ELECTROMAGNETIC SCATTERING FROM A HOMOGENEOUS
BODY OF REVOLUTION

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ABSTRACT

This report considers plane-wave scattering by a homogeneous material body of revolution. The problem is formulated in terms of equivalent electric and magnetic currents over the surface which defines the body. Application of the boundary conditions leads to four simultaneous surface integral equations to be satisfied by the two unknown equivalent currents, electric and magnetic. The set of four equations is reduced to a coupled pair of equations by taking linear combinations of the original four equations. Because many pairs of linear combinations are possible, there are many surface integral equation formulations for the problem. Two formulations commonly encountered in the literature are discussed and solved by the method of moments. A general computer program for material bodies of revolution is developed, listed, and documented. Examples of numerical computations are given for dielectric spheres and a finite dielectric cylinder. The computed results for the sphere are compared to the exact series solution obtained by separation of variables.
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I. INTRODUCTION

The problem of plane-wave scattering by a homogeneous material body of revolution is formulated in terms of equivalent electric and magnetic currents over the body surface. Application of boundary conditions leads to a set of four integral equations to be satisfied. Linear combinations of these four equations lead to a coupled pair of equations to be solved. One choice of combination constants gives the formulation described by Poggio and Miller [1]. This formulation has been applied to material cylinders by Chang and Harrington [2], and to material bodies of revolution by Wu [3]. We will call this choice the PMCHW formulation (formed by the initials of the above cited investigators).

Another choice of combination constants gives the formulation obtained by Müller [4]. This formulation has been applied to dielectric cylinders by Solodukhov and Vasil'ev [5] and by Morita [6], and to bodies


of revolution by Vasil'ev and Materikova [7]. We will call this choice the Müller formulation. Conditions for the uniqueness of solutions are established in terms of the combination constants. It is found that solutions to both the PMCHW formulation and to Müller's formulation are unique at all frequencies.

Numerical solutions to the coupled pair of equations are obtained by the method of moments [8]. It is relatively easy to obtain numerical solutions to these equations because the required operators are the same as those evaluated in earlier reports [9, 10]. An exemplary computer program capable of obtaining the solution to both the PMCHW formulation and the Müller formulation is described and listed. This is a main program which uses subroutines similar to those in [10] to compute the equivalent electric and magnetic currents and the two principal plane scattering patterns for a loss-free homogeneous body of revolution excited by an axially incident electromagnetic plane wave. Computed results for the equivalent currents and principal plane scattering patterns of a dielectric sphere whose relative dielectric constant is four show reasonable agreement between our solution to the PMCHW formulation, our solution to the Müller formulation, and the "exact" series [11] solution in the resonance region. Computer program subroutines which calculate the "exact" series solution for perfectly conducting spheres as well as for loss-free homogeneous spheres will be described and listed in a subsequent report.


II. SURFACE INTEGRAL EQUATION FORMULATION

An electromagnetic field propagating in a homogeneous medium of permeability $\mu_e$ and permittivity $\varepsilon_e$ is incident on the surface $S$ of a homogeneous obstacle of permeability $\mu_d$ and permittivity $\varepsilon_d$. The subscript $e$ denotes exterior medium and the subscript $d$ denotes diffracting medium. We wish to calculate the scattered electromagnetic field $\mathbf{E}_s^s, \mathbf{H}_s^s$ outside $S$ and the diffracted electromagnetic field $\mathbf{E}_d^i, \mathbf{H}_d^i$ inside $S$ in terms of the electromagnetic field $\mathbf{E}_e^i, \mathbf{H}_e^i$ which would exist on $S$ in the absence of the obstacle. This original problem is shown in Fig. 1 where $\mathbf{J}_e^i, \mathbf{M}_e^i$ are the electric and magnetic sources of $\mathbf{E}_e^i, \mathbf{H}_e^i$ and $\mathbf{n}$ is the unit normal vector which points outward from $S$.

The equivalence principle (stated in Appendix A) is used to piece together an outside situation consisting of medium $\mu_e, \varepsilon_e$ and field $\mathbf{E}_s^s, \mathbf{H}_s^s$ outside $S$ and an inside situation consisting of medium $\mu_e, \varepsilon_e$ and field $-\mathbf{E}_e^i, -\mathbf{H}_e^i$ inside $S$. This composite situation is shown in Fig. 2. Since $\mathbf{E}_s^s, \mathbf{H}_s^s$ is source-free outside $S$ and $\mathbf{E}_e^i, \mathbf{H}_e^i$ is source-free inside $S$, the only sources in Fig. 2 are the equivalent electric surface current $\mathbf{J}$ and the equivalent magnetic surface current $\mathbf{M}$ on $S$.

As a second application of the equivalence principle, we combine an outside situation consisting of medium $\mu_d, \varepsilon_d$ and zero field with an inside situation consisting of medium $\mu_d, \varepsilon_d$ and field $\mathbf{E}_d^i, \mathbf{H}_d^i$. This combination of situations is shown in Fig. 3. Since $\mathbf{E}_e^i, \mathbf{H}_e^i$ is source-free inside $S$, the only sources in Fig. 3 are the equivalent electric surface current $-\mathbf{J}$ and the equivalent magnetic surface current $-\mathbf{M}$ on $S$. By using (A-1) and (A-2) to express the surface currents in terms of the discontinuities of the tangential fields across $S$ and by using

$$\mathbf{n} \times \mathbf{E} = \mathbf{n} \times (\mathbf{E}_e^s + \mathbf{E}_e^i) \quad \text{(1)}$$

$$\mathbf{n} \times \mathbf{H} = \mathbf{n} \times (\mathbf{H}_e^s + \mathbf{H}_e^i) \quad \text{(2)}$$

on $S$, the interested reader can verify that the surface currents in Fig. 3 are indeed the negatives of those in Fig. 2. Equations (1) and (2) are the boundary conditions that the tangential components of the fields in the
Fig. 1. Original problem.

Fig. 2. Outside equivalence.

Fig. 3. Inside equivalence.
original problem as shown in Fig. 1 are continuous across $S$.

The scattered field $E^s$, $H^s$ outside $S$ and the diffracted field $E$, $H$ inside $S$ could easily be calculated if $J$ and $M$ were known because the media into which $J$ and $M$ radiate is homogeneous in Figs. 2 and 3. We have to determine $J$ and $M$. The equivalence principle states that there exist $J$ and $M$ which radiate the fields in Figs. 2 and 3, but the equivalence principle does not tell what $J$ and $M$ are. The equivalence principle does state that

$$J = n \times H$$

$$M = E \times n$$

but this is not very useful because $E$ and $H$ are unknown.

From Figs. 2 and 3,

$$-n \times E_e^- = n \times E_i^-$$

$$-n \times H_e^- = n \times H_i^-$$

$$-n \times E_d^+ = 0$$

$$-n \times H_d^+ = 0$$

where

$E_e^-$ is the electric field just inside $S$ due to $J$, $M$, radiating in $\epsilon_e$

$H_e^-$ is the magnetic field just inside $S$ due to $J$, $M$, radiating in $\epsilon_e$

$E_d^+$ is the electric field just outside $S$ due to $J$, $M$, radiating in $\epsilon_d$

$H_d^+$ is the magnetic field just outside $S$ due to $J$, $M$, radiating in $\epsilon_d$

The equivalent currents $J$, $M$ which appear in Figs. 2 and 3 satisfy (5)-(8) because (5)-(8) were obtained from Figs. 2 and 3. It is shown in Appendix B that the solution to (5)-(8) is unique. Therefore, (5)-(8) uniquely determine the equivalent currents $J$, $M$ of Figs. 2 and 3.

Equations (5)-(8) form a set of four equations in the two unknowns $J$ and $M$. The usual methods of equation solving apply only when the number of equations is equal to the number of unknowns. We want to reduce the set of four equations (5)-(8) to two equations. One way to do this is to form the linear combination

5
\[-n \times \left( E_e^- + \alpha E_d^+ \right) = n \times E_i \]  \hspace{1cm} (9)

of (5) and (7) and the linear combination

\[-n \times \left( H_e^- + \beta H_d^+ \right) = n \times H_i \]  \hspace{1cm} (10)

of (6) and (8) where \( \alpha \) and \( \beta \) are complex constants.

The solution \( J, M \) to (5)-(8) satisfies (9) and (10). This \( J, M \) will be the only solution to the pair of equations (9) and (10) if

\[-n \times \left( E_e^- + \alpha E_d^+ \right) = 0 \]  \hspace{1cm} (11)

\[-n \times \left( H_e^- + \beta H_d^+ \right) = 0 \]  \hspace{1cm} (12)

have only the trivial solution \( J = M = 0 \). From (11) and (12),

\[ P_e = - \alpha \beta^* P_d \]  \hspace{1cm} (13)

where \( P_e \) is the complex power flow of \( E_e^- \), \( H_e^- \) inside \( S \) and \( P_d \) is the complex power flow of \( E_d^+, H_d^+ \) outside \( S \). The asterisk in (13) denotes complex conjugate. If \( \alpha \beta^* \) is real, then the real part of (13) reduces to

\[ \text{Real}(P_e) = - \alpha \beta^* \text{Real}(P_d) \]  \hspace{1cm} (14)

If \( \alpha \beta^* \) is not only real but also positive, then

\[ \text{Real}(P_d) = 0 \]  \hspace{1cm} (15)

because both \( \text{Real}(P_e) \) and \( \text{Real}(P_d) \) are greater than or equal to zero. Equation (15) implies that

\[ n \times E_d^+ = n \times H_d^+ = 0 \]  \hspace{1cm} (16)

Substitution of (16) into (11) and (12) yields

\[ n \times E_e^- = n \times H_e^- = 0 \]  \hspace{1cm} (17)
The system of equations (16) and (17) is precisely the homogeneous system of equations associated with (5)-(8). It was shown in Appendix B that this homogeneous system of equations has only the trivial solution \( \mathbf{J} = \mathbf{M} = 0 \). Therefore, if \( \alpha \beta^* \) is real and positive, then the coupled pair of equations (11) and (12) has only the trivial solution \( \mathbf{J} = \mathbf{M} = 0 \) so that the solution \( \mathbf{J}, \mathbf{M} \) to (5)-(8) is the only solution to the coupled pair of equations (9) and (10).

If \( \alpha = \beta = 1 \), then (9) and (10) become

\[
-\mathbf{n} \times (\mathbf{E}^- + \mathbf{E}^+) = \mathbf{n} \times \mathbf{E}^i \tag{18}
\]

\[
-\mathbf{n} \times (\mathbf{H}^- + \mathbf{H}^+) = \mathbf{n} \times \mathbf{H}^i \tag{19}
\]

The set of equations (18) and (19) is the coupled pair of surface integral equations transcribed by Poggio and Miller [1]. We call these equations the PMCHW equations. Since \( \alpha = \beta = 1 \) implies that \( \alpha \beta^* \) is real and positive, the argument consisting of (11)-(17) and involving real power flow shows that (18) and (19) uniquely determine the desired \( \mathbf{J}, \mathbf{M} \) of Figs. 2 and 3.

That (18) and (19) uniquely determine \( \mathbf{J}, \mathbf{M} \) of Figs. 2 and 3 can also be shown as follows. The desired \( \mathbf{J}, \mathbf{M} \) of Figs. 2 and 3 satisfies (18) and (19) because (18) and (19) were obtained from Figs. 2 and 3.

This desired \( \mathbf{J}, \mathbf{M} \) will be the only solution to (18) and (19) if the associated set of homogeneous equations

\[
-\mathbf{n} \times (\mathbf{E}^- + \mathbf{E}^+) = 0 \tag{20}
\]

\[
-\mathbf{n} \times (\mathbf{H}^- + \mathbf{H}^+) = 0 \tag{21}
\]

has only the trivial solution \( \mathbf{J} = \mathbf{M} = 0 \).

The following argument shows that (20) and (21) have only the trivial solution \( \mathbf{J} = \mathbf{M} = 0 \). Let \( \mathbf{E}_d, \mathbf{H}_d \) be the electromagnetic field outside \( S \) due to \( \mathbf{J}, \mathbf{M} \) radiating in \( \mu_d, \varepsilon_d \). Let \( \mathbf{E}_e, \mathbf{H}_e \) be the electromagnetic field inside \( S \) due to \( \mathbf{J}, \mathbf{M} \) radiating in \( \mu_e, \varepsilon_e \). Use the equivalence principle to form the composite situation consisting of medium \( \mu_d, \varepsilon_d \) and field \( \mathbf{E}_d, \mathbf{H}_d \) outside \( S \) and medium \( \mu_e, \varepsilon_e \) and field \( -\mathbf{E}_e, -\mathbf{H}_e \) inside \( S \) as shown in Fig. 4. In Fig. 4, \( \mathbf{E}_d, \mathbf{H}_d \) is a source-free Maxwellian field outside \( S \). Since \( \mathbf{E}_e, \mathbf{H}_e \) is
Fig. 4. Composite situation used to prove that (20) and (21) have only the trivial solution $J = M = 0$.

Fig. 5. Composite situation used to prove that (26) and (27) have only the trivial solution $J = M = 0$. 
a source-free Maxwellian field inside \( S \), the field \(-E_e, -H_e\) appearing in Fig. 4 is also a source-free Maxwellian field inside \( S \). Now, (20) and (21) state that the tangential components of the field in Fig. 4 are continuous across \( S \). Thus, Fig. 4 is entirely source-free so that the field in Fig. 4 is zero everywhere in which case \((B-1)-(B-4)\) are satisfied. But, as shown in Appendix B, \((B-1)-(B-4)\) have only the trivial solution \( J = M = 0 \). Hence, (20) and (21) have only the trivial solution \( J = M = 0 \).

If

\[
\alpha = -\frac{\rho_d}{\rho_e} \quad (22)
\]

\[
\beta = -\frac{\mu_d}{\mu_e} \quad (23)
\]

then (9) and (10) become

\[
-n \times (E_e - \frac{\varepsilon_d}{\varepsilon_e} E^+_e) = n \times E^+_e \quad (24)
\]

\[
-n \times (H_e - \frac{\mu_d}{\mu_e} H^+_e) = n \times H^+_e \quad (25)
\]

The set of equations (24) and (25) is the coupled pair of surface integral equations obtained by Müller [4]. We call these equations the Müller equations. The singularity that the kernels of the integral equations (24) and (25) exhibit as the source point passes through the field point is not as pronounced as the singularity of the kernels of (18) and (19). If \( \varepsilon_e, \varepsilon_d, \mu_d \), and \( \varepsilon_d \) are real in (22)-(25), then \( \alpha \beta^* \) is real and positive. In this case, the argument consisting of (11)-(17) shows that (24) and (25) uniquely determine the desired \( J, M \) of Figs. 2 and 3.

An alternate proof, valid for lossy media, that (24) and (25) uniquely determine the desired \( J, M \) is presented. This proof is similar to the argument which used Fig. 4 to show that (18) and (19) uniquely determine the desired \( J, M \) and is as follows. The desired \( J, M \) of Figs. 2 and 3 satisfies (24) and (25) because (24) and (25) were obtained from Figs. 2 and 3. This desired \( J, M \) will be the only solution to (24) and (25) if the associated set of homogeneous equations
has only the trivial solution $J = M = 0$.

The following argument shows that (26) and (27) have only the trivial solution $J = M = 0$. Let $E_d, H_d$ be the electromagnetic field outside $S$ due to $J, M$ radiating in $\nu_d, \varepsilon_d$. Because the electromagnetic field $E_d, H_d$ is a source-free Maxwellian field in $\nu_d, \varepsilon_d$ outside $S$, the dual electromagnetic field $\eta_d H_d, -\frac{1}{\eta_d} E_d$ where

$$\eta_d = \sqrt{\frac{\nu_d}{\varepsilon_d}}$$

is also a source-free Maxwellian field outside $S$. Let $E_e, H_e$ be the electromagnetic field inside $S$ due to $J, M$ radiating in $\nu_e, \varepsilon_e$. Because the electromagnetic field $E_e, H_e$ is a source-free Maxwellian field inside $S$, the dual electromagnetic field $\eta_e H_e, -\frac{1}{\eta_e} E_e$ where

$$\eta_e = \sqrt{\frac{\nu_e}{\varepsilon_e}}$$

is also a source-free Maxwellian field inside $S$. Use the equivalence principle to form the composite situation consisting of medium $\nu_d, \varepsilon_d$ and field $\eta_d H_d, -\frac{1}{\eta_d} E_d$ outside $S$ and medium $\nu_e, \varepsilon_e$ and field

$$\sqrt{\frac{\nu_e \varepsilon_e}{\mu_d \varepsilon_d}} \left( \eta_e H_e, -\frac{1}{\eta_e} E_e \right)$$

inside $S$ as shown in Fig. 5. Now, (26) and (27) state that the tangential components of the field in Fig. 5 are continuous across $S$. Thus, Fig. 5 is entirely source-free so that the field in Fig. 5 is zero everywhere in which case (B-1)-(B-4) are satisfied. But, as shown in Appendix B, (B-1)-(B-4) have only the trivial solution $J = M = 0$. Hence, (26) and (27) have only the trivial solution $J = M = 0$. 
III. METHOD OF MOMENTS SOLUTION FOR A BODY OF REVOLUTION

In this section, a method of moments solution to (9) and (10) is developed for a homogeneous loss-free body of revolution. Special cases of (9) and (10) are the PMCHW equations (18) and (19) and the Müller equations (24) and (25).

For compatibility with equation (40) on page 14 of [9], we rewrite (9) as

\[- \frac{1}{\eta_e} (E^- + \alpha E^+) \tan = \frac{1}{\eta_e} E^\dagger \tan\]

where \(\tan\) denotes tangential components on \(S\) and \(\eta_e\) is given by (29). The fields on the left-hand sides of (30) and (10) are written as the sum of fields due to \(J\) and fields due to \(M\). Advantage is taken of the fact that the operator which gives the electric field due to a magnetic current is the negative of the operator which gives the magnetic field due to an electric current and that the operator which gives the magnetic field due to a magnetic current is the square of the reciprocal of the intrinsic impedance times the operator which gives the electric field due to an electric current. In view of the above considerations, (30) and (10) become

\[- \frac{1}{\eta_e} \hat{E}_e(J) + \frac{1}{\eta_e} \hat{H}_e(M) - \frac{a}{\eta_e} \hat{E}_d(J) + \frac{a}{\eta_e} \hat{H}_d^+(M) \tan = \frac{1}{\eta_e} E^\dagger \tan \]

\[- \eta \times (\hat{H}_e(J) + \frac{1}{2} \hat{H}_e(M) + \beta \hat{H}_d^+(J) + \frac{\beta}{2} \hat{H}_d^+(M)) = \eta \times H^\dagger \]

where \(\hat{E}\) denotes the operator which gives the electric field due to an electric current. The subscript \(e\) or \(d\) on \(\hat{E}\) denotes radiation in either \(\mu_e, \varepsilon_e\) or \(\mu_d, \varepsilon_d\). The superscript + or -, if present on \(\hat{E}\), denotes field evaluation either just outside \(S\) or just inside \(S\). The \(\hat{H}\)'s in (31) and (32) are the corresponding magnetic field due to electric current operators. We stress that all \(\hat{E}\)'s and \(\hat{H}\)'s in (31) and (32) are, by definition, operators which give electric and magnetic fields due to electric currents, even though these operators act on both electric and magnetic currents \(J\) and \(M\) in (31) and (32).
Let

\[ J = \sum_{n=-\infty}^{\infty} \sum_{j=1}^{N} (I_{nj}^{t} J_{nj}^{t} + I_{nj}^{\phi} J_{nj}^{\phi}) \]  \hspace{1cm} (33) \]

\[ M = n_e \sum_{n=-\infty}^{\infty} \sum_{j=1}^{N} (V_{nj}^{t} J_{nj}^{t} + V_{nj}^{\phi} J_{nj}^{\phi}) \]  \hspace{1cm} (34) \]

where \( I_{nj}^{t}, I_{nj}^{\phi}, V_{nj}^{t}, \) and \( V_{nj}^{\phi} \) are coefficients to be determined and

\[ J_{nj}^{t} = u_t f_j(t)e^{jn\phi} \]  \hspace{1cm} (35) \]

\[ J_{nj}^{\phi} = u_\phi f_j(t)e^{jn\phi} \]  \hspace{1cm} (36) \]

In (35) and (36), \( t \) is the arc length along the generating curve of the body of revolution and \( \phi \) is the longitudinal angle. \( u_t \) and \( u_\phi \) are unit vectors in the \( t \) and \( \phi \) directions respectively such that \( u_\phi \times u_t = n \) and \( f_j(t) \) is the scalar function of \( t \) defined on page 10 of [9]. The body of revolution and coordinate system are shown in Fig. 6. Substitution of (33) and (34) into (31) and (32) yields

\[ \sum_{n=-\infty}^{\infty} \sum_{j=1}^{N} \left\{ \left( \frac{H_e J_{nj}^{t}}{e_{-nj}} + \frac{\alpha n d}{n_e} \frac{E_t J_{nj}^{t}}{d_{-nj}} \right) \tan n_j + \left( \frac{H_e J_{nj}^{\phi}}{e_{-nj}} + \frac{\alpha n d}{n_e} \frac{E_d J_{nj}^{\phi}}{d_{-nj}} \right) \tan n_j \right\} = \frac{1}{n_e} E_t \tan \]  \hspace{1cm} (37) \]

\[ \sum_{n=-\infty}^{\infty} \sum_{j=1}^{N} \left\{ \left( \frac{H_e J_{nj}^{t}}{e_{-nj}} + \frac{\beta n d}{n_e} \frac{E_d J_{nj}^{t}}{d_{-nj}} \right) v_{nj} + \left( \frac{H_e J_{nj}^{\phi}}{e_{-nj}} + \frac{\beta n d}{n_e} \frac{E_d J_{nj}^{\phi}}{d_{-nj}} \right) v_{nj} \right\} = n \times \frac{H_t}{n} \]  \hspace{1cm} (38) \]

Define the inner product of two vector functions on \( S \) to be the integral over \( S \) of the dot product of these two vector functions. Because the field operators in (37) and (38) are the same as those considered in [9],
only the nth term of the sum (37) or (38) contributes to the inner product of (37) or (38) with either \( J_{-ni}^t \) or \( J_{-ni}^\phi \). Hence, the inner product of (37) with \( J_{-ni}^t \), \( i=1,2,...,N \), and \( J_{-ni}^\phi \), \( i=1,2,...,N \), successively and the inner product of (38) with \( J_{-ni}^t \), \( i=1,2,...,N \), and \( J_{-ni}^\phi \), \( i=1,2,...,N \), successively gives the matrix equation

\[
\begin{bmatrix}
(Y_{ne}^{\phi t} + \alpha Y_{nd}^{\phi t}) & (Y_{ne}^{\phi \phi} + \alpha Y_{nd}^{\phi \phi}) & (Z_{ne}^{tt} + \frac{\alpha_n d}{\eta_e} Z_{nd}^{tt}) & (Z_{ne}^{t\phi} + \frac{\alpha_n d}{\eta_e} Z_{nd}^{t\phi}) \\
(-Y_{ne}^{tt} - \alpha Y_{nd}^{tt}) & (-Y_{ne}^{t\phi} - \alpha Y_{nd}^{t\phi}) & (Z_{ne}^{\phi t} + \frac{\alpha_n d}{\eta_e} Z_{nd}^{\phi t}) & (Z_{ne}^{\phi \phi} + \frac{\alpha_n d}{\eta_e} Z_{nd}^{\phi \phi}) \\
(Z_{ne}^{t\phi} + \frac{\beta_n e}{\eta_d} Z_{nd}^{t\phi}) & (Z_{ne}^{\phi \phi} + \frac{\beta_n e}{\eta_d} Z_{nd}^{\phi \phi}) & (Y_{ne}^{tt} + \beta Y_{nd}^{tt}) & (Y_{ne}^{t\phi} + \beta Y_{nd}^{t\phi}) \\
(-Z_{ne}^{tt} - \frac{\beta_n e}{\eta_d} Z_{nd}^{tt}) & (-Z_{ne}^{t\phi} - \frac{\beta_n e}{\eta_d} Z_{nd}^{t\phi}) & (Y_{ne}^{\phi t} + \beta Y_{nd}^{\phi t}) & (Y_{ne}^{\phi \phi} + \beta Y_{nd}^{\phi \phi}) 
\end{bmatrix} \begin{bmatrix}
\hat{v}_{n}^t \\
\hat{v}_{n}^\phi \\
\hat{t}_{n}^t \\
\hat{t}_{n}^\phi 
\end{bmatrix} = \begin{bmatrix}
\hat{v}_{n}^t \\
\hat{v}_{n}^\phi \\
\hat{t}_{n}^t \\
\hat{t}_{n}^\phi 
\end{bmatrix}
\]  

(39)

for \( n = 0, \pm 1, \pm 2, \ldots \). In (39), \( \hat{v}_{n}^t, \hat{v}_{n}^\phi, \hat{t}_{n}^t, \) and \( \hat{t}_{n}^\phi \) are column vectors of the coefficients appearing in (33) and (34). Also,

\[
(Y_{nf}^{pq})_{ij} = - \int \int_{S} J_{ni}^p \cdot \eta_e \times H_{f}^i (J_{nj}^q) ds 
\]  

(40)

\[
(Z_{nf}^{pq})_{ij} = - \frac{1}{\eta_f} \int \int_{S} J_{ni}^p \cdot E_{f}^i (J_{nj}^q) ds 
\]  

(41)

\[
\hat{v}_{ni}^p = \frac{1}{\eta_e} \int \int_{S} J_{ni}^p \cdot E_{i}^i ds 
\]  

(42)

\[
\hat{t}_{ni}^p = \int \int_{S} J_{ni}^p \cdot \eta_e \times H_{i}^i ds 
\]  

(43)

where \( p \) may be either \( t \) or \( \phi \), \( q \) may be either \( t \) or \( \phi \), and \( f \) may be either \( e \) or \( d \). If \( p=q \) in (40), it matters whether the magnetic field \( H_{f}^i (J_{nj}^q) \) is evaluated just outside or just inside \( S \). The \( Y \)'s without
carets in (39) are given by the right-hand side of (40) in which the magnetic field is evaluated just inside S. The Y's with carets in (39) are given by the right-hand side of (40) with magnetic field evaluation just outside S.

The Y and Z submatrices on the left-hand side of (39) are the same as in equation (88) on page 24 of [9] with the reservations that the caret on Y denotes magnetic field evaluation just outside S, and the extra subscript e or d denotes radiation in either \( \nu_e \), \( \nu_d \) or \( \nu_e \), \( \nu_d \). The \( \hat{t} \) column vectors on the right-hand side of (39) are the same as in equation (88) on page 24 of [9] whereas the \( \hat{v} \) column vectors in (39) are the same as the Y's without carets in equation (88) on page 24 of [9].

The solution \( \hat{v}_n \), \( \hat{v}_n \), \( \hat{t}_n \), and \( \hat{t}_n \) to the matrix equation (39) determines the equivalent electric and magnetic currents \( J \) and \( M \) according to (33) and (34). From Fig. 2, these currents radiate in \( \nu_e \), \( \nu_d \) to produce the scattered field outside S.

IV. FAR FIELD MEASUREMENT AND PLANE WAVE EXCITATION

In this section, measurement vectors are used to obtain the far field of the equivalent surface currents \( J \) and \( M \) radiating in \( \nu_e \), \( \nu_d \). This far field is the far field scattered by the homogeneous body of revolution. For plane wave excitation, the composite vector on the right-hand side of (39) is expressed in terms of these measurement vectors.

By reciprocity,

\[
E^s \cdot I_{\nu e}^r = \iint_S (J(r) \cdot E(I_{\nu e}^r) - M(r) \cdot H(I_{\nu d}^r))ds \quad (44)
\]

where \( E^s \) is the far electric field due to \( J \) and \( M \), \( I_{\nu e}^r \) is a receiving electric dipole at the far field measurement point, \( E(I_{\nu e}^r) \) is the electric field due to \( I_{\nu e}^r \), and \( H(I_{\nu d}^r) \) is the magnetic field due to \( I_{\nu d}^r \). Both \( E(I_{\nu e}^r) \) and \( H(I_{\nu d}^r) \) are evaluated at point \( r \) on S where \( r \) is the point at which the differential portion of surface ds is located. If \( \nu_e \) is tangent to the radiation sphere,
Fig. 6. Body of revolution and coordinate system.

Fig. 7. Plane wave scattering by a dielectric body of revolution.
where \( r_r \) is the distance between the measurement point and the origin in the vicinity of \( S \). Also, \( k_r \) is the propagation vector of the plane wave coming from \( L_r \), \( k \) is the propagation constant and \( n \) is the intrinsic impedance of the medium outside \( S \). To simplify the notation in this section, we have omitted the subscript \( e \) from all parameters dependent on the medium. It is understood that all far field measurement vectors and plane wave excitation vectors depend only on the external medium \( \mu_e', \varepsilon_e' \).

Substitution of (33), (34), (45), and (46) into (44) gives

\[
E^s_\theta = \frac{-jk_r}{4\pi r_r} \sum_{n=-\infty}^{\infty} (\vec{R}_n^{\theta} \vec{V}_n + \vec{R}_n^{\phi} \vec{V}_n + \vec{R}_n^{t} \vec{T}_n + \vec{R}_n^{\phi} \vec{T}_n) e^{-jkr} \tag{47}
\]

for \( L_r = \vec{u}_\theta^r \) and

\[
E^s_\phi = \frac{-j e}{4\pi r_r} \sum_{n=-\infty}^{\infty} (-\vec{R}_n^{\theta} \vec{V}_n + \vec{R}_n^{\phi} \vec{V}_n + \vec{R}_n^{t} \vec{T}_n + \vec{R}_n^{\phi} \vec{T}_n) e^{-jkr} \tag{48}
\]

for \( L_r = \vec{u}_\phi^r \) where \( \vec{u}_\theta^r \) and \( \vec{u}_\phi^r \) are unit vectors in the \( \theta_r \) and \( \phi_r \) directions respectively. As shown in Fig. 7, \( \theta_r \) and \( \phi_r \) are the angular coordinates of the receiver location at which \( L_r \) is placed. In (47) and (48), \( E^s_\theta \) and \( E^s_\phi \) are the \( \theta_r \) and \( \phi_r \) components of \( E^s \). Also, \( \vec{V}_n, \vec{T}_n, \vec{R}_n, \) and \( \vec{R}_n^{\phi} \) are column vectors of the coefficients appearing in (33) and (34). Furthermore, \( R_{pq}^n \) is a row vector whose \( j \)th element is given by

\[
R_{pq}^n = k e^{-j n \phi} \int_5 \vec{r}_p \cdot \vec{u}_q \ e^{-j k_r \cdot \vec{r}_r} \ ds \tag{49}
\]

where \( p \) may be either \( t \) or \( \phi \) and \( q \) may be either \( \theta \) or \( \phi \). In view of (35) and (36), (49) is the same as equation (92) on page 26 of [9]. It is shown in [9] that the right-hand side of (49) does not depend on \( \phi_r \).
For plane wave incidence and expansion functions \( J^t_{\omega n} \) and \( J^\phi_{\omega n} \) given by (35) and (36), the equivalent currents (33) and (34) and the fields (47) and (48) have special forms. To obtain these forms, assume that the incident electromagnetic field \( E^i, H^i \) is either a \( \theta \) polarized field defined by

\[
E^i = k \eta \, u^t_\theta \, e^{\frac{-j k \cdot r}{r}}
\]

or a \( \phi \) polarized field defined by

\[
E^i = k \eta \, u^y_\theta \, e^{\frac{-j k \cdot r}{r}}
\]

where \( k \) is the propagation vector and, as shown in Fig. 7, \( u^t \) and \( u^y \) are unit vectors in the \( \theta \) and \( y \) directions respectively. Here, \( \theta \) is the colatitude of the direction from which the incident wave comes. \( k \) is in the \( xz \) plane. No generality is lost by putting \( k \) in the \( xz \) plane because if \( k \) were shifted out of the \( xz \) plane by an angle \( \phi \), the response would also be shifted by the same angle \( \phi \).

Substituting (50) and (51) into (42) and (43), then substituting (52) and (53) into (42) and (43), next taking advantage of the relationships

\[
J^t_{\omega n} \times n = J^\phi_{\omega n}
\]

\[
J^\phi_{\omega n} \times n = -J^t_{\omega n}
\]

which are apparent from (35), (36) and Fig. 6, then comparing the results with (49), and finally using equation (104) on page 29 of [9], we obtain
The first superscript on $V_n$ and $I_n$ in (56) is the superscript which appears on the right-hand side of (39). The second superscript on $V_n$ and $I_n$ in (56) denotes the polarization of the incident plane wave. If this second superscript is $\theta$, the $\theta$ polarized field given by (50) and (51) is incident. If this second subscript is $\phi$, the $\phi$ polarized field given by (52) and (53) is incident. The $j$th element of the column vector $\vec{R}^{pq}_n$ on the right-hand side of (56) is given by (49) with $\theta_r$ replaced by $\theta_t$. Conceding that $\theta_r$ does not appear explicitly in (49), we really mean that $\theta_r$ is replaced by $\theta_t$ after the surface integral in (49) is evaluated. In other words, $\theta_r$ is replaced by $\theta_t$ in equation (95) on page 27 of [9].

For plane wave incidence, the $+n$ and $-n$ terms in formulas (33) and (34) for the equivalent currents can be combined as follows. According to equations (102) and (103) of [9], the $Y$ and $Z$ submatrices in (39) are either even or odd in $n$. The even-odd properties in $n$ of the submatrices of the square matrix on the left-hand side of (39) are tabulated as

$$
\begin{array}{cccc}
- & + & + & - \\
+ & - & - & + \\
- & + & + & - \\
+ & - & - & + \\
\end{array}
$$

where $+$ denotes an even submatrix and $-$ denotes an odd submatrix. It follows that the submatrices of the inverse of the square matrix on the left-hand side of (39) have even-odd properties in $n$ given by
From (56) and the even-odd properties of $R_{pq}$ given by equation (104) on page 29 of [9], the column vectors $\hat{V}_n$ and $\hat{I}_n$ on the right-hand side of (39) are either even or odd in $n$. The even-odd properties of the submatrices on the left-hand side of (56) are tabulated as

$$
\begin{bmatrix}
- & + & - & + \\
+ & - & + & - \\
+ & - & + & - \\
- & + & - & + \\
\end{bmatrix}
$$

Because of the above even-odd properties of the square matrix on the left-hand side of (39) and the column vector on the right-hand side of (39), the solutions to (39) satisfy

$$
\begin{bmatrix}
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\end{bmatrix}
\begin{bmatrix}
\hat{V}_n \\
\hat{V}_n \\
\hat{V}_n \\
\hat{V}_n \\
\end{bmatrix}
= 
\begin{bmatrix}
\hat{V}^t_0 \\
\hat{V}^t_0 \\
\hat{V}^t_0 \\
\hat{V}^t_0 \\
\end{bmatrix}
\begin{bmatrix}
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\hat{t}^{t0} \\
\end{bmatrix} + 
\begin{bmatrix}
\hat{I}^t_0 \\
\hat{I}^t_0 \\
\hat{I}^t_0 \\
\hat{I}^t_0 \\
\end{bmatrix}
\begin{bmatrix}
\hat{I}_n \\
\hat{I}_n \\
\hat{I}_n \\
\hat{I}_n \\
\end{bmatrix}
$$

(57)

The first superscript on the column vectors $\hat{V}_n$ and $\hat{I}_n$ in (57) is that which appears on the column vectors $\hat{V}_n$ and $\hat{I}_n$ on the left-hand side of (39). The second superscript on the column vectors in (57) denotes either the $\theta$ or the $\phi$ polarized incident plane wave. Substitution of (57), (35), and (36) into (33) and (34) yields
For the $\theta$ polarized incident wave and

\[
\begin{align*}
\mathbf{J}_{\theta}^t &= (\mathbf{F}_{\theta}^t)^{\mathbf{u}}_{\mathbf{t}} + \sum_{n=1}^{\infty} 2(\mathbf{F}_{\theta}^t)^{\mathbf{u}}_{n} \cos(n\theta) + 2(\mathbf{F}_{\theta}^t)^{\mathbf{u}}_{n} \sin(n\theta) \\
\frac{1}{n} \mathbf{M}_{\theta}^t &= (\mathbf{V}_{\theta}^t)^{\mathbf{u}}_{\mathbf{t}} + \sum_{n=1}^{\infty} 2j(\mathbf{V}_{\theta}^t)^{\mathbf{u}}_{n} \sin(n\theta) + 2(\mathbf{V}_{\theta}^t)^{\mathbf{u}}_{n} \cos(n\theta)
\end{align*}
\]  

(58)

(59)

for the $\phi$ polarized incident wave. In (58)-(61), $\mathbf{F}$ is a row vector of the $f_j(t)$. The superscript $\theta$ or $\phi$ on $\mathbf{J}$ or $\mathbf{M}$ in (58)-(61) differentiates the equivalent currents for the $\theta$ polarized incident wave from those for the $\phi$ polarized incident wave.

The far scattered fields (47) and (48) are specialized to the $\theta$ polarized incident plane wave by appending the additional subscript $\theta$ to $\mathbf{E}$ on the left-hand sides of (47) and (48) and the additional superscript $\theta$ to $\mathbf{V}_n$, $\mathbf{v}_n$, $\mathbf{t}_n$, and $\mathbf{t}_n$ on the right-hand sides of (47) and (48). Moreover, in view of equation (104) on page 29 of [9] and (57), the $+n$ and $-n$ terms in (47) and (48) can be combined. As a result, (47) and (48) become

\[
E_{\theta 0}^t = -\frac{jkr}{4\pi r} \left\{ \mathbf{R}_{\phi 0}^{\phi \theta} \phi_{\theta 0} + \mathbf{R}_{\theta 0}^{\theta \theta} \theta_{\theta 0} + 2 \sum_{n=1}^{\infty} (\mathbf{R}_{n n}^{\theta \theta} \theta_{n n} \phi_{\theta n} + \mathbf{R}_{n n}^{\phi \phi} \phi_{\theta n} \theta_{\theta n}) \cos(n\phi) \right\} + \mathbf{R}_{n n}^{\phi \phi} \phi_{\theta n} \theta_{\theta n} \cos(n\phi) + \mathbf{R}_{n n}^{\phi \phi} \phi_{\theta n} \theta_{\theta n} \sin(n\phi)
\]  

(62)

(63)

for the $\theta$ polarized incident plane wave. Similarly, (47) and (48) become
for the φ polarized incident plane wave. The first subscript on \( E^s \) on the left-hand sides of (62)-(65) denotes the receiver polarization and the second subscript on \( E^s \) denotes the transmitter polarization.

The scattering cross section \( \sigma_{pq} \) is defined by

\[
\sigma_{pq} = \frac{4\pi^2 r^2 |E^s_{pq}|^2}{|E^i|^2}
\]

(66)

where \( p \) is either \( \theta \) or \( \phi \) and \( q \) is either \( \theta \) or \( \phi \). In (66), \( E^s_{pq} \) is a component of the scattered field given by (62)-(65) and \(|E^i|^2\) is the magnitude of the electric field of the incident plane wave. According to (50) and (52),

\[
|E^i|^2 = k\eta
\]

(67)

for both polarizations so that

\[
\sigma_{pq} = \frac{4\pi^2 r^2 |E^s_{pq}|^2}{k^2 \eta^2}
\]

(68)

Normalized versions of (68) are

\[
\frac{\sigma_{pq}}{\pi a^2} = \frac{4\pi^2 r^2 |E^s_{pq}|^2}{k^2 a^2 \eta^2}
\]

(69)

\[
\frac{\sigma_{pq}}{\lambda^2} = \frac{r^2 |E^s_{pq}|^2}{\pi \eta^2}
\]

(70)

where \( a \) is some characteristic length associated with the scatterer and \( \lambda \) is the wavelength in the external medium.
V. EXAMPLES

A computer program has been written to calculate the equivalent currents and scattering patterns for a dielectric body of revolution excited by an axially incident plane wave. This program is described and listed in Part Two. Some computational results obtained with this program are given in this section.

Figures 8 and 9 show the magnitude and phase of the normalized equivalent currents \( \frac{J_\theta}{H_y} \), \( \frac{J_\phi}{H_y} \), \( \frac{M_\theta}{E_x} \), and \( \frac{M_\phi}{E_x} \) on the surface of a dielectric sphere for which \( ka = 3 \) and \( \varepsilon_r = 4 \). Here, \( k \) is the propagation constant in free space, \( a \) is the radius of the sphere and \( \varepsilon_r \) is the relative dielectric constant of the sphere. Figure 8 represents our solution of the PMCHW formulation. Figure 9 depicts our solution of the Müller formulation.

In Figs. 8 and 9, the incident field is a plane wave traveling in the positive z direction. THETA = 0° is the forward scattering direction and THETA = 180° is the backscattering direction. The incident field is given by (50) and (51) with \( \theta_0 = 180° \). The origin \( r = 0 \) is at the center of the sphere. In Figs. 8 and 9, \( J_\theta \) is the \( u_\theta \) component of electric current (58) versus \( \theta \) in the \( \phi = 0° \) plane, \( J_\phi \) is the \( u_\phi \) component of (58) versus \( \theta \) in the \( \phi = 90° \) plane, \( M_\theta \) is the \( u_\theta \) component of magnetic current (59) versus \( \theta \) in the \( \phi = 90° \) plane, and \( M_\phi \) is the \( u_\phi \) component of (59) versus \( \theta \) in the \( \phi = 0° \) plane. For axial incidence, only the \( n=1 \) term is present in (58) and (59). The symbols \( \times \) and \( + \) denote respectively magnitude and phase of the method of moments solution for the pertinent component of the electric or magnetic current. The solid curves are the exact equivalent currents obtained from the Mie series solution [11]. The normalizing constants \( E_x \) and \( H_y \) are defined in terms of the incident field (50) and (51) by

\[
E_x = u_x \cdot \frac{E_x}{u_x} \Big|_{r=0} = -kn
\]

\[
H_y = u_y \cdot \frac{H_y}{u_y} \Big|_{r=0} = -k
\]

Fig. 8. Equivalent electric and magnetic currents for dielectric sphere, $ka = 3$, $\varepsilon_r = 4$, PMCHW solution. Symbols $\times$ and $+$ denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. 6. Equivalent electric and magnetic currents for dielectric sphere, $ka = 3$, $\varepsilon = 4$, Müller solution. Symbols × and + denote magnitude and phase respectively. Solid line denotes exact solution.
where \( k \) and \( n \) are respectively the propagation constant of free space and the impedance of free space.

The currents of Figs. 8 and 9 were obtained by using a 20 point Gaussian quadrature formula for all integrations in \( \phi \). All integrations over the functions \( \{ f_j(t) \} \) in \( t = a(n-\theta) \) were done by sampling each \( f_j(t) \) four times. The \( \{ f_j(t) \} \) consisted of 14 overlapping triangle (divided by the cylindrical coordinate radius) functions equally spaced in \( \theta \). More precisely,

\[
\text{NP} = 31
\]
\[
\text{NPHI} = 20
\]
\[
\text{MT} = 2
\]

where the above variables are input data for the computer program described and listed in Part Two, Section V.

Figures 10 and 11 show the scattering patterns radiated by the currents of Figs. 8 and 9 respectively. The symbols \( \times \) and \( + \) denote \( \sigma_{\theta\theta}/\pi \) and \( \sigma_{\phi\phi}/\pi \), respectively. The solid curves are the exact patterns obtained from the Mie series solution \([11]\). The patterns \( \sigma_{\theta\theta} \) and \( \sigma_{\phi\phi} \) are given by (69), (62) and (63). Here \( \sigma_{\theta\theta} \) is the \( \theta \) polarized pattern versus \( \theta_r \) in the \( \phi = 0 \) plane and \( \sigma_{\phi\phi} \) is the \( \phi \) polarized pattern versus \( \theta_r \) in the \( \phi = 90^\circ \) plane. The THETA in Figs. 10 and 11 refers to \( \theta_r \). For axial incidence, only the \( n=1 \) terms are present in (62) and (63). Elsewhere \([3, 12]\), the pattern \( \sigma_{\theta\theta} \) is called the horizontal polarization because it is polarized parallel to the scattering plane. Similarly, the pattern \( \sigma_{\phi\phi} \) is called the vertical polarization because it is polarized perpendicular to the scattering plane.

Figures 12-17 show the scattering patterns for three other dielectric spheres. Figures 12 and 13 are for relative dielectric constant \( \varepsilon_r = 1.1 \), Figs. 14 and 15 for \( \varepsilon_r = 10.0 \), and Figs. 16 and 17 for \( \varepsilon_r = 20.0 \). All other

Fig. 10. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $\epsilon = 4$, FMCHW solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 11. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $\varepsilon = 4$, Müller solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 12. Plane wave scattering patterns for dielectric sphere, \( k \alpha = 3, \epsilon = 1.1 \), PMCHW solution. Symbols \( \times \) and + denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 13. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $\varepsilon_r = 1.1$, Müller solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 14. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $\epsilon_r = 10$, PMCHW solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 15. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $\varepsilon_r = 10$, Müller solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 16. Plane wave scattering patterns for dielectric sphere, $ka = 3$, $r = 20$, PMCHW solution. Symbols $-$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 17. Plane wave scattering patterns for dielectric sphere, \( ka = 3, r = 20 \), Müller solution. Symbols - and + denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. 18. Plane wave scattering patterns for finite dielectric cylinder of radius $a$ and height $2a$. $a = 0.25$ free space wavelengths, $\varepsilon = 4$, PMCHW solution. Symbols $<$ and $+$ denote horizontal polarization and vertical polarization respectively.
Fig. 19. Plane wave scattering patterns for finite dielectric cylinder of radius $a$ and height $2a$. $a = 0.25$ free space wavelengths, $c_r = 4$, Müller solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively.
parameters in Figs. 12-17 are the same as in Figs. 10 and 11. In Figs. 12 and 13, values less than 0.0001 are plotted at 0.0001.

Figures 18 and 19 show the computed scattering patterns of a finite dielectric cylinder of radius \( a \) and height \( 2a \) when \( a \) is 0.25 free space wavelengths. The relative dielectric constant of the cylinder is \( \epsilon_r = 4 \). The incident field is a plane wave traveling in the positive z direction, the same field which was incident upon the previous dielectric spheres. 

Figure 18 shows \( \frac{\sigma_{\theta}}{\pi a^2} \) and \( \frac{\sigma_\phi}{\pi a^2} \) as obtained from our solution of the PMCHW formulation. Figure 19 shows \( \frac{\sigma_{\theta}}{\pi a^2} \) and \( \frac{\sigma_\phi}{\pi a^2} \) as obtained from our solution of the Müller formulation. The patterns \( \frac{\sigma_{\theta}}{\pi a^2} \) and \( \frac{\sigma_\phi}{\pi a^2} \) are plotted with the symbols \( \times \) and \( + \) respectively.

The equivalent currents which radiate the patterns of Figs. 18 and 19 were obtained by using a 48 point Gaussian quadrature formula for all integrations in \( \phi \). All integrations in \( t \) over the functions \( \{f_j(t)\} \) were done by sampling each \( f_j(t) \) four times. The \( \{f_j(t)\} \) consisted of 11 overlapping triangle (divided by the cylindrical coordinate radius) functions equally spaced in \( t \). More precisely,

\[
\begin{align*}
NP &= 25 \\
NPHI &= 48 \\
MT &= 2
\end{align*}
\]  

where the above variables are input data for the computer program described and listed in Part Two, Section V.

VI. DISCUSSION

According to Figs. 12 and 13, the scattering patterns obtained from our solution of the Müller formulation are more accurate than those obtained from our solution of the PMCHW formulation for the dielectric sphere with \( ka = 3 \) and \( \epsilon_r = 1.1 \). From plots not included in this report, we observed that both our PMCHW solution and our Müller solution for the equivalent
currents on the dielectric sphere were reasonably accurate. However, the following argument shows that when $\varepsilon_r$ is near one, a slight inaccuracy in the equivalent currents could affect the scattering patterns drastically. As $\varepsilon_r$ approaches one, the equivalent electric and magnetic currents approach $n \times H^i$ and $E^i \times n$ respectively whereas the scattering patterns approach zero. This means that the equivalent electric and magnetic currents produce fields which nearly cancel each other. Hence, a slight inaccuracy in the equivalent currents could cause a large percentage inaccuracy in the scattering patterns.

We believe that our Müller solution is more accurate than our PMCHW solution whenever $\varepsilon_r$ is close to one. When $\alpha$ and $\beta$ are given by (22) and (23) as in the Müller formulation, the left-hand sides of (9) and (10) approach $-\mathbf{M}$ and $\mathbf{J}$ respectively as $\varepsilon_r$ approaches one. In this case, the expected solution

$$\mathbf{J} = n \times H^i \quad \mathbf{M} = E^i \times n$$

can be obtained by inspection of (9) and (10). However, if $\alpha = \beta = 1$ as in the PMCHW formulation, the solution to (9) and (10) is not obvious when $\varepsilon_r = 1$ because the field operators on the left-hand sides of (9) and (10) are not diagonal. With our Müller solution, the matrix on the left-hand side of (39) would become tridiagonal for $\varepsilon_r = 1$ if its first two rows of submatrices were interchanged. With our PMCHW solution, no such simplification of this matrix is possible for $\varepsilon_r = 1$.

We recommend at least 10 expansion functions per wavelength per component along the generating curve of the dielectric body of revolution. For example, if the generating curve were one wavelength long, the order of the square matrix on the left-hand side of (39) should be at least 36. The number 36 is arrived at as follows. There should be at least 9 expansion functions per component of current. We say 9 expansion functions rather than 10 because we are using overlapping triangle functions with no peak of triangle function at either ends of the generating curve. There are two components of electric current and two components of magnetic current.
According to equations (20)-(23) of [9] and (58)-(61) of [9], each element of the square matrix in (39) is a triple integral consisting of one integration with respect to $\phi$ and two integrations with respect to $t$. The $\phi$ integral is evaluated by using a Gaussian quadrature formula. Each $t$ integration is done by crude sampling akin to the trapezoid rule. In any case, there should be at least 10 sample points per wavelength in the media in question. For instance, if $\rho_{\text{max}}$ is the largest cylindrical coordinate radius of the dielectric body of revolution and $t_{\text{max}}$ is the length of the generating curve, then

$$NP \geq \frac{10t_{\text{max}}}{\lambda} + 1$$

$$NPHI \geq \frac{10\pi\rho_{\text{max}}}{\lambda}$$

where $NP$ and $NPHI$ are input arguments of the subroutine YZ described and listed in Part Two, Section II and $\lambda$ is the wavelength in the media in question. If $f = e$ in (40)-(41) then $\lambda$ is the wavelength in the external media, but if $f = d$ in (40)-(41) then $\lambda$ is the wavelength inside the diffracting body of revolution. The main program in Part Two, Section V is oversimplified in that it uses the same values of $NP$ and $NPHI$ for both $f = e$ and $f = d$.

Loss of accuracy in the computed patterns of Figs. 16 and 17 may be due to the fact that (74) was violated. According to (74), the values of $NP$ and $NPHI$ for $f = d$ should be nearly 70 or greater instead of the low values appearing in (72). Unfortunately, increasing the values of the variables $NP$ and $NPHI$ increases the computer time required to solve the problem.

We have been trying to obtain accurate numerical results for the dielectric sphere for which $a = 0.2$ free space wavelengths and $\varepsilon_r = 80$. from our general dielectric body of revolution program. We have not been able to obtain clear-cut convergence with respect to the variables on the left-hand sides of (72) because we could not afford to increase them as much as desired. Our PMCHW solution and our Müller solution for the equivalent currents and scattering patterns differ from each other and from the exact solution.
For the sphere, each element of the square matrix or the left-hand side of (39) can be written as a sum over the infinite set of spherical modes. So far, we have not been able to successfully implement this alternate evaluation of the matrix elements in terms of spherical modes. The major difficulty seems to be lack of agreement of a few matrix elements for which both expansion and testing functions are near one of the poles of the sphere.

Both the PMCIIW solution and the Müller solution are obtained by taking a linear combination of (5) and (7) and a linear combination of (6) and (8). There are two other possibilities which are

(1) A linear combination of (5) and (6) and a linear combination of (7) and (8).

(2) A linear combination of (5) and (8) and a linear combination of (6) and (7).

These other two possibilities give rise to alternative numerical solutions which may compare favorably with the PMCIIW solution and the Müller solution.
APPENDIX A

THE EQUIVALENCE PRINCIPLE

Let \( E_e, \ H_e \) be an electromagnetic field defined outside a closed surface \( S \). The permeability, permittivity, and electric and magnetic source currents outside \( S \) are \( \mu_e, \ \epsilon_e, \ J_e, \) and \( M_e \) respectively. This outside situation where the subscript \( e \) stands for "exterior medium" is shown in Fig. A-1. In Fig. A-1, the media and sources inside \( S \) are undisclosed. Let \( E_d, \ H_d \) be an electromagnetic field defined inside \( S \) where the permeability, permittivity, and electric and magnetic sources are \( \mu_d, \ \epsilon_d, \ J_d, \) and \( M_d \) respectively. This inside situation where the subscript \( d \) stands for "diffracting medium" is shown in Fig. A-2. In Fig. A-2, the media and sources outside \( S \) are undisclosed. The equivalence principle states that the solution to the composite radiation problem consisting of medium \( \mu_e, \ \epsilon_e \) and sources \( J_e, M_e \) outside \( S \), medium \( \mu_d, \ \epsilon_d \) and sources \( J_d, M_d \) inside \( S \), and electric and magnetic surface currents \( J, M \) on \( S \) given by

\[
\begin{align*}
J &= n \times (H_e - H_d) \\
M &= (E_e - E_d) \times n
\end{align*}
\]

where \( n \) is the exterior unit normal vector on \( S \) is the composite electromagnetic field \( E, \ H \) defined by

\[
\begin{align*}
E, \ H &= E_e, H_e \quad \text{outside } S \\
E, \ H &= E_d, H_d \quad \text{inside } S
\end{align*}
\]

The composite radiation problem is shown in Fig. A-3 which is entitled composite situation.

The equivalence principle is proved by showing that the configuration of media and sources in Fig. A-3 gives rise to the composite field \( E, H \) defined by (A-3) and (A-4). Now, \( E, H \) will be the field generated by the media and sources of Fig. A-3 if \( E, H \) satisfies Maxwell's equations with source terms included and the radiation condition at infinity. \( E, H \) satisfies Maxwell's equations outside \( S \) and inside \( S \) because \( E_e, H_e \) and
Fig. A-1. Outside Situation.

Fig. A-2. Inside Situation.

Fig. A-3. Composite Situation.

\[ j = n \times (H_e - H_d) \]
\[ m = (E_e - E_d) \times n \]
$E_d$, $H_d$ satisfy Maxwell's equations outside $S$ and inside $S$ respectively. $E$, $H$ also satisfies the radiation condition at infinity because $E_e$, $H_e$ satisfies the radiation condition at infinity. It remains to be shown that Maxwell's equations for $E$, $H$ exhibit the surface current sources $J$ and $M$ given by (A-1) and (A-2).

It is well known that a surface current source on $S$ gives rise to a discontinuity across $S$ of the tangential component of the field. The preceding statement is easily verified by means of an argument based on the integral forms of Maxwell's equations. Now, this same argument can be construed to imply that a discontinuity across $S$ of the tangential component of the field gives rise to a surface current source on $S$. Hence Maxwell's equations for $E$, $H$ exhibit the electric and magnetic surface current sources $J$ and $M$ on $S$ given by (A-1) and (A-2). Therefore, $E$, $H$ is the solution to the composite radiation problem shown in Fig. A-3 because $E$, $H$ satisfies Maxwell's equations with sources and the radiation condition at infinity.
APPENDIX B

PROOF THAT THE SOLUTION TO EQUATIONS (5)-(8) IS UNIQUE

The solution $J, M$ to (5)-(8) will be unique if the associated set of homogeneous equations

\begin{align*}
-n \times E_e^- &= 0 \quad \text{(B-1)} \\
-n \times H_e^- &= 0 \quad \text{(B-2)} \\
-n \times E_d^+ &= 0 \quad \text{(B-3)} \\
-n \times H_d^+ &= 0 \quad \text{(B-4)}
\end{align*}

has only the trivial solution $J = M = 0$.

From (B-1) and (B-2), $J, M$ radiate in $\nu_e, \varepsilon_e$ to produce a field whose tangential components are zero just inside $S$. Hence, according to the relation between $J, M$ and the discontinuity of tangential field across $S$ as exemplified by (A-1) and (A-2), the field $E_e, H_e$ radiated by $J, M$ in $\nu_e, \varepsilon_e$ outside $S$ satisfies

\begin{align*}
-n \times H_e = J 
\end{align*}

just outside $S$. See Fig. B-1.

From (B-3) and (B-4) the electric and magnetic currents $-J, -M$ radiate in $\nu_d, \varepsilon_d$ to produce a field whose tangential components are zero just outside $S$. Hence, according to the relation between $-J, -M$ and the discontinuity of tangential field across $S$, the field $E_d, H_d$ radiated by $-J, -M$ in $\nu_d, \varepsilon_d$ satisfies

\begin{align*}
-n \times (-H_d) &= J 
\end{align*}

just inside $S$. See Fig. B-2.

The equivalence principle is used to combine the outside situation in Fig. B-1 with the inside situation in Fig. B-2 to obtain the composite situation shown in Fig. B-3. Because of (B-5)-(B-8), the composite
Fig. B-1. Radiation of $J, M$ according to (B-1) and (B-2).

Fig. B-2. Radiation of $-J, -M$ according to (B-3) and (B-4).

Fig. B-3. Composite Situation.
The situation in Fig. B-3 is source-free. Therefore, the field in Fig. B-3 is zero everywhere. Hence, the fields in Figs. B-1 and B-2 are zero everywhere which implies that \( J = M = 0 \).

Thus, the solution to (5)-(8) is unique because the associated set (B-1)-(B-4) of homogeneous equations has only the trivial solution.
APPENDIX C

MORE EXAMPLES

The equivalent currents and scattering patterns for the dielectric spheres for which \( \varepsilon_r = 4 \) and \( ka = 4, 5, \) and \( 6 \) are plotted in Appendix C. Figures C-1 to C-4 are for \( ka = 4 \), Figs. C-5 to C-8 are for \( ka = 5 \), and Figs. C-9 to C-12 are for \( ka = 6 \). All other parameters are the same as in Figs. 8-11 in Section V. In particular, the input data for the computer program which generated the method of moments results plotted in Figs. C-1 to C-12 is given by (72).

It is evident from Figs. 8-11 and Figs. C-1 to C-12 that the method of moments solutions for the equivalent currents and scattering patterns are not as accurate at \( ka = 4, 5, \) and \( 6 \) as at \( ka = 3 \). Loss of accuracy at the higher values of \( ka \) may be due to violation of (73).
Fig. C-1. Equivalent electric and magnetic currents for dielectric sphere, $ka = 4$, $e_r = 4$. PMCHW solution. Symbols $\times$ and $+$ denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-2. Equivalent electric and magnetic currents for dielectric sphere, $ka = 4$, $\varepsilon_r = 4$, Müller solution. Symbols $\times$ and $+$ denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-3. Plane wave scattering patterns for dielectric sphere, $ka = 4$, $e = 4$, PMCHW solution. Symbols $\circ$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. C-4. Plane wave scattering patterns for dielectric sphere, $ka = 4$, $\varepsilon_r = 4$, Müller solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. C-5. Equivalent electric and magnetic currents for dielectric sphere, \( ka = 5, \varepsilon_r = 4 \), PMCHW solution. Symbols × and + denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-6. Equivalent electric and magnetic currents for dielectric sphere, $ka = 5$, $c_r = 4$, Müller solution. Symbols × and + denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-7. Plane wave scattering patterns for dielectric sphere, $ka = 1, r_p = 4$, PMLHW solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. C-8. Plane wave scattering patterns for dielectric sphere, $ka = 5$, $c = 4$, Müller solution. Symbols - and + denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. C-9. Equivalent electric and magnetic currents for dielectric sphere, $ka = 6$, $\varepsilon_r = 4$, PMCHW solution. Symbols $\times$ and $+$ denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-10. Equivalent electric and magnetic currents for dielectric sphere, $ka = 6$, $\varepsilon_r = 4$, Müller solution. Symbols $\times$ and $+$ denote magnitude and phase respectively. Solid line denotes exact solution.
Fig. C-11. Plane wave scattering patterns for dielectric sphere, $ka = 6$, $r_a = 4$, PMCHW solution. Symbols $\times$ and $+$ denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
Fig. C-12. Plane wave scattering patterns for dielectric sphere, \(ka = 6, \epsilon = 4\), Müller solution. Symbols \(\times\) and + denote horizontal polarization and vertical polarization respectively. Solid line denotes exact solution.
PART TWO
COMPUTER PROGRAM

I. INTRODUCTION

The computer program calculates the equivalent electric and magnetic currents (58) and (59) and the scattering patterns (70) for a loss-free homogeneous dielectric and/or magnetic material body of revolution immersed in an axially incident plane wave. This computer program consists of a main program and the subroutines YZ, PLANE, DECOMP and SOLVE.

Part Two consists of definitions of the input and output for the subroutines YZ, PLANE, DECOMP, and SOLVE, listings of these subroutines, definitions of the input and output for the main program, a verbal flow chart of the main program, and a listing of the main program with sample input and output. The subroutines YZ and PLANE are similar to subroutines of the same name in [10]. The subroutines DECOMP and SOLVE are, except for dimension statements, exactly the same as in [13]. Hence, the insides of the subroutines YZ, PLANE, DECOMP, and SOLVE are not described in detail in Part Two. Because these subroutines are quite complicated, a black box approach is suggested wherein the user is concerned with just the input and output of these subroutines. However, the user is encouraged to delve inside the main program and to make any changes therein that he deems necessary to suit his needs.

II. THE SUBROUTINE YZ

Description:

The subroutine YZ(NN, NP, NPHI, M, MT, RH, ZH, X, A, Y, Z) stores the matrices \( Y_{nf} \) and \( Z_{nf} \) defined by

\[
Y_{nf} = \begin{bmatrix}
Y_{tt}^{nf} & Y_{t\phi}^{nf} \\
\phi_{tt}^{nf} & \phi_{t\phi}^{nf}
\end{bmatrix}
\]  

(75)

\[
Z_{nf} = \begin{bmatrix}
Z_{tt}^{nf} & Z_{t\phi}^{nf} \\
\phi_{tt}^{nf} & \phi_{t\phi}^{nf}
\end{bmatrix}
\]  

(76')
by columns in Y and Z respectively. The submatrices on the right-hand sides of (75) and (76) are given by (40) and (41). The first 9 arguments of YZ are input variables. Except for the new input variables M and MT, the subroutine YZ is the same as the old subroutine YZ on pages 17-21 of [10]. If M = -1 and MT = 2, these subroutines are exactly the same as far as the calculation of Y and Z in terms of the rest of the input variables is concerned.

\[ M = -1 \] for field evaluation just inside S and \( M = +1 \) for field evaluation just outside S. \( M = -1 \) if \( f = e \) in (75) and \( M = +1 \) if \( f = d \) in (75). The value of \( M \) is not used in calculating (76) because the tangential components of the electric field operator \( E_f \) in (41) are continuous across S. All numerical integrations over \( t \) of \( f_j(t) \) appearing in (35) and (36) are done by sampling each \( f_j(t) \) \( 2*MT \) times. The representations of \( \rho f_i(t) \) and \( \frac{d}{dt} (\rho f_i(t)) \) given by (66) and (67) of [9] are replaced by representations which contain \( 2*MT \) impulse functions instead of 4 impulse functions. For instance, (66) of [9] is replaced by

\[ \rho f_i(t) = \frac{1}{k} \sum_{p=1}^{2*MT} T_{p+(i-1)*2*MT} \delta(t-t_p+(i-1)*MT), \]

(77)

\( i=1,2,\ldots,N \) where \( N \) will be defined in the paragraph which follows the next paragraph. The T's appearing in (77) will be defined by (78).

Seven of the input variables are the same as in the old subroutine YZ on pages 17-21 of [10]. These variables are defined in terms of variables appearing in [9] by

\[ NN = n, \quad NP = P, \quad \text{page 9} \]
\[ NPHI = N_k, \quad \text{page 13} \]
\[ RH(i) = k\phi_i, \quad \text{page 9} \]
\[ ZH(i) = kZ_i, \quad \text{page 9} \]
\[ X(k) = x_k, \quad \text{page 13} \]
\[ A(k) = A_k, \quad \text{page 13} \]
In summary, \( n \) denotes the \( e^{-in\phi} \) dependence in (35) and (36). \((\rho_1^-, z_1^-)\), \(i=1,2,\ldots, P\) are coordinates on the generating curve, the \( k \) which multiplies \( \rho_1^- \) and \( z_1^- \) is the propagation constant, and \( \chi_k \) and \( A_k \) are the abscissas and weights for the \( N \) point Gaussian quadrature integration in \( \phi \). Note that \( NP-1 \) should be an integer multiple of \( MT \). If \( NP-1 \) is not an integer multiple of \( MT \), the program will ignore \( RH(m) \) and \( ZH(m) \) for

\[ [(NP-1)/MT]*MT+1 < m \leq NP \]

where \([NP-1)/MT]\) is the largest integer which does not exceed \((NP-1)/MT\).

Minimum allocations are given by

```
COMPLEX Y(4*N*N), Z(4*N*N)
DIMENSION RH(NP), ZH(NP), X(NPHI), A(NPHI), D(NG)
PD(NG), TF(2*MT*N), CR(NPHI), C1(NPHI), C2(NPHI),
C3(NPHI), C4(NPHI)
COMMON RS(NG), ZS(NG), SV(NG), CV(NG), T(2*MT*N)
```

where

\[ N = [(NP-1)/MT] - 1 \]

\[ NG = (N+1)\ast MT \]

The variables in common make the results of some intermediate calculations done in YZ available to the subroutine PLANE described in Section III of Part Two.

We mention a few portions of YZ which differ from the subroutine listed on pages 18-21 of [10]. Equation (29) of [9] has been generalized to

\[
T_{2*MT*(J-1)+1} = \frac{k^2}{\Lambda_1} \sum_{q=1}^{L} d_{MT*(J-1)+q} - \frac{1}{2} \sum_{d_{MT*(J-1)+1}} \frac{1}{2} d_{MT*(J-1)+1}
\]  

\[
T_{2*MT*(J-1)+1} = \frac{k}{\Lambda_2} \left( \sum_{q=1}^{L} d_{MT*(J-1)+q} + \sum_{d_{MT*(J-1)+1}} \frac{1}{2} d_{MT*(J-1)+1} \right)
\]

for \( J \) = 1, 2, \ldots, \( N \) and \( 1 \) = 1, 2, \ldots, \( MT \) where

\[
\Lambda_1 = \sum_{q=1}^{MT} d_{MT*(J-1)+1}
\]

\[
\Lambda_2 = k \sum_{j=1}^{MT} d_{MT*(J-1)+1}
\]
and the d's are given by equation (28) of [9]. Note that $\Delta_1$ is the electrical length of generating curve over which the first half of $f_j(t)$ exists and that $\Delta_2$ is the electrical length of generating curve over which the second half of $f_j(t)$ exists. The generalization of equation (68) of [9] is

$$T'_{2*MT*(J-1)+I} = \frac{kd_{MT*J-1+I}}{\Delta_1}$$ (82)

$$T'_{2*MT*J-MT+I} = -\frac{kd_{MT*J+I}}{\Delta_2}$$ (83)

for $J = 1, 2, \ldots N$, and $I = 1, 2, \ldots MT$ where $\Delta_1$ and $\Delta_2$ are given by (80) and (81). Expressions (78) - (83) are calculated in DO loop 68. DO loop 12 accumulates $\Delta_1$ in DEL. DO loop 19 puts (78) in $T(2*MT*(J-1)+I)$ and (82) in $TP(2*MT*(J-1)+I)$. DO loop 15 accumulates $\Delta_2$ in DEL. DO loop 16 puts (79) in $T(2*MT*J-MT+I)$ and (83) in $TP(2*MT*J-MT+I)$.

The subscripts $KT$, $LT$, and $J1$ inside DO loop 32 are obtained as follows. Since the generating curve consists of $NG = (N+1)*MT$ small intervals, it is composed of $(N+1)$ large intervals where the $m$th large interval consists of the $((m-1)*MT+1)$th through the $(m*MT)$th small intervals. The index $I$ of DO loop 60 denotes the $I$th small interval. The $I$th small interval is contained in the $(I9+1)$th large interval where $I9 = [(I-1)/MT]$. The second half of $f_{I9}(t)$ and the first half of $f_{I9+1}(t)$ are in this large interval. The index $K$ of DO loop 32 denotes $f_{I9+K-1}(t)$. Since $T((m-1)*2*MT+1)$ through $T(m*2*MT)$ is allotted to $f_{m}(t)$, $m=1, 2, \ldots K$, $f_{m}(t)$ is preceded by $(m-1)$ overlaps. For each overlap the subscript of $T$ increases by an amount MT not accounted for by $I$. Hence, replacing $m$ by $I9+K-1$, we arrive at the subscript

$$KT = I + (I9+K-2)*MT$$

for $T$. Here, $KT$ is the field subscript which refers to the testing function. By analogy, the source subscript $LT$ which refers to the expansion function is given by

$$LT = J + (J9+1-2)*MT$$
Retaining $f_{I9+K-1}(t)$ as the testing function, we take the analogously subscripted function $f_{J9+L-1}(t)$ to be the expansion function and arrive at the matrix subscript

$$J1 = (J9+L-2)*N2 + I9+K-1$$

where, as in the program, $N2 = 2*N$.

DO loop 17 accumulates in $R1$ the contribution to $(Y^tt)^{J,J}$ of equation (31) of [9] due to the $\pi / k^2d_{i}^{1}0_{i}$ term in equation (32) of [9].

This contribution is given by

$$R1 = \sum_{I=1}^{2*MT} I+2*MT*(J-1) * PD(I+MT*(J-1))$$

where, as in the program,

$$PD(i) = -M * \pi / k^2d_{i}^{1}0_{i}$$

Here, the factor $-M$ not included in equation (32) of [9] provides for the choice of field evaluation either outside or inside $S$.

DO loop 18 accumulates in $R1$ the contribution to $(Y^tt)_{J-1,J}$ of equation (31) of [9] due to the $\pi / d_{i}^{1}0_{i}$ term in equation (32) of [9].

This contribution is given by

$$R1 = \sum_{I=1}^{MT} I+2*MT*(J-1) * MT * PD(I+MT*(J-1))$$

LISTING OF THE SUBROUTINE YZ

SUBROUTINE YZ(NN,NP,NI,HI,M,MT,RH,ZH,X,A,Y,Z)

COMPLEX UY(784),Z(784),G1,G2,G3,G4,G5,G6,Y1,Y2,Y3,Y4,Z1,Z2,Z3,Z4
DIMENSION RH(161),ZH(161),X(48),A(48),D(160),PD(160),TP(320)
DIMENSION CR(48),CI(48),C2(48),C3(48),C4(48)
COMMON RS(160),ZS(160),SV(160),CV(160),T(320)
PI=3.141593
PIM=-M*PI
N=(NP-1)/MT-1
N2=2*N
NG=(N+1)*MT
NGM=NG-MT
MT2=MT*2
DC .57 I=1,NG
I2=I+1
DZ=ZH(I2)-ZH(I)
DZ=DZ+DZ*DS
DI=SRT(DR*DR+DZ*DZ)
RS(I)=0.5*(RH(I2)+RH(I))
ZS(I)=0.5*(ZH(I2)+ZH(I))
SV(I)=DR/D(I)
CV(I)=DZ/D(I)
PD(I)=PIM/(D(I)*RS(I))

57 CONTINUE
J1=0
J5=0
DO 68 J=1,N
DEL=0.
DO 12 I=1,MT
J1=J1+1
DEL=DEL+D(J1)
12 CONTINUE
J1=J1-MT
SN=0.
DO 19 I=1,MT
J5=J5+1
J1=J1+1
SN=SN+D(J1)
TP(J5)=D(J1)/DEL
T(J5)=(SN+2*D(J1))*TP(J5)
19 CONTINUE
DEL=0.
DO 15 I=1,MT
J1=J1+1
DEL=DEL+D(J1)
15 CONTINUE
J1=J1-MT
SN=DEL
DO 16 I=1,MT
J5=J5+1
J1=J1+1
SN=SN+D(J1)
TP(J5)=-D(J1)/DEL
T(J5)=-(SN+2*D(J1))*TP(J5)
16 CONTINUE
J1=J1-MT
68 CONTINUE
P12=A*(B)
PH=NN
DO 25 K=1,NPHI
PH=PH+PH
PHN=PH+PH
SN=SN+2*PH
C1(K)=2*SN*SN
C1=12*A(K)
C5=1*CO3(PH)
C1(K)=2*P(K)*C5
C2(K)=R1*SN(PH)*SN(PH)
C4(K) = CS

25 CONTINUE
N2N = N2*N
N4N = N2N*2
DO 62 J = 1, N4N
Y(J) = 0.
Z(J) = 0.
62 CONTINUE
U = (0. + 1.)
DO 59 J = 1, NG
FJ = FN/RS(J)
L1 = 1
L2 = 2
IF(J .LE. MT) L1 = 2
IF(J .GT. NGM) L2 = 1
J9 = (J - 1)/MT
JT = J + MT*(J9 - 2)
J5 = (J9 - 2)*N2 - 1
S1 = 1.
DO 60 I = 1, J
J9 = (I - 1)/MT
IT = I + MT*(19 - 2)
J6 = 19 + J5
FJ = FN/RS(I)
R2 = R*RS(J) - RS(I)
ZP = ZS(J) - ZS(I)
R2 = R*RP*ZP*ZP
IF(I .NE. J) GO TO 41
S1 = 0.5
K2 = 0.0625*D(J)*D(J)
41 R3 = RS(1)*RS(J)
G1 = 0.
G2 = 0.
G3 = 0.
G4 = 0.
G5 = 0.
G6 = 0.
DO 61 K = 1, NPHI
P4 = R2 + R3*CR(K)
R5 = SQRTR4
Z1 = S1/4.5*(COS(R5) - U*SIN(R5))
Y1 = Z*I*(1. + U*R5)/R4
G1 = C1(K)*Y1 + G1
G2 = C2(K)*Y1 + G2
G3 = C3(K)*Y1 + G3
G4 = C4(K)*Z1 + G4
G5 = C2(K)*Z1 + G5
G6 = C3(K)*Z1 + G6
61 CONTINUE
G = U*G3
Y1 = (R*CV(J) - ZP*SV(J)) * G2 - RS(1)*CV(J)*G1
Y2 = (RS(J)*SV(J) - RS(1)*SV(J) - ZP*SV(J))*G3
Y3 = ZP*G3
Y4 = (R*CV(J) - ZP*SV(J)) * G2 - RS(J)*CV(J)*G1
Z2 = U*SV(J)*G5 + CV(J)*G4
G1 = -U*G4
Z2 = -SV(J)*G6
G2 = -FI*G4
Z3 = SV(J)*G6
G3 = FI*G4
65
Z4=U*(G5-FI*G3)
K1=1
K2=2
IF(I. LE. MT) K1=2
IF(I. GT. NGM) K2=1
DO 31 L=L1+L2
LT=JT+MT*L
J7=J6+L*N2
DO 32 K=K1,K2
KT=JT+MT*K
TT=T(LT: T(KT)
J1=J7+K
J2=J1+N
J3=J1+N2N
J4=J3+N
Y(J1)=TT*Y1+Y(J1)
Y(J2)=TT*Y2+Y(J2)
Y(J3)=TT*Y3+Y(J3)
Y(J4)=TT*Y4+Y(J4)
Z(J1)=TT*Z1*TP(LT)*TP(KT)*G1+Z(J1)
Z(J2)=TT*Z2*TP(LT)*T(KT)*G2+Z(J2)
Z(J3)=TT*Z3*TP(KT)*T(LT)*G3+Z(J3)
Z(J4)=TT*Z4+Z(J4)
32 CONTINUE
31 CONTINUE
60 CONTINUE
59 CONTINUE
N2P=N2+1
KD1=1
J1=0
J5=0
DO 11 J=1,N
R1=0.
DO 17 I=1,M12
J5=J5+1
R1=R1+T(JJ)*T(J1)*P(J5)
17 CONTINUE
J1=J1-MT2
J5=J5-MT2
KD2=KD1+N
KD3=KD1+N2N
KD4=KD3+N
G1=Y(KD1)-Y(KD4)
Y(KD1)=P1+G1
Y(KD2)=0.
Y(KD3)=0.
Y(KD4)=K1-G1
Z(KD1)=Z(KD1)+Z(KD1)
Z(KD2)=Z(KD2)-Z(KD3)
Z(KD3)=-Z(KD2)
Z(KD4)=Z(KD4)+Z(KD4)
IF(J-1) 26,27,26
27 J1=J1+MT2
J5=J5+M17
GO TO 22
26 KU1=KD1-1
KU2=KD2-1
KU3=KD3-1

66
KU4=KD4-1
KL1=KD1-N2
KL2=KD2-N2
KL3=KD3-N2
KL4=KD4-N2
R1=0.
DO 18 I=1,MT
    J1=J1+1
    J5=J5+1
    R1=R1+T(J1)*T(J1-MT)*D(J5)
18 CONTINUE
    J1=J1+MT
    G1=Y(KU1)-Y(KL4)
    G2=Y(KU4)-Y(KL1)
    Y(KU1)=R1+G1
    Y(KU2)=Y(KU2)-Y(KL2)
    Y(KU3)=Y(KU3)-Y(KL3)
    Y(KU4)=R1+G2
    Y(KL1)=R1-G2
    Y(KL2)=-Y(KL2)
    Y(KL3)=-Y(KL3)
    Y(KL4)=R1-G1
    Z(KU1)=Z(KU1)+Z(KL1)
    Z(KU2)=Z(KU2)-Z(KL3)
    Z(KU3)=Z(KU3)-Z(KL2)
    Z(KU4)=Z(KU4)+Z(KL4)
    Z(KL1)=Z(KU1)
    Z(KL2)=-Z(KU3)
    Z(KL3)=-Z(KU2)
    Z(KL4)=Z(KU4)
22 KD1=KD1+N2P
11 CONTINUE
    IF(N.LT.3) RETURN
    J2=N2
    DO 13 I=3,N
        J2=J2+N2
        J1=I-2
        K1=I
        DO 14 J=1,J1
            KU1=J2+J
            KU2=KU1+N
            KU3=KU1+N2N
            KU4=KU3+N
            KL2=KL1+N
            KL3=KL1+N2N
            KL4=KL3+N
            Y(KL1)=-Y(KU4)
            Y(KL2)=-Y(KU2)
            Y(KL3)=-Y(KU3)
            Y(KL4)=-Y(KU1)
            Z(KL1)=Z(KU1)
            Z(KL2)=-Z(KU3)
            Z(KL3)=-Z(KU2)
            Z(KL4)=Z(KU4)
14 CONTINUE
13 CONTINUE
RETURN
END
III. THE SUBROUTINE PLANE

Description:

The subroutine PLANE(NN, N, MT, NT, THR, R) puts $R^p_q$ of (49) in

\[ R(j + (m-1)N+(L-1)4N) \]

where $j=1,2,\ldots N$, $m=1$ denotes $pq = t\theta$, $m=2$ denotes $pq = t\phi$, $m=3$ denotes $pq = r\phi$, and $L=1,2,\ldots NT$. Here, $L$ denotes the $L$th value of the receiver angle $\theta_r$. The first 5 arguments of PLANE are input variables. Except for the new input variable $MT$, the subroutine PLANE is the same as the old subroutine PLANE on pages 22-26 of [10]. If $MT = 2$, these subroutines are exactly the same as far as the calculation of $R$ in terms of the rest of the input variables is concerned.

The integration of $f_j(t)$ over $t$ inherent in (49) is approximated by sampling $f_j(t)$ $2*MT$ times instead of 4 times. The representation of $\rho f_j(t)$ given by (66) of [9] is replaced by (77). $NN$ is the value of $n$ appearing in $R^p_q$. It is required that $NN \geq 0$ but this requirement causes no real loss of generality because $R^p_q$ is either even or odd in $n$. $N$ is the number of expansion functions lying on the generating curve. Specifically, $N$ is the maximum value of $i$ in (77). THR(L) is the $L$th value of the receiver angle $\theta_r$ where $L=1,2,\ldots NT$. The variables RS, ZS, SV, CV, and T appearing in the common statement early in the subroutine PLANE are input variables calculated by calling the subroutine YZ beforehand. The calculated values of these variables depend only on the second, fifth, sixth, and seventh arguments ($NP, MT, RH, ZH$) of YZ.

Minimum allocations are given by

- \text{COMPLEX R}(4*NT*N)
- \text{DIMENSION THR(NT), BJ(M)}
- \text{COMMON RS(NG), ZS(NG), SV(NG), CV(NG), T(?*MT*N)}

where

\[ NG = (N+1)*M \]

and $M$ is the largest of the values of $M$ calculated by PLANE. The suggested allocation BJ(50) will work if the maximum circumference of the body of revolution is less than 26 wavelengths.
Most of the statements in the subroutine PLANE are the same as or very similar to statements in the old subroutine PLANE listed on pages 25-26 of [10]. The major difference is in the calculation of the subscript \((T + MT K)\) for \(T\) and the subscript \(J1\) for \(R\). Using reasoning similar to that used to obtain the subscripts \(KT\) and \(J1\) in the subroutine YZ, we arrive at

\[
(IT + MT K) = I + (I9 + K-1) \times MT
\]

\[
J1 = I9 + K-1
\]

where

\[
I9 = [(I-1)/MT]
\]

The above \(J1\) is valid only for \(L=1\). If \(L \neq 1\), then \((L-1) \times 4 \times N\) must be added to this \(J1\).

LISTING OF THE SUBROUTINE PLANE

```fortran
SUBROUTINE PLANE(NN, N, MT, NT, THR, R)
COMPLEX R(1064), U1, U2, R1, R2, R3, R4
DIMENSION THR(37), BJ(50)
COMMON RS(160), ZS(160), SV(160), CV(160), T(320)
NG=(N+1)*MT
U=(0., 1.)
U1=3.141593*U*NN
N4=4*N
JR=N4*NT
DO 22 J=1, JR
R(J)=0.
22 CONTINUE
J5=1
DO 12 I=1, N
CS=CS(SIN(THR(I))
SN=2.*SIN(THR(I))
DO 13 I=1, NG
X=2.5*RS(I)*SN
IF(X .LE. 5E-7) GO TO 10
M=2.6*X+13.-2./X
IF(X .LE. 5E-6) M=10.6+A/A;D; J(X)
IF(M .GE. (NN+2)) GO TO 19
18 BJ1=0.
BJ2=0.
BJ3=0.
IF(NN .GE. 1) BJ1=1.
IF(NN .GE. 0) BJ2=1.
GO TO 24
19 BJ1(M)-=0.
JM=4-1
BJ1(JM)+=1.
```

69
DO 14 J=3,M
J=M-1
BJ(JM) = JM/X*BJ(JM+1) - BJ(JM+2)
14 CONTINUE
S=0.
DO 15 J=3,M,2
S=S+BJ(J)
15 CONTINUE
S=BJ(1)+2.*S
BJ2=BJ(NN+1)/S
BJ3=BJ(NN+2)/S
BJ1=-BJ3
IF(NN.GT.0) BJ1=BJ(NN)/S

24 ARG=ZS(I)*CS
U2=U1*(CCS(ARG)+U*SIN(ARG))
R4=(BJ3-BJ1)*U*U2
R2=(BJ3+B1)*U2
R1=-BJ2*C1(I)*SN*U2+CS*SV(I)*R4
R3=SV(I)*R2
R2=-CS*R2
I9=(I-1)/MT
IT=I+MT*(I9-2)
J7=I9+J5
K1=1
K2=2
IF(I9.EQ.0) K1=2
IF(J9.EQ.N) K2=1
DO 20 K=K1,K2
TT=IT+MT*K
J1=J7+K
J2=J1+N
J3=J2+N
J4=J3+N
R(J1)=TT*R1+R(J1)
R(J2)=TT*R2+R(J2)
R(J3)=TT*R3+R(J3)
R(J4)=TT*R4+R(J4)
20 CONTINUE
13 CONTINUE
12 CONTINUE
RETURN
END
IV. THE SUBROUTINES DECOMP AND SOLVE

Description:

The subroutines DECOMP(N,IPS,UL) and SOLVE(N,IPS,UL,B,X) solve a system of N linear equations in N unknowns. These subroutines will be used in Section V to solve the matrix equation (39). The input to DECOMP consists of N and the N by N matrix of coefficients on the left-hand side of the matrix equation stored by columns in UL. The output from DECOMP is IPS and UL. This output is fed into SOLVE. The rest of the input to SOLVE consists of N and the column of coefficients on the right-hand side of the matrix equation stored in B. SOLVE puts the solution to the matrix equation in X.

Minimum allocations are given by

\[
\text{COMPLEX } \text{UL}(N^*N) \\
\text{DIMENSION SCL(N), IPS(N)}
\]

in DECOMP and by

\[
\text{COMPLEX } \text{UL}(N^*N), \text{B}(N), \text{X}(N) \\
\text{DIMENSION IPS(N)}
\]

in SOLVE.

More detail concerning DECOMP and SOLVE is on pages 46-49 of [13]

LISTING OF THE SUBROUTINES DECOMP AND SOLVE

SUBROUTINE DECOMP(N,IPS,UL)
COMPLEX UL(3136),PIVOT,EM
DIMENSION SCL(56),IPS(56)
DO 5 I=1,N
IPS(I)=I
RN=0.
J1=1
DO 2 J=1,N
ULM=ABS(REAL(UL(J1)))+ABS(AIMAG(UL(J1)));
J1=J1+N
IF(RN--ULM) 1,2,2
1 RN=ULM
2 CONTINUE
SCL(1)=1./RN
5 CONTINUE
NM1=N-1
K2=0
DO 17 K=1,NM1
BIG=0.
DO 11 I=K,N
IP=IPS(I)
IPK=IP+K2
SIZE=(ABS(REAL(UL(IPK))) + ABS(AIMAG(UL(IPK)))) * SCL(IP)
IF(SIZE-BIG) 10, 11, 10
10 BIG=SIZE
IPV=1
11 CONTINUE
IF(IPV-K) 14, 15, 14
14 J=IPS(K)
IPS(K)=IPS(IPV)
IPS(IPV)=J
15 KPP=IPS(K)+K2
PIVOT=UL(KPP)
KP1=K+1
DO 16 I=KP1,N
KP=KP
IP=IPS(I)+K2
EM=-UL(IP)/PIVOT
18 UL(IP)=-EM
DO 16 J=KP1,N
IP=IP+N
KP=KP+N
UL(IP)=UL(IP)+EM*UL(KP)
16 CONTINUE
K2=K2+N
17 CONTINUE
RETURN
END
SUBROUTINE SOLVE(N, IPS, UL, B, X)
COMPLEX UL(3136), B(56), X(56), SUM
DIMENSION IPS(56)
NP1=N+1
IP=IPS(1)
X(1)=B(IP)
DO 2 I=2,N
IP=IPS(I)
IPB=IP
IM1=I-1
SUM=0.
DO 1 J=1, IM1
SUM=SUM+UL(IP)*X(J)
1 IP=IP+N
2 X(I)=B(IPB)+SUM
K2=N*(N-1)
IP=IPS(N)+K2
X(N)=X(N)/UL(IP)
DO 4 I=1, N
I=NP1-IRACK
K2=K2-N
IP1=IPS(I)+K2
IP1=I+1
SUM=0.
IP=IP1
DO 3 J=IP1, N
IP=IP+N
3 SUM=SUM+UL(IP)*X(J)
4 X(I)=(X(I)-SUM)/UL(IP)
RETURN
END

V. THE MAIN PROGRAM

Description:

The main program calculates the electric and magnetic currents (58)-(59) and the normalized scattering cross sections $\frac{\sigma^{el}}{\lambda^2}$ and $\frac{\sigma^{m}}{\lambda^2}$ (70) for the $\theta$ polarized axially incident ($\theta_t = 180^\circ$) plane wave (50)-(51).

The main program calls the subroutines YZ, PLANE, DECOMP, and SOLVE. The main program is short and simple. It is a representative application of the theory in Part One of this report.

Input data is read early in the main program according to

```
READ(1,10) NP, NPHI, MT, NT
10 FORMAT(4I3)
READ(1,11) BK, UR, ER, ALP, BET
11 FORMAT(5E14.7)
READ(1,12)(RH(I), ZH(I), I=1, NP)
12 FORMAT(IOF8.4)
READ(1,11)(X(K), K=1, NPHI)
READ(1,11)(A(K), K=1, NPHI)
```

The input variables NP, NPHI, MT, RH, ZH, X, and A are very similar to variables of the same names in the argument list of the subroutine YZ. In summary, $(\rho_i, z_i)$ on the generating curve, $2*MT$ is the number of values of $t$ at which $f_j(t)$ is sampled for the purpose of numerical integration, and X and A are respectively the abscissas and weights for the NPHI point Gaussian quadrature integration in $\phi$.

The scattering cross sections are evaluated at receiver angles $\theta_r = (J-1)\pi/(NT-1)$ radians for $J=1,2,\ldots NT$. BK is the propagation constant $k$ in the external medium. This $k$ appears in (50)-(51). UR and ER are respectively the relative permeability $\mu_d$ and the relative permittivity $\varepsilon_d$ of the body of revolution. Here, $\mu_d$ and $\varepsilon_d$ are the permeability and permittivity of the (diffracting) body of revolution.
and \( \mu_e \) and \( \varepsilon_e \) are those of the external medium. ALP and BET are respectively the constants \( \alpha \) and \( \beta \) appearing in (39). The PMCHW solution is obtained by setting \( \alpha = \beta = 1 \). The Müller solution is obtained if \( \alpha \) and \( \beta \) are given by (22) and (23).

Minimum allocations are given by

- \( \text{COMPLEX YE}(4\times\text{N}\times\text{N}), ZE(4\times\text{N}\times\text{N}), R(4\times\text{NT}\times\text{N}), \)
- \( B(4\times\text{N}), YD(4\times\text{N}\times\text{N}), ZD(4\times\text{N}\times\text{N}), Y(16\times\text{N}\times\text{N}), \)
- \( C(4\times\text{N}) \)

\( \text{DIMENSION RH(NP), ZH(NP), X(NPHI), A(NPHI),} \)
- \( \text{THR(NT), RC(N), IPS(4\times\text{N})} \)

where

\[ N = \left\lfloor \frac{(NP-1)}{MT} \right\rfloor - 1 \]

Statement 38 puts \( Y_{le} \) defined in accordance with (75) by

\[ Y_{le} = \begin{bmatrix} Y_{le}^{\dagger} & Y_{le}^{\phi} \\ Y_{le}^{\phi \dagger} & Y_{le}^{\phi \phi} \end{bmatrix} \quad (84) \]

in YE and \( Z_{le} \) defined in accordance with (76) by

\[ Z_{le} = \begin{bmatrix} Z_{le}^{\dagger} & Z_{le}^{\phi} \\ Z_{le}^{\phi \dagger} & Z_{le}^{\phi \phi} \end{bmatrix} \quad (85) \]

in ZE. Storage of \( Y_{le} \) and \( Z_{le} \) is by columns.

Statement 39 puts the matrix \( R_1 \) defined by

\[ R_1 = \begin{bmatrix} R_1^{\dagger \theta} & R_1^{\phi \theta} \\ R_1^{\phi \dagger \theta} & R_1^{\phi \phi} \end{bmatrix} \]

in R. The column vectors on the right-hand side of (86) appear in (56). For receiver angle \( \theta_r = \text{THR}(J) \), \( R_1 \) of (86) is stored by columns in \( R((J-1)\times4\times\text{N}+1) \) through \( R((J\times4\times\text{N}) \).
DO loop 22 uses (56) to store the right-hand side of (39) in B. DO loop 22 also puts 2/(kp\_I*M*T+1) in RC where \( \rho_{I*M*T+1} \) is the cylindrical coordinate radius evaluated at the peak of the triangle function inherent in \( f_1(t) \).

Statement 40 puts \( Y_{1d} \) defined by

\[
Y_{1d} = \begin{bmatrix}
\dot{Y}_{1d} & Y_{1d}^t \\
Y_{1d}^\phi & \phi_{1d}
\end{bmatrix}
\]  

(87)

in YD and \( Z_{1d} \) defined by

\[
Z_{1d} = \begin{bmatrix}
\dot{Z}_{1d} & Z_{1d}^t \\
Z_{1d}^\phi & \phi_{1d}
\end{bmatrix}
\]  

(88)

in ZD. Storage of \( Y_{1d} \) and \( Z_{1d} \) is by columns.

Nested DO loops 26 and 27 put the first two columns of submatrices on the left-hand side of (39) in Y. The index J of DO loop 26 denotes the Jth column of the composite square matrix on the left-hand side of (39).

Nested DO loops 28 and 29 put the third and fourth columns of submatrices on the left-hand side of (39) in Y. The index J of DO loop 28 denotes the (2*N+J)th column of the composite square matrix on the left-hand side of (39).

Statements 41 and 42 solve the matrix equation (39) for the composite column vector consisting of \( \dot{V}_1^t, \dot{V}_1^\phi, \dot{I}_1^t, \) and \( \dot{I}_1^\phi \). This composite column vector is stored in C.

At the peak of the Jth triangle function, the n=1 term of the equivalent electric current (58) reduces to
\[ J^0 = \frac{2}{\rho_{MT*J+1}} I_{tJ}^0 u_t \cos \phi + \frac{2j}{\rho_{MT*J+1}} I_{\phi J}^0 u_{\phi} \sin \phi \quad (89) \]

DO loop 31 prints the real and imaginary parts of

\[ U_1 = \frac{2}{(k_0 MT*J+1)} I_{tJ}^0 , \]

the real and imaginary parts of

\[ U_2 = \frac{2j}{(k_0 MT*J+1)} I_{\phi J}^0 \]

and the magnitudes of \( U_1 \) and \( U_2 \). Here, \( U_1 \) is the \( t \) component of the equivalent electric current in the E plane and \( U_2 \) is the \( \phi \) component of the equivalent electric current in the H plane when the \( y \) component of the incident magnetic field is minus one at the origin.

At the peak of the \( J \)th triangle function, the \( n=1 \) term of the equivalent magnetic current (59) reduces to

\[ B(J) = \frac{1}{\eta} M_\theta = \frac{2j}{(k_0 MT*J+1)} V_{tJ}^\theta u_t \sin \phi + \frac{2}{(k_0 MT*J+1)} V_{\phi J}^\phi u_{\phi} \cos \phi \quad (90) \]

DO loop 34 prints the real and imaginary parts of

\[ B(J) = \frac{2j}{(k_0 MT*J+1)} V_{tJ}^\theta \]

and the real and imaginary parts of

\[ U_2 = \frac{2}{(k_0 MT*J+1)} V_{\phi J}^\phi \]

and the magnitudes of \( B(J) \) and \( U_2 \). Here, \( B(J) \) is the \( t \) component of the equivalent magnetic current in the H plane and \( U_2 \) is the \( \phi \) component of the equivalent magnetic current in the E plane when the \( x \) component of the incident electric field is minus one at the origin.

Nested DO loops 35 and 36 calculate and print \( \frac{\sigma_{\theta \theta}}{\lambda^2} \) and \( \frac{\sigma_{\phi \phi}}{\lambda^2} \) of (70). Inner DO loop 36 accumulates the portion
\[ R_1^t V_1^\theta + R_1^t V_1^\phi + R_1^t I_1^\theta + R_1^t I_1^\phi \]

of (62) in ET and the portion

\[ -R_1^t V_1^\theta - R_1^t V_1^\phi + R_1^t I_1^\theta + R_1^t I_1^\phi \]

of (63) in EP. The W1 and W2 printed in DO loop 35 are respectively \( \sigma_0/\lambda^2 \) and \( \sigma_\theta/\lambda^2 \) for receiver angle \( \theta_r = (J-1)*\pi/(NT-1) \) radians.

Suggested modifications of the main program are:

1) Changing the normalization of the scattering patterns.
   For example, one could replace (70) by (69). All scattering patterns in Part One, Section V are plots of (69).

2) Removing the restriction that the values of the input arguments NP, NPHI, MT, RH, ZH, X, and A of the subroutine YZ be the same in call statements 38 and 40. This modification is indicated by (74).

3) Generalizing from axial plane wave incidence to oblique plane wave incidence.

The above three modifications can be realized without tampering with any of the subroutines YZ, PLANE, DECOMP, and SOLVE.

The sample input and output accompanying the listing of the main program is for the dielectric sphere with \( \omega = 1 \) and \( \varepsilon_r = 4.77 \).
LISTING OF THE MAIN PROGRAM

//P3M JOB (XXXX, XXXX, 1, 2), 'MAUTZ, JOE', REGION=200K
// EXEC WATFIV
//G3.SYSIN DD *
$JOB
MAUTZ, TIME=1, PAGES=40

C
SUBROUTINES YZ, PLANE, DECOMP, AND SOLVE ARE CALLED.
COMPLEX YE(784), ZE(784), R(1064), B(56), YD(784), ZD(784), Y(3136)
COMPLEX C(56), U, U1, U2, ET, EP, CONJG
DIMENSION RH(161), ZH(161), X(48), A(48), THR(37), RC(19), IPS(56)
READ(1, 10) NP, NPHI, MT, NT
10 FORMAT(4I3)
READ1,11) BK, UR, ER, ALP, BET

11 FORMAT(5E14.7)
READ1,12) (RH(I), I=1, NP)
READ(1, 12) (ZH(I), I=1, NPHI)

12 FORMAT(10F8.4)
READ1,11) (X(K), K=1, NPHI)
READ1,11) (A(K), K=1, NPHI)
WRITE(3, 13) NP, NPHI, MT, NT

13 FORMAT1 NP, NPHI, MT, NT/(1X, I3, I5, 2I3))
WRITE(3, 14) BK, UR, ER, ALP, BET

WRITE3, 15) (RH(I), I=1, NP)

15 FORMAT1 RH'/ (1X, 10F8.4))
WRITE3, 16) (ZH(I), I=1, NP)

16 FORMAT1 ZH'/ (1X, 10F8.4))
WRITE3, 17) (X(K), K=1, NPHI)

17 FORMAT1 X'/ (1X, 5E14.7))
WRITE3, 18) (A(K), K=1, NPHI)

18 FORMAT1 A'/ (1X, 5E14.7))
DO 19 J=1, NP
RH(J)= BK*RH(J)
ZH(J)= BK*ZH(J)

19 CONTINUE

38 CALL YZ(1, NP, NPHI, -1, MT, RH, ZH, X, A, YE, ZE)
WRITE3, 20) (YE(I), I=1, 4), (ZE(I), I=1, 4)

20 FORMAT1 SOME ELEMENTS OF YE AND ZE'/ (1X, 4E14.7))
PI=3.141593
DT=PI/(NT-1)
DO 21 J=1, NT
THF(J)=(J-1)*DT

21 CONTINUE

N=(NP-1)/MT-1

39 CALL PLANE(1, NP, NPHI, MT, NT, THR, R)
N2=2*N
N3=3*N
N4=4*N
NTN=(NT-1)*N4
DO 22 I=1, N
J3=I+NTN
B(I)=R(J3)
B(I+N)=R(J3+N)
B(I+N2)=-R(J3+N3)
B(I+N3)=-R(J3+N2)
RC(I)=2./RH(MT*I+1)

22 CONTINUE
WRITE3, 23) (B(I), I=1, N2)

23 FORMAT1 HALF OF THE ELEMENTS OF B'/ (1X, 4E14.7))
EM=SQR(T(UR=ER)
DO 24 I=1,NP
  H(I)=EM*RH(I)
  ZH(I)=EM*ZH(I)

24 CONTINUE

DO 40 J=1,N
  CALL YZ(1,NP,NFHI,1,MT,RH,ZH,X,A,YD,ZA)
  WRITE(3,25)(YD(I),I=1,4),(ZD(I),I=1,4)

25 FORMAT(* SOME ELEMENTS OF YD AND ZD*/(1X,4E14.7))

40 CONTINUE

CALL DECOMP(N4,IPSY)
CALL SOLVE(N4,IPSY,Y,BC)
WRITE(3,30)

30 FORMAT(* REAL JT REAL JP IMAG JT IMAG JP MAG JT
IMAG JP*)
  U=(0.,1.)
  DO 31 J=1,N
      RK=RC(J)
      ET=U*RR
      B(J)=ET*C(J)
      J1=J+N
      B(J1)=RK*C(J1)
      U1=KR*C(J1+N2)
      U2=ET*C(J1+N3)
      W1=CIABS(U1)
      W2=CIABS(U2)
      WRITE(3,32)U1,U2,W1,W2

32 FORMAT(1X,6E11.4)

31 CONTINUE

WRITE(3,33)

33 FORMAT(* REAL MT IMAG MT REAL MP IMAG MP MAG MT
IMAG MP*)
  DO 34 J=1,N
      U2=B(J+N)
      W1=CIABS(R(J))

34 CONTINUE
W2=CABS(U2)
WRITE(3,32) B(J),U2,W1,W2
34 CONTINUE
WRITE(3,37)
37 FORMAT(0<SIGTHETA SIGPHI*)
CON=.25/PI**3
JR1=0
DO 35 J=1,NT
ET=0.
EP=0.
DO 36 I=1,N2
JC2= I+N2
JR1=JR1+1
JR2=JR1+N2
ET=ET+R(JR2)*C(I)*R(JR1)*C(JC2)
EP=EP-R(JR1)*C(I)+R(JR2)*C(JC2)
36 CONTINUE
JR1=JR1+N2
WI=CON*ET*CONJG(ET)
W2=CCN*EP*CCNJG(EP)
WRITE(3,32) WI,W2
35 CONTINUE
STOP
END

$DATA
21 20 2.19
0.1000000E+01 0.1000000E+01 0.4000000E+01 0.1000000E+01 0.1000000E+01
0.0000 0.1564 0.3090 0.4540 0.5878 0.7071 0.8090 0.8910 0.9511 0.9877
1.0000 0.9877 0.9511 0.8910 0.8090 0.7071 0.5878 0.4540 0.3090 0.1564
0.0000 -1.0000 -0.9877 -0.9511 -0.8910 -0.8090 -0.7071 -0.5878 -0.4540 -0.3090 -0.1564
0.0000 0.1564 0.3090 0.4540 0.5878 0.7071 0.8090 0.8910 0.9511 0.9877
1.0000 -0.9931286E+00 -0.9639719E+00 -0.9122344E+00 -0.8391170E+00 -0.7463319E+00
-0.6360537E+00 -0.5108670E+00 -0.3737061E+00 -0.2277859E+00 -0.1527534E+00
-0.7652652E+00 -0.5108670E+00 -0.3737061E+00 -0.2277859E+00 -0.1527534E+00
-0.1761401E+00 -0.4060143E+00 -0.6267205E+00 -0.832674E+00 -0.1761401E+00
-0.1527534E+00 -0.1420961E+00 -0.1316866E+00 -0.1181945E+00 -0.1019301E+00
-0.1019301E+00 -0.832674E+00 -0.6267205E-01 -0.4060143E-01 -0.1761401E-01
*STOP

PRINTED OUTPUT
NP NPNI MT NT
21 20 2 19

BK
0.1000000E+01 0.1000000E+01 0.4000000E+01 0.1000000E+01 0.1000000E+01
0.0000 0.1564 0.3090 0.4540 0.5878 0.7071 0.8090 0.8910 0.9511 0.9877
1.0000 0.9877 0.9511 0.8910 0.8090 0.7071 0.5878 0.4540 0.3090 0.1564
0.0000

RH
-0.9931286E+00 -0.9639719E+00 -0.9122344E+00 -0.8391170E+00 -0.7463319E+00
-0.6360537E+00 -0.5108670E+00 -0.3737061E+00 -0.2277859E+00 -0.1527534E+00
-0.7652652E+00 -0.5108670E+00 -0.3737061E+00 -0.2277859E+00 -0.1527534E+00
-0.1761401E+00 -0.4060143E+00 -0.6267205E+00 -0.832674E+00 -0.1761401E+00
-0.1527534E+00 -0.1420961E+00 -0.1316866E+00 -0.1181945E+00 -0.1019301E+00
-0.1019301E+00 -0.832674E+00 -0.6267205E-01 -0.4060143E-01 -0.1761401E-01

ZH
-0.9931286E+00 -0.9639719E+00 -0.9122344E+00 -0.8391170E+00 -0.7463319E+00
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REFERENCES


