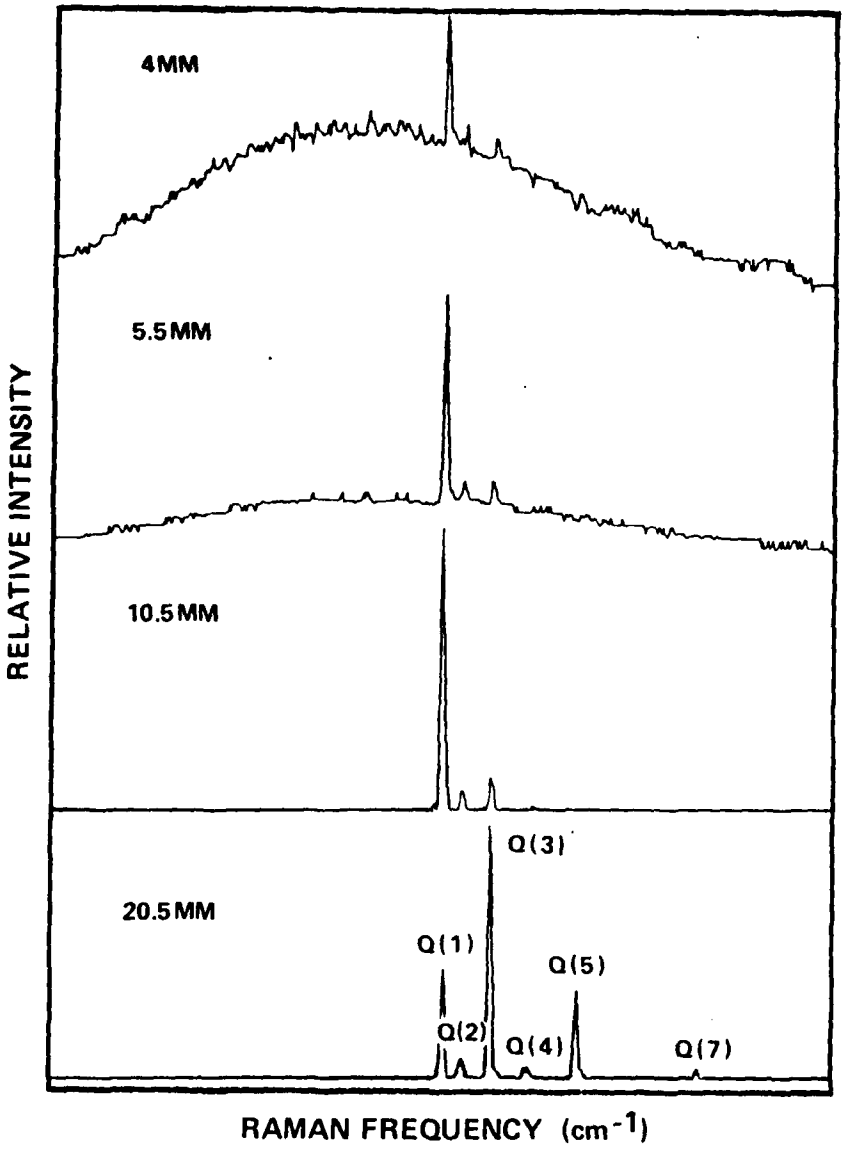

Figure 1. CARS spectra of the H_2 Q branch ($v'' = 0$) in a $\phi = 3.2$ CH_4/N_2O flame

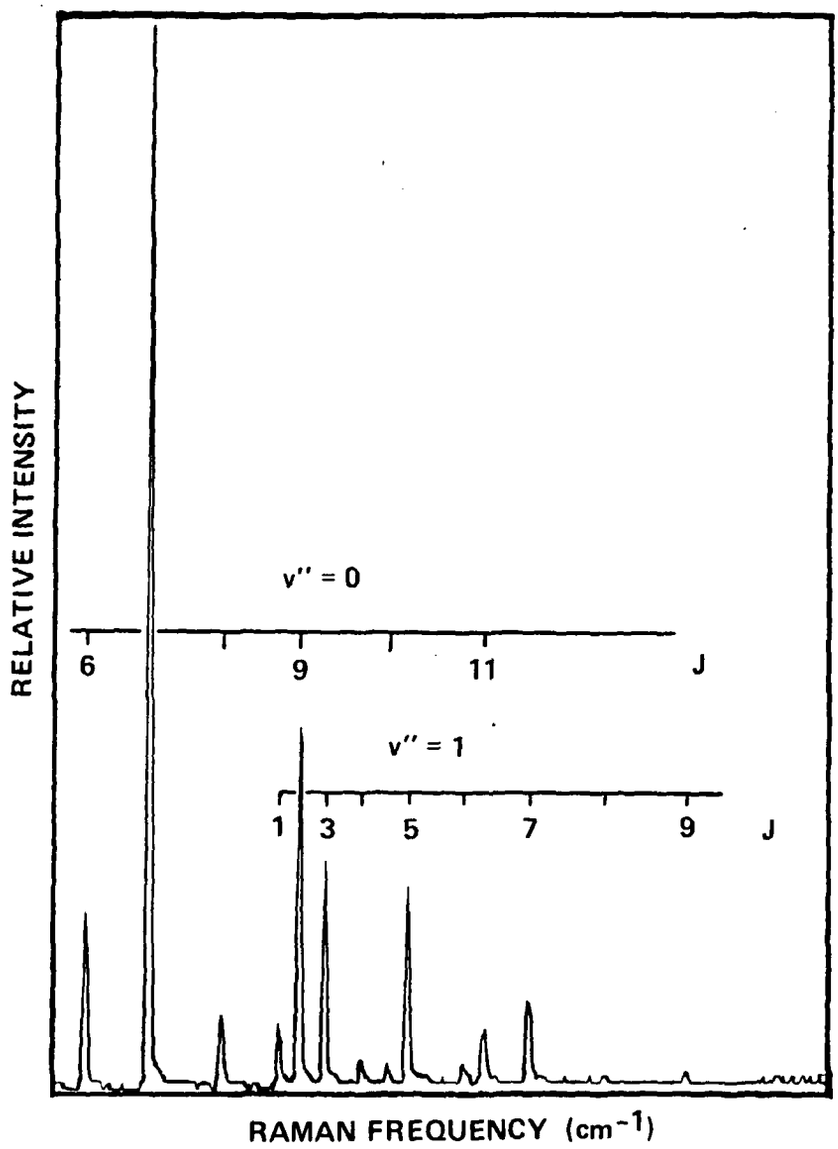

Figure 2. CARS spectra of the H_2 Q branch (v'' - 0 and v'' = 1) in a ϕ = 1.8 CH_4/N_2O flame

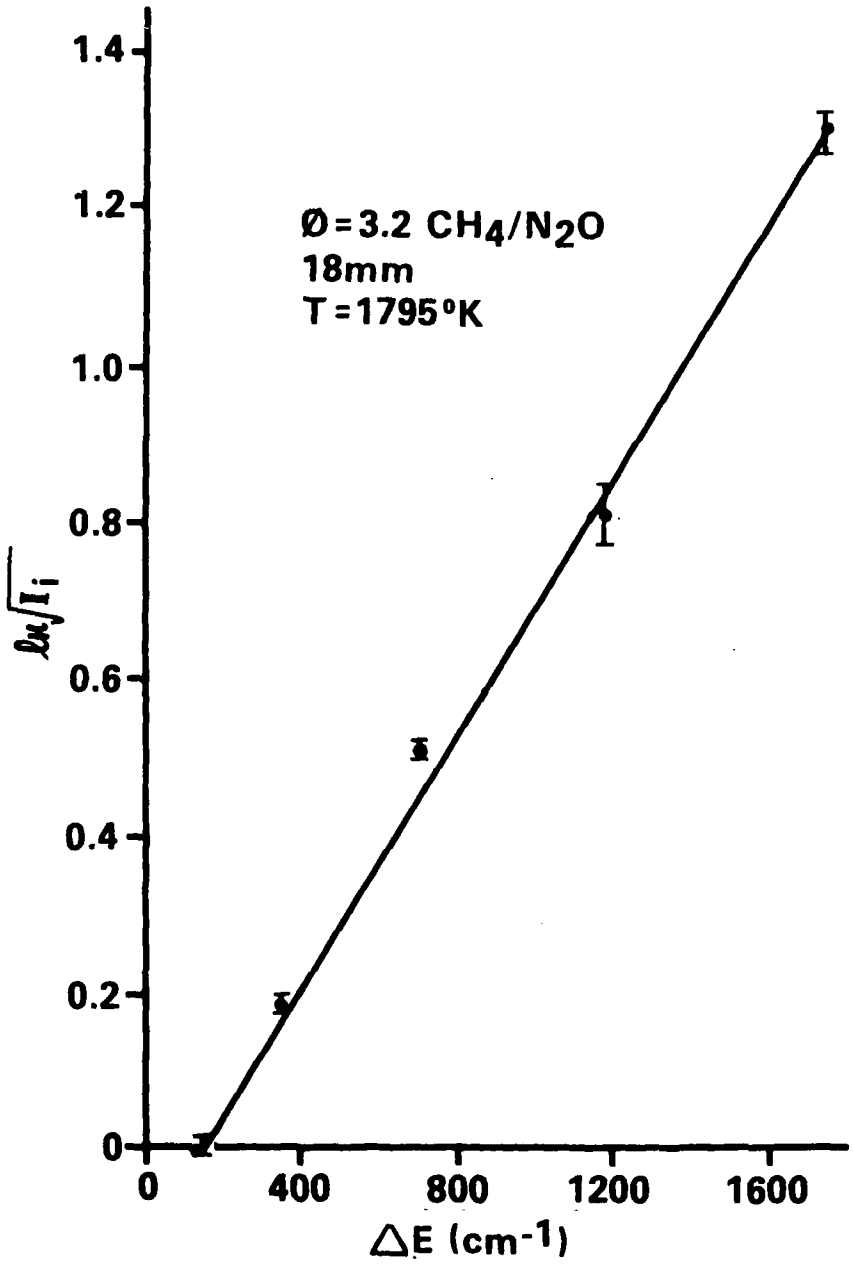

Figure 3. $\ln \sqrt{I_i}$ versus ΔE (cm^{-1}) for a ϕ = 3.2 CH_4/N_2O flame at 18 mm above the burner surface, T = 1795K

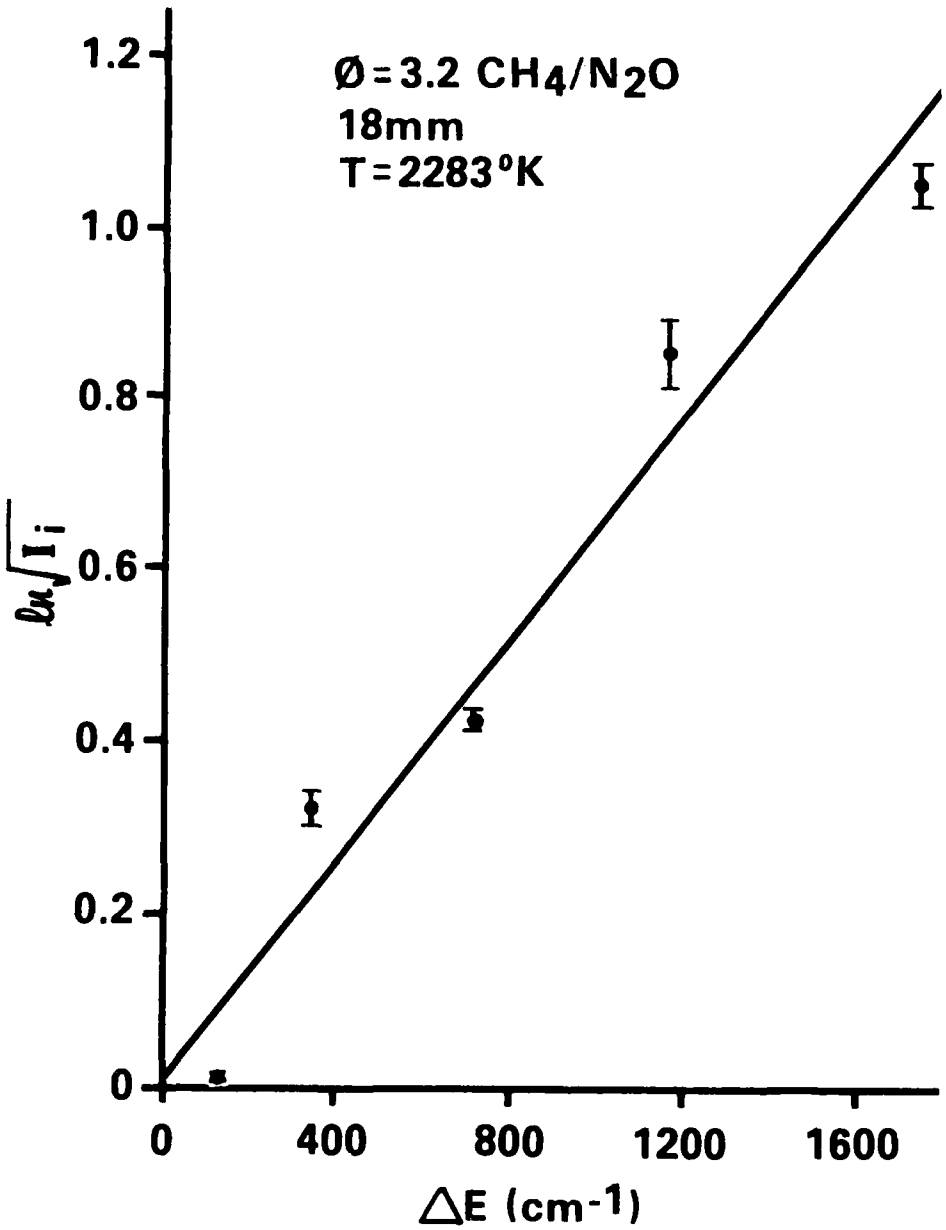

Figure 4. $\ln \sqrt{I_i}$ versus E (cm^{-1}) for a ϕ = 3.2 CH_4/N_2O flame at 18 mm above the burner surface, T = 2283K

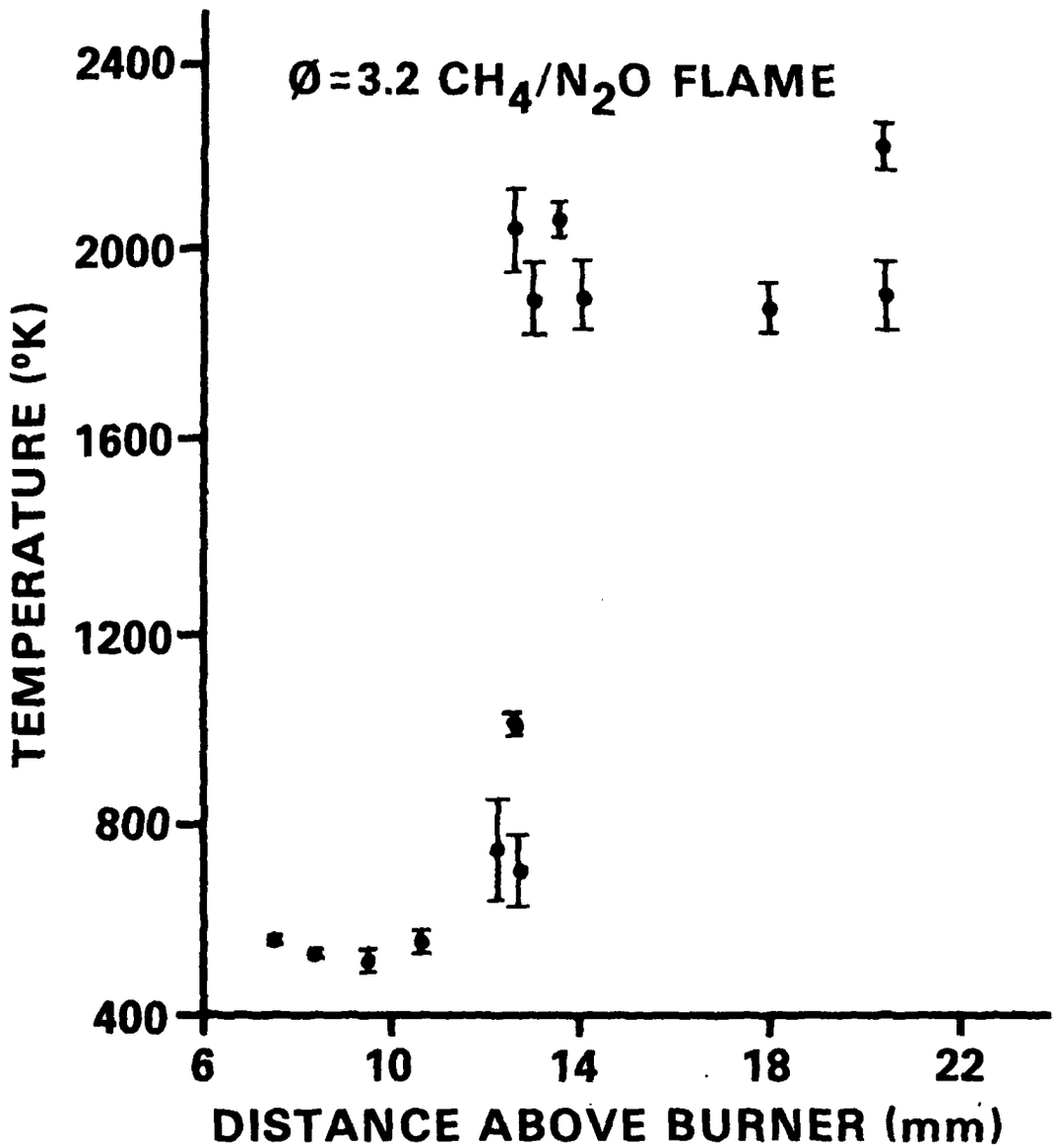

Figure 5. Temperature versus distance above burner for a $\phi = 3.2$ CH_4/N_2O flame

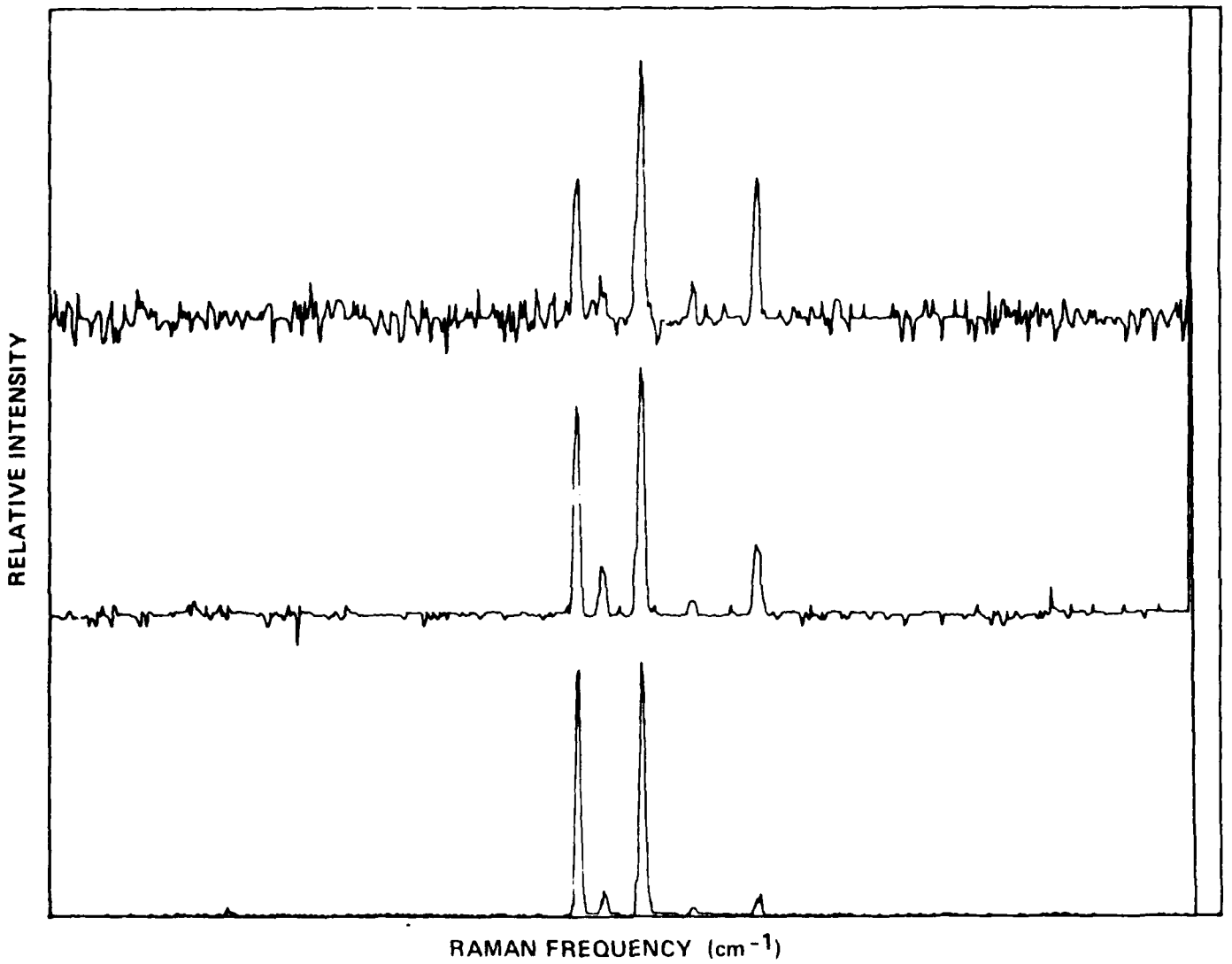

Figure 6. Time resolved (single-shot) CARS spectra of the H_2 Q branch taken at nominal 6-sec intervals after ignition and shown sequentially from bottom to top

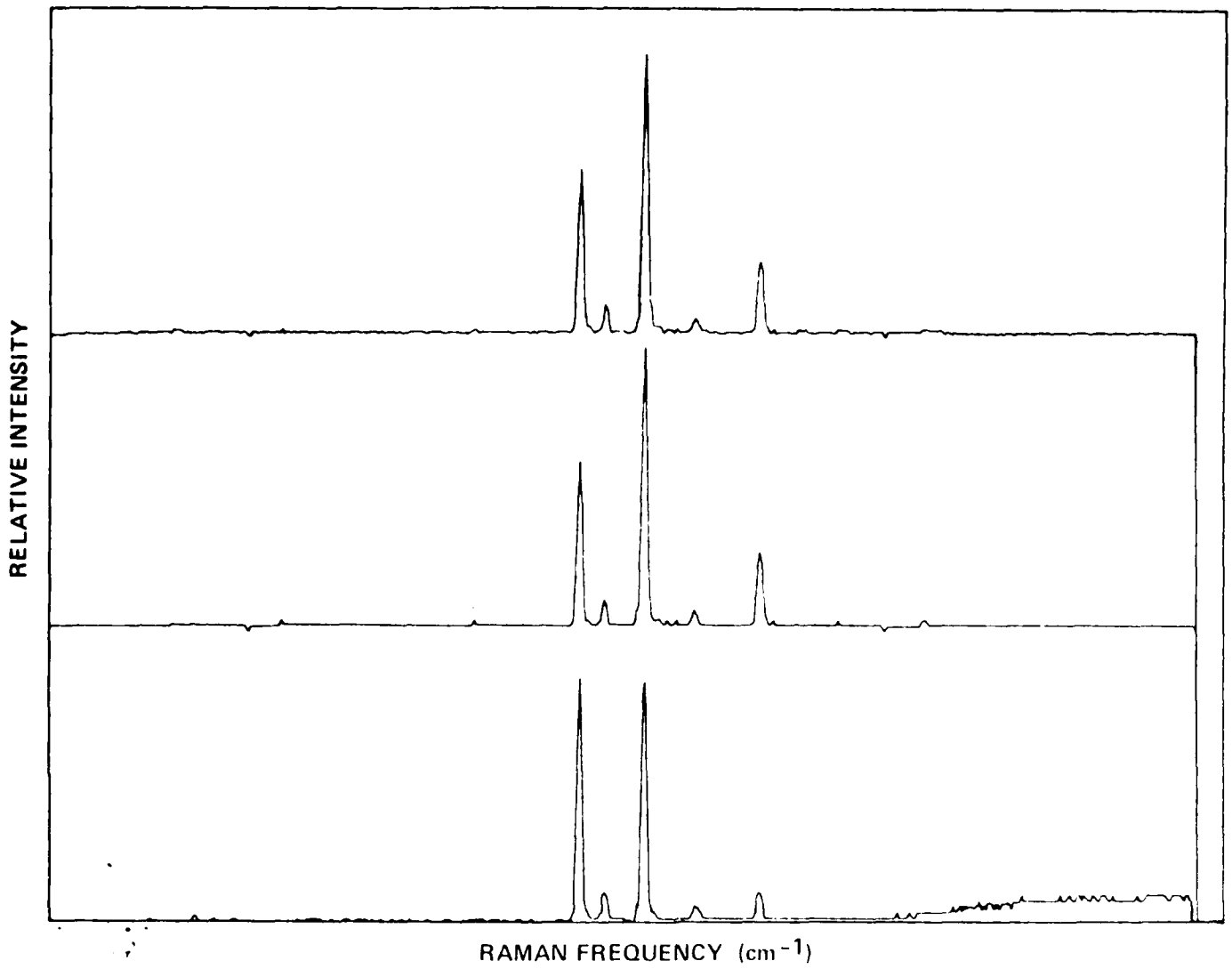

Figure 7. Time resolved (10-shot average) CARS spectra of the H_2 Q branch taken at nominal 6-sec intervals after ignition and shown sequentially from bottom to top

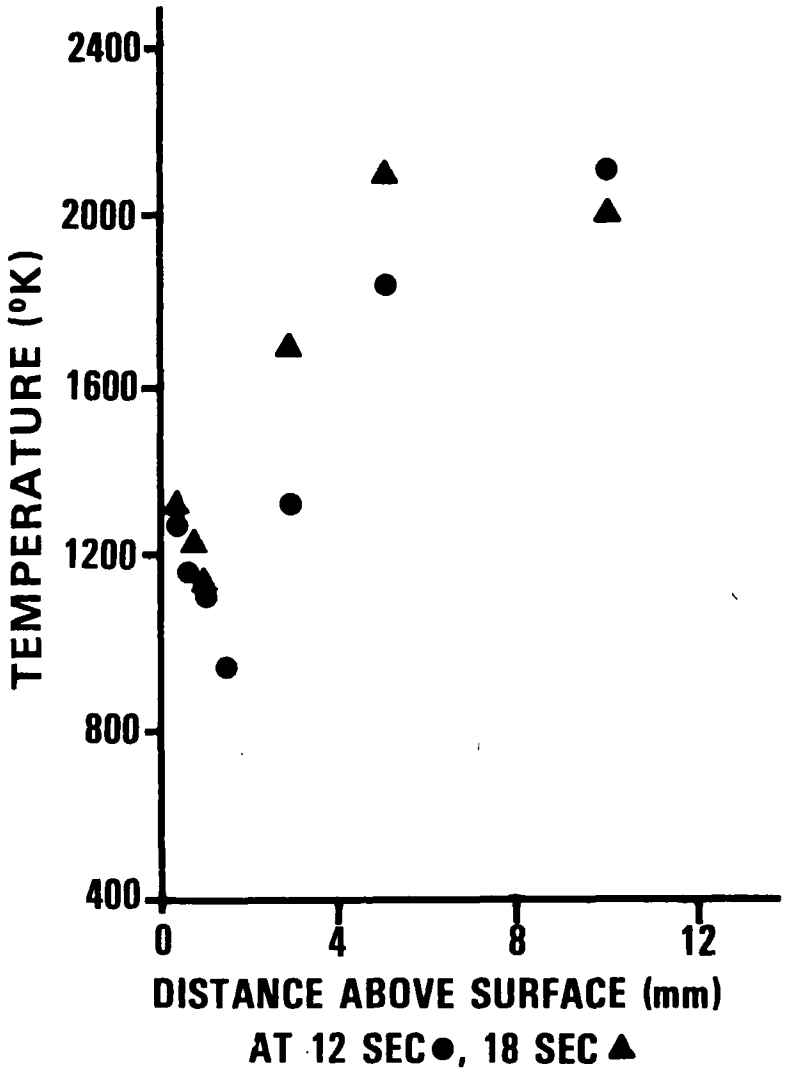

Figure 8. Temperature versus distance above surface at 12 and 18 seconds after nitramine composite ignition

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