PREDICTED NO$_2$ IR CHEMILUMINESCENCE IN THE NATURAL ATMOSPHERE

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Scientific Report No. 1

13 May 1986

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This technical report has been reviewed and is approved for publication

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Predicted NO₂ IR Chemiluminescence in the Natural Atmosphere

Results from the author's recent theoretical study of NO₂ IR Chemiluminescence from the NO + O and NO + O₂ reactions are employed in a new upper atmospheric NO₂ IR radiance model. Predictions are made for altitude-dependent and column-integrated NO₂ chemiluminescent radiance in four vibrational bands, and comparison is made with previous model predictions.
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1. INTRODUCTION

It has been proposed that NO$_2$ chemiluminescence from the NO + O and NO + O$_3$ reactions would be a significant source of atmospheric background IR radiance in a disturbed atmosphere.$^{(1,2)}$ Predicting this radiance requires a reliable model of NO$_2$ IR emission in addition to knowledge of the species concentrations. Several radiance models have been proposed, each of which uses different assumptions and gives significantly different results. Archer's model$^{(1)}$ assumes a spectral intensity used in the WOE code,$^{(3)}$ based on "theoretical considerations of energy conservation." This spectrum gives all 3 fundamental bands a comparable intensity, and also includes numerous overtone and combination bands, resulting in substantial radiance throughout the LWIR. A second model, due to Kofsky et al.,$^{(2,4)}$ is based on experimental data which observes roughly a 10:1 ratio of $v_3$ to $v_1+v_3$ intensity, with no other bands included. Finally, a more sophisticated, room temperature model for $v_3$ and $v_1+v_3$ chemiluminescence has recently been developed by the author.$^{(5,6)}$ which gives excellent agreement with laboratory data at pressures in the 0.02 - 2 torr range.

In this Report a simple NO$_2$ IR chemiluminescence model is developed for use at the pressures and temperatures found in the upper atmosphere. The model is used to calculate radiance in the natural atmosphere in the 80 - 135 km range, where comparison can be made with the predictions of Kofsky et al.$^{(2)}$ Since the radiance is proportional to the NO density, the calculation may be extrapolated to the conditions of elevated NO found in a disturbed atmosphere.

The current model is based on the author's previous work,$^{(5,6)}$ and gives similar results at laboratory pressures. At very low pressures there are factor-of-3 differences for the $v_3$ and $v_1+v_3$ bands, however, and furthermore the current model includes additional bands ($v_1$ and $v_2$) not included previously.
2. CURRENT KNOWLEDGE OF THE NO₂ CHEMILUMINESCENCE MECHANISM

As discussed elsewhere,(5,6) available data on NO₂ chemiluminescence is consistent with the following mechanism:

1. The NO + O reaction initially yields an activated complex that appears to possess energy equipartition between its electronic and vibrational modes. This complex relaxes either radiatively (via continuum emission) or collisionally to the ground electronic state, with vibrational equipartition maintained in either case as a consequence of chaotic dynamics in the excited electronic (2B₂) state. In the low pressure limit the vibrational distribution following continuum emission is characterized by an average energy of about 5000 cm⁻¹. These conclusions are derived from the shapes and absolute intensities of the continuum and IR band spectra in the NO + O reaction.

2. The NO + O₃ reaction directly yields NO₂ in vibrational equipartition. The nascent state distribution is close to a statistical RRHO distribution, with an average energy of about 6000 cm⁻¹. These conclusions are derived from trajectory calculations(7) and are fully consistent with the shapes and absolute intensities of the continuum and IR band spectra in the NO + O₃ reaction.

3. During collisional relaxation of vibrationally excited NO₂, equipartition between the v₁ and v₃ modes is approximately maintained, resulting in a -10:1 ratio of v₃ to v₁+v₃ emission. This indicates that v₁ and v₃ quanta are collisionally removed with comparable rates. This behavior is consistent with either rapid intermode V-V coupling, or similar V-R.T rates for v₁ and v₃, or, as suggested by studies of vibrational deactivation in other triatomics,(8-10) a compromise between the two mechanisms. No information is available on the relaxation of v₂. In triatomics such as O₃,(8) CO₂,(9) and OCS,(10) the rate of intermode coupling is roughly comparable to the rate of relaxation of the v₂ mode, leading to similar effective quenching rates for all 3 vibrational modes. Similar behavior would be expected for NO₂.
4. NO$_2$ reacts with atomic oxygen at a rate of 0.9-1.5x10$^{-11}$ cm$^3$ molec$^{-1}$ sec$^{-1}$, depending on vibrational excitation. Using typical N$_2$, O$_2$ and O atom densities in the 100-135 km altitude range, where NO$_2$ vibrational emission should peak, and assuming NO$_2$ vibrational quenching rates measured in the laboratory for N$_2$ and O$_2$.\textsuperscript{(11,12)} it is seen that most of the collisional removal of NO$_2$(v) proceeds via reaction with O rather than physical quenching. This means that even if the assumption of equal quenching rates for the 3 vibrational modes is inaccurate, the impact on the column-integrated atmospheric radiance is small.

5. The Einstein A coefficients for emission are consistent with harmonic oscillator scaling rules; i.e., A - v for fundamental bands. In this approximation, and neglecting overtone and combination bands, the rate constant for radiative removal of energy from a given mode is simply the Einstein coefficient of the v=1 fundamental.
3. MODEL DESCRIPTION

The model employs equations similar to Archer's Eq. (B-4), (1)

$$I(j) = [NO][O](k_2 \sum_{j} k_3[M_j]) Q(j)$$

for the NO + O reaction, and

$$I(j) = [NO][O_3] k_1 Q(j)$$

for the NO + O reaction, where $I(j)$ is the radiance of band $j$ in photons/cm$^3$, $k_1$ is the NO + O$_3$ rate constant, $k_2$ and $k_3$ are the 2-body and 3-body NO + O rate constants, $Q(j)$ is the quantum yield, and $Q(j)$ is the quenching factor, given by

$$Q(j) = A_j / (\sum_{j} q_i[M_i] + k_0[O] + A_j)$$

where $A_j$ is the Einstein coefficient of band $j$, $q_i$ is the vibrational quenching rate constant for species $M_i$, and $k_0$ is the rate of reaction with O atoms, taken as $1 \times 10^{-11}$ cm$^3$ molec$^{-1}$ sec$^{-1}$.

The rate constants $k_1$, $k_2$ and $k_3$ have room temperature values given previously. (5) The temperature dependences are taken from Sharp(13) for the NO + O reaction, and Borders and Birks(14) for the NO + O$_3$ reaction.

Assuming that 5000 cm$^{-1}$ of vibrational energy is deposited in NO$_2$ in both reactions, and assuming equipartition among the 3 vibrational modes, the quantum yields are 1.1, 2.2, and 1.0 for $v_1$, $v_2$, and $v_3$, respectively. $q_i$ is taken as $5 \times 10^{-13}$ cm$^3$ molec$^{-1}$ sec$^{-1}$ for $M = O_2$ and $2 \times 10^{-13}$ for $M = N_2$. (11, 12) The Einstein coefficients are taken as 110 sec$^{-1}$ for $v_3$ (Ref. 5) and 0.02 sec$^{-1}$ for $v_2$, based on the measured band strength. (15) The Einstein coefficient for the very weak $v_1$ band is unknown, since a reliable
band strength is not available. For the sake of argument we take the value 0.01 sec\(^{-1}\), which probably is an upper limit.

Unfortunately, we do not have a simple way of estimating effective quantum yields for overtone or combination bands. For the present we shall take the \(v_1^2 + v_3\) band intensity as 0.1 of the \(v_3\) intensity, consistent with both experiments and modeling at laboratory pressures. This ratio might change slightly under upper atmospheric conditions, however. We assume that emission from other 2-quantum bands is much smaller than \(v_1^2 + v_3\) emission, and can be neglected.

The current formulation is fully consistent with the previous model calculations\(^{(5)}\) at laboratory pressures. However, it yields factor-of-3 smaller \(v_3\) and \(v_1^2 + v_3\) radiances in the low pressure limit. This is because the previous model assumed essentially instantaneous intermode vibrational mixing, causing virtually all of the emitted vibrational energy to appear in \(v_3\). The present model represents the opposite point of view, and assumes that the energy initially apportioned to each mode (i.e., \(1/3\) of the available vibrational energy) cannot be radiated by another mode. This also appears to be the point of view of the WOE spectrum used by Archer. Put another way, the previous model assumed an effective \(v_3\) quantum yield 3x as large, and an effective Einstein coefficient \(1/3\) as large. The neglect of intermode mixing in the present model means that the \(v_3\) radiance may be underestimated and that the predicted \(v_1\) and \(v_2\) radiances may be overestimated in the low pressure, low O atom limit. However, as explained in paragraph 4, of Section 2, the column-integrated \(NO_2\) radiance should be insensitive to details of vibrational relaxation.
4. RESULTS

\(v_1, v_2, v_3\) and \(v_1-v_3\) band radiances were calculated using the altitudes and nighttime O, NO and O\textsubscript{3} densities shown in Table II-3 of Kofsky et al.\(^{(2)}\) The \(N_2\) and \(O_2\) densities and temperature were taken from the U. S. Standard Atmosphere.\(^{(16)}\) The input data are given in Table 1, and the computed radiances are given in Table 2. Integrated zenith and 95 km limb radiances were computed for the \(v_3, v_1+v_3,\) and \(v_2\) bands, and are given in Table 3 in Kofsky's units of \(w/\text{km}^2\text{-sr}.\) For the \(v_3\) and \(v_1+v_3\) bands these values turn out to be a factor of 3 to 5 lower than Kofsky's. We note that Kofsky assumed slightly different \(v_3\) and \(v_1+v_3\) quantum yields, and neglected \(\text{NO}_2(v)\) quenching and reactive removal as well as temperature dependence of \(k_2.\) Both studies find that the contribution from the \(\text{NO} + \text{O}_3\) reaction is negligible at these altitudes.

<table>
<thead>
<tr>
<th>Altitude/km</th>
<th>(N_2/\text{cm}^{-3})</th>
<th>(O_2/\text{cm}^{-3})</th>
<th>(O/\text{cm}^{-3})</th>
<th>(O_3/\text{cm}^{-3})</th>
<th>(\text{NO}/\text{cm}^{-3})</th>
<th>(T/^\circ\text{K})</th>
</tr>
</thead>
<tbody>
<tr>
<td>80</td>
<td>3.0E14</td>
<td>8.1E13</td>
<td>3.6E10</td>
<td>2.8E8</td>
<td>8.6E6</td>
<td>199</td>
</tr>
<tr>
<td>90</td>
<td>5.5E13</td>
<td>1.5E13</td>
<td>2.6E11</td>
<td>2.8E8</td>
<td>1.5E7</td>
<td>187</td>
</tr>
<tr>
<td>100</td>
<td>9.2E12</td>
<td>2.2E12</td>
<td>5.6E11</td>
<td>3.7E7</td>
<td>4.4E7</td>
<td>195</td>
</tr>
<tr>
<td>110</td>
<td>1.6E12</td>
<td>2.6E11</td>
<td>2.6E11</td>
<td>2.9E7</td>
<td>5.7E7</td>
<td>240</td>
</tr>
<tr>
<td>120</td>
<td>3.7E11</td>
<td>4.4E10</td>
<td>1.5E11</td>
<td>2.5E7</td>
<td>3.7E7</td>
<td>360</td>
</tr>
<tr>
<td>135</td>
<td>8.7E10</td>
<td>8.6E9</td>
<td>4.6E10</td>
<td>2.8E4</td>
<td>8.6E6</td>
<td>517</td>
</tr>
</tbody>
</table>
### TABLE 2. CALCULATED IR BAND RADIANCE (photons/cm$^3$-sec)

<table>
<thead>
<tr>
<th>Alt./km.</th>
<th>$v_1$</th>
<th>$v_2$</th>
<th>$v_3$</th>
<th>$v_1+v_3$</th>
<th>$v_1$</th>
<th>$v_2$</th>
<th>$v_3$</th>
<th>$v_1+v_3$</th>
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</thead>
<tbody>
<tr>
<td>80</td>
<td>0.0</td>
<td>0.1</td>
<td>12</td>
<td>1.2</td>
<td>0.0</td>
<td>0.0</td>
<td>2</td>
<td>0.2</td>
</tr>
<tr>
<td>90</td>
<td>0.3</td>
<td>0.9</td>
<td>37</td>
<td>3.7</td>
<td>0.0</td>
<td>0.1</td>
<td>2</td>
<td>0.2</td>
</tr>
<tr>
<td>100</td>
<td>1.6</td>
<td>5.3</td>
<td>92</td>
<td>9.2</td>
<td>0.0</td>
<td>0.1</td>
<td>3</td>
<td>0.3</td>
</tr>
<tr>
<td>110</td>
<td>1.6</td>
<td>5.3</td>
<td>31</td>
<td>3.1</td>
<td>0.0</td>
<td>0.1</td>
<td>1</td>
<td>0.1</td>
</tr>
<tr>
<td>120</td>
<td>1.9</td>
<td>5.9</td>
<td>17</td>
<td>1.7</td>
<td>0.0</td>
<td>0.1</td>
<td>0</td>
<td>0.0</td>
</tr>
<tr>
<td>135</td>
<td>0.7</td>
<td>2.0</td>
<td>3</td>
<td>0.3</td>
<td>0.0</td>
<td>0.0</td>
<td>0</td>
<td>0.0</td>
</tr>
</tbody>
</table>

### TABLE 3. CALCULATED COLUMN-INTEGRATED RADIANCE (w/km$^2$-sr)

<table>
<thead>
<tr>
<th></th>
<th>$v_2$</th>
<th>$v_3$</th>
<th>$v_1+v_3$</th>
</tr>
</thead>
<tbody>
<tr>
<td>zenith</td>
<td>0.00025</td>
<td>0.005</td>
<td>0.0008</td>
</tr>
<tr>
<td>Ref. 2</td>
<td>0.02</td>
<td>0.002</td>
<td></td>
</tr>
<tr>
<td>95 km limb</td>
<td>0.008</td>
<td>0.2</td>
<td>0.03</td>
</tr>
<tr>
<td>Ref. 2</td>
<td>1.0</td>
<td>0.1</td>
<td></td>
</tr>
</tbody>
</table>
5. DISCUSSION

Although the present model uses equations similar to Archer's, and yields a similar spectrum in the collision-free limit, a very important difference is the use of different Einstein coefficients (hence, different quenching factors) for each band. This results in spectra which are strongly dominated by \( v_3 \) and \( v_1-v_3 \), as predicted by Kofsky et al. The WOE-type spectrum, which shows comparable emission in all bands, is achieved only at the highest altitudes, where the intensity is weak due to the falloff of NO and O densities.

A qualitative explanation for the altitude dependence of the vibrational band radiance in the 100-120 km range is as follows. Because of the large \( v_3 \) Einstein coefficient, this band is essentially unquenched. Hence it roughly tracks the \([O][NO]\) product, which peaks at 110 km. On the other hand, the \( v_1 \) and \( v_2 \) bands are efficiently removed by O atoms, hence these bands roughly track \([NO]\), which is nearly constant at these altitudes. At lower and higher altitudes different considerations apply. At 135 km all the bands are unquenched. Below 100 km all the bands are quenched, and in addition the 3-body contribution to the NO - O reaction is important.

Since the \( v_3 \) and \( v_1+v_3 \) radiances are governed by parameters which are ultimately tied to laboratory data, the uncertainty in these bands is probably comparable to the factor of 2 estimated in the previous model study of NO\(_2\) vibrational emission.\(^{(5)}\) The uncertainty for the \( v_2 \) band is probably larger, since this band has not been observed in chemiluminescence. The uncertainty arises mainly in our estimate of the \( v_2 \) quantum yield, which may lead to an overall uncertainty of perhaps a factor of 4, the lower end of the range being more probable than the higher end.

The \( v_1 \) band radiance is even more uncertain because of the unknown Einstein coefficient, but in any case it will be small enough to be of negligible importance.
The $v_3$ and $v_1-v_3$ band radiances may be combined with experimental or calculated spectral shapes (e.g., Ref. 5) to yield approximate radiances per micron. It should be noted that these spectra, particularly $v_1\cdot v_3$, are broadened and red shifted from the corresponding cold bands. The $v_2$ band should exhibit similar effects.
6. REFERENCES


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