Renormalization of the Graphene Dispersion Velocity Determined from Scanning Tunneling Spectroscopy

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(Received 12 July 2012; published 11 September 2012)

In graphene, as in most metals, electron-electron interactions renormalize the properties of electrons but leave them behaving like noninteracting quasiparticles. Many measurements probe the renormalized properties of electrons right at the Fermi energy. Uniquely for graphene, the accessibility of the electrons at the surface offers the opportunity to use scanned probe techniques to examine the effect of interactions at energies away from the Fermi energy, over a broad range of densities, and on a local scale. Using scanning tunneling spectroscopy, we show that electron interactions leave the graphene energy dispersion linear as a function of excitation energy for energies within \( \pm 200 \text{ meV} \) of the Fermi energy. However, the measured dispersion velocity depends on density and increases strongly as the density approaches zero near the charge neutrality point, revealing a squeezing of the Dirac cone due to interactions.

DOI: 10.1103/PhysRevLett.109.116802
PACS numbers: 73.22.Pr, 68.37.Ef, 73.20.At

In the absence of interactions the low energy excitations in graphene are described by massless Dirac quasiparticles with a linear energy-momentum dispersion \( E = \hbar v_k \), where the constant of proportionality is the carrier group velocity \( v \). When many-body interactions are included, the energy dispersion can change [1–12]. In general, this change depends on the parameters \( \{n, k, r_s\} \), where \( n \) is the carrier density, \( k \) is the momentum relative to the Brillouin zone corner, and \( r_s \) is the interaction parameter or coupling constant describing the relative strength of the Coulomb interactions and is given by the ratio of potential to kinetic energies [13]. In graphene, both the kinetic and potential energies scale as the square root of the density, so that \( r_s \) is independent of density. Therefore, in graphene, the dependence of the dispersion renormalization on \( r_s \) and \( n \) can be separated. For monolayer graphene \( r_s \) is given by \( r_s = e^2/4\pi\varepsilon_0\varepsilon_m\hbar v \), where \( e \) is the electron charge, \( \hbar \) is Planck’s constant divided by \( 2\pi \), \( \varepsilon_0 \) is the permittivity of free space, and \( \varepsilon_m \) is the effective dielectric constant of the medium in which the graphene sheet is embedded.

Previously, measurements of the graphene dispersion renormalization have either examined solely the density dependence [11] or the dependence on energy [6,9,10] but not both over wide ranges of density and energy due to limitations in experimental techniques. Transport measurements, which are sensitive to the behavior of electrons at the Fermi energy \( (E_F) \), have shown the dispersion velocity at \( E_F \) increases at low density [11]. Angle-resolved photoemission spectroscopy has examined the energy dependence of the dispersion and found deviations from linearity at high energies below approximately \( -0.4 \text{ eV} \) [9]. Angle-resolved photoemission spectroscopy measurements are complicated by the appearance of plasmonic and phonon structures which can distort the spectrum at the Dirac point [6], and they have not been able to examine the density dependence in detail due to the lack of large area gated graphene devices. Scanning tunneling spectroscopy measurements of back-gated devices allows the graphene dispersion to be examined as a function of energy away from the Fermi energy [8,14], so that the renormalization can be determined separately as a function of excitation energy and density.

For excitations at the Fermi level, i.e., at \( k = k_F \), many-body theory predicts that the renormalized spectrum of graphene can be characterized by a velocity \( v^* \) given by [1–3]

\[
\frac{v^*}{v} = 1 - r_s \left[ \frac{5}{3} + \ln(r_s) \right] + \frac{r_s}{8} \ln\left( \frac{n_c}{n} \right). \quad (1)
\]

The Fermi velocity enhancement described in Eq. (1) was calculated within the random phase approximation, where \( v \) is the bare dispersion velocity in the absence of interactions and \( n_c \) is a density corresponding to the ultraviolet cutoff energy, which is \( \approx 3 \text{ eV} \) [2,11]. The 2nd and 3rd terms in Eq. (1) result from intraband and interband contributions, respectively. The expression in Eq. (1) describes the reshaping of the Dirac cone at energies equal to the Fermi energy \( E_F \). Velocity enhancements obtained from Shubnikov–de Haas oscillations in transport measurements

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**Performing Organization:** Columbia University, Department of Electrical Engineering, New York, NY, 10027

**Dates Covered:** 00-00-2012 to 00-00-2012

**Security Classification:** Unclassified

**Distribution/Availability Statement:** Approved for public release; distribution unlimited

**Number of Pages:** 5
have considered only the interband contribution in Eq. (1) and found the need for an additional density-dependent parameter to get agreement with experiment [11]. The question on how the dispersion velocity depends on energy over a broad range of densities remains open.

In this Letter, we present the first experimental measurements of the Landau level (LL) spectroscopy of graphene on hexagonal boron nitride (hBN) spacer layers. The lower disorder in graphene with hBN spacer layers allows a significant improvement in LL lifetimes, with the appearance of many LLs over a wide energy range. This enhanced spectrum, compared to previous results without hBN [8,14], allows us to separate the energy and density contributions to the renormalization of the graphene band structure. We show how to quantitatively examine the Landau level tunneling spectrum taking into account probe tip and graphene quantum capacitances and to extract interaction driven velocity renormalization with good accuracy. We find that interactions do not significantly distort the Dirac cones at low energies. Instead, they preserve the linear dispersion while parametrically renormalizing the dispersion velocity at a given fixed density. The measured renormalized velocity is satisfactorily described by the random phase approximation theory incorporating the interaction strength $r_e$ and the ultraviolet cutoff density $n_C$ [2,12] described by Eq. (1) without the need of additional parameters.

The experiments were performed in a custom designed cryogenic scanning tunneling microscope system operating at 4 K with applied magnetic fields from 0 to 8 T [15]. The graphene devices were fabricated by the method detailed in Dean et al. [16]. A heavily doped silicon substrate was used as a back gate to control the carrier density $n$ in the sample by applying back gate voltage $V_G$. In our experiments, a bias voltage $V_B$ is applied to the graphene, and the tunneling current $I$ is measured from the tip to obtain tunneling spectra $dI/dV$ as a function of $V_B$, $V_G$, and applied magnetic field $B$. Topography measurements of graphene on hBN [Fig. 1(a)] are characterized by a moiré pattern formed by the relative rotation of the graphene sheet with respect to the underlying hBN crystal [17,18]. The topographic modulation in Fig. 1(a) is consistent with the thin BN spacer layer (5 nm) used in these studies [16].

Upon applying a magnetic field, the tunneling spectra develops sharp peaks as a function of $V_B$ as shown in Fig. 1(b), indicating the formation of the LLs [14,19–21]. The reduced disorder in graphene on hBN is evident by the high fidelity of the LL formation compared to previous measurements on SiO$_2$ [8,14]. The LL index $N$ can be identified from the peak distance in bias $V_B$ compared to the expected LL energy spectra proportional to $|N|^{1/2}$. The widths of the LLs, determined by the lifetime of the Dirac quasiparticles, approaches those measured in the low-disorder graphene grown epitaxially on SiC [19,21]. Figure 1(c) shows the map of the $N = -1$ LL peak position, which reflects the spatial variation of disorder potential. Within the area of interest, two local potential extrema are indicated as A and B. At low density, A becomes an electron puddle (blue) while B a hole puddle (yellow).

The tunneling spectra in Fig. 1(b) reflect the graphene LL density of states at a gate voltage of $V_G = -30$ V, which induces additional holes to lower the Fermi level by $\approx 200$ meV with respect to the Dirac point. Repeating the same tunneling spectra measurements while varying the back gate voltage, thus controlling the relative positions of the Fermi level and the Dirac point, we can obtain a complete data set, which we refer to as a “gate map.” The resulting tunneling spectra, represented in a two-dimensional plot of $V_B$ and $V_G$, are shown in Figs. 2(a) and 2(b), for 2 and 5 T, respectively. The LL spectra variation as a function of energy and density can be investigated by tracing peaks of the tunneling spectra (bright curve traces) in these gate maps. At a lower magnetic field [Fig. 2(a)], the smooth curvature of the LLs with gate voltage approximately corresponds to the energy variation of the charge neutrality point (Dirac point), which varies as $E_D \propto n^{1/2}$. In larger fields [Fig. 2(b)], LL formation leads to the familiar stair case pattern in the gate maps, previously seen in graphene on SiO$_2$ [8,14], or GaAs two-dimensional electron gases [22]. Every LL becomes pinned at the Fermi level until it fills with carriers, and upon filling
the next unfilled LL makes a quick transition to become pinned at $E_F$.

To accurately analyze the energy dispersion, we simulate the single particle properties in the gate maps [Fig. 2(c)] by calculating the tunneling conductance spectral map using a capacitor model [23] that includes the graphene-back-gate capacitance ($C_G$), graphene-probe capacitance ($C_P$), the graphene quantum capacitance, and a constant velocity $v$. The simulation in Fig. 2(c) shows that the LL staircase transitions do not occur at fixed $V_G$ but at different values as $V_B$ changes, appearing as a line with an angle in the gate map. This angle is due to the local gating of the graphene by the bias potential between the probe tip and graphene and is determined by the ratio of the back gate and probe capacitances, $C_G/C_P$. The constant density (or constant chemical potential) axis is therefore along the $V_B$ spectra taken at multiple fixed $V_G$. The yellow overlays are portions of the simulation in (c). A comparison [Fig. 2(c)] by calculating the tunneling conductance spectra at $B = 5$ T with a constant density $dI/dV$ magnitude. (a) $dI/dV$ gate map at $B = 2$ T obtained in the hole puddle labeled $B$ in Fig. 1(c). (b) $dI/dV$ gate map spectra at $B = 5$ T. The tilted red line indicates an axis of constant density (chemical potential) following the transitions of the LLs as discussed in the text, while the vertical blue line indicates constant $V_G$. The yellow overlays are portions of the simulation in (c). (c) Simulation of the $dI/dV$ gate map spectra at $B = 5$ T with a constant dispersion velocity of $1.1 \times 10^6$ m/s and the probe-sample capacitance determined from the slope of the LL transitions in (b) [24]. Portions of the simulation are overlaid in (b) (yellow lines) showing agreement at high densities and deviations at low densities from using a constant velocity. The white arrow points to the simulated LL$_2$ position which lies in between the experimental LL$_1$ and LL$_2$ peaks at low density.

FIG. 2 (color online). Gate mapping tunneling spectroscopy of the Landau level density of states of graphene on hBN. Each map is built of individual $dI/dV$ vs $V_B$ spectra taken at multiple fixed $V_G$. The color scale is the $dI/dV$ magnitude. (a) $dI/dV$ gate map at $B = 2$ T obtained in the hole puddle labeled $B$ in Fig. 1(c). (b) $dI/dV$ gate map spectra at $B = 5$ T. The tilted red line indicates an axis of constant density (chemical potential) following the transitions of the LLs as discussed in the text, while the vertical blue line indicates constant $V_G$. The yellow overlays are portions of the simulation in (c). (c) Simulation of the $dI/dV$ gate map spectra at $B = 5$ T with a constant dispersion velocity of $1.1 \times 10^6$ m/s and the probe-sample capacitance determined from the slope of the LL transitions in (b) [24]. Portions of the simulation are overlaid in (b) (yellow lines) showing agreement at high densities and deviations at low densities from using a constant velocity. The white arrow points to the simulated LL$_2$ position which lies in between the experimental LL$_1$ and LL$_2$ peaks at low density.

To accurately determine the interaction driven velocity enhancements, we define a constant density (or chemical potential) axis in the $(V_B, V_G)$ space using the transitions of the LLs, which is determined from the $C_G/C_P$ ratio [24]. Therefore the LL energy $E_N$ can be obtained at a fixed density for the $N$th LL (LL$_N$) by measuring the LL spectral peak positions in the plateau regions along lines of constant density [the tilted red line in Fig. 2(b)] thereby converting $dI/dV(V_B, V_G)$ to $dI/dV(E, n)$. A comparison of the $dI/dV$ spectra at constant $V_G$ vs constant density $n$ [Fig. 3(a)] shows that the spectrum at constant gate voltage overestimates the energy scale by as much as 5% leading to a corresponding error in an estimate of the dispersion velocity. The importance of examining the correct LL$_N$ peak energies at constant density is compounded by the fact that the energy scale error at constant gate voltage...
varies with density, since the \( C_G/C_F \) ratio varies. This will result in errors as large as 30% in determining \( r_s \) parameters from the velocity incorrectly determined at constant gate voltage. Below, we determine the many-body corrections to the renormalized velocity by correctly analyzing the LL spectra at constant density with a correct energy scale [24].

In graphene, the linear dispersion yields the LL spectra:

\[
E_N = E_D + \text{sgn}(N) \sqrt{2e\hbar^2 N B}
\]

The LL spectroscopy presented in Fig. 2 thus allows us to probe the dispersion velocity at fixed density and magnetic field by inspecting \( E_N \) for different \( N \) [Fig. 3(a)]. Figure 3(b) shows \( E_N - E_D \) vs \( \sqrt{NB} \) for different density values. Here \( E_D \) is obtained from \( E_{N=0} \) at a given density. A remarkable feature in Fig. 3(b) is that the LL dispersion is highly linear [see the inset in

![Fig. 3 (color online). Landau level dispersion. (a) \( dI/dV \) spectra (red line) at constant density obtained along the red line in Fig. 2(b). A fit of LL peaks \( N = 1 \) to \( N = 6 \) is shown in green using a series of Lorentzians. (Inset) A comparison of the \( dI/dV \) spectra at constant gate voltage (\( V_G = 12 \) V) (blue line) and constant density (red line), obtained along the lines indicated in Fig. 2(b). The spectra at constant gate voltage overestimates the energy scale by 5% at this density. (b) Determination of the graphene dispersion velocity from the LL peak energy positions (obtained along lines of constant density in the gate maps) for the electron puddle \( A \) at \( B = 2 \) T. The peak positions are plotted as a function of the square root of the Landau orbital index \( N \) and magnetic field \( B \). A linear fit (solid lines) is used to determine the dispersion velocity. (Inset) Residuals from the linear fit showing very good linearity in the LL dispersion.]

![Fig. 4 (color online). Renormalized graphene dispersion. (a) Schematic of the Dirac cone variation as a function of density. The data presented in this Letter show that the graphene energy-momentum dispersion remains linear at low energy in the presence of electron interactions, while the Dirac cone angle, which is inversely proportional to the velocity, decreases (gets squeezed) at low density. (b) Renormalized velocity determined from the linear fitting of LL peak positions as described in Fig. 3(b) as a function of density for the electron (red symbols) and hole (blue symbols) puddle locations \( A \) and \( B \) at \( B = 2 \) T. The solid line shows a fit to the combined data from puddle \( A \) and \( B \) using Eq. (1) with \( v = (0.957 \pm 0.003) \times 10^6 \) m/s and \( r_s = (0.69 \pm 0.03) \) [25].

Fig. 3(b)] at a fixed density but changes slope as the density decreases. The energy dispersion is thus described by the Dirac cone in this energy range, but the cone is squeezed at low density, as schematically illustrated in Fig. 4(a).

The velocity extracted from the LL dispersion fits is shown in Fig. 4(b) as a function of density for the electron and hole puddle locations, \( A \) and \( B \). We observe a significant increase in velocity near the charge neutrality point. We fit the renormalized velocity using Eq. (1). The solid line in Fig. 4(b) shows a two-parameter fit, to the combined data of puddles \( A \) and \( B \), with a bare velocity \( v = (0.957 \pm 0.003) \times 10^6 \) m/s and an interaction parameter \( r_s = 0.69 \pm 0.03 \) [25]. The bare velocity is in agreement with previous measurements by a large number of different experiments [19,21,26–29]. The fit parameter \( r_s \) is equivalent to a measurement of the effective dielectric constant of \( \varepsilon_{BN/SiO_2} = 5.3 \pm 0.3 \) [25], using \( \varepsilon_m = (\varepsilon_{BN/SiO_2} + 1)/2 \).

In summary, we present measurements of the LL density of states of graphene on hBN as a function of spatial
position, density, and magnetic field. Tunneling spectroscopy at energies away from $E_F$ allow us to separately analyze the graphene dispersion dependence on excitation energy and density. The dispersion of the LLs as a function of orbital index shows that the linear graphene energy-momentum dispersion is retained at low energies even as the charge neutrality point is approached. However, electron interactions cause an increase of the dispersion velocity, effectively squeezing the Dirac cone, with decreasing density. The dispersion of the LLs as a function of energy and density. The dispersion dependence on excitation energy and density.

We thank S. Das Sarma, S. Adam, and M. D. Stiles for valuable discussions and S. Blankenship and A. Band for technical assistance. J. C. and S. J. acknowledge support under the Cooperative Research Agreement between the University of Maryland and the National Institute of Standards and Technology Center for Nanoscale Science and Technology, Grant No. 70NANB10H193, through the University of Maryland. Columbia work is supported under the Cooperative Research Agreement between the National Institute of Standards and Technology, Grant No. 70NANB10H193, through FCRP through C2S2 and FENA.

P. K. and A. Y. F. thank NRI INDEX program for their support.

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