

**AFRL-ML-WP-TP-2007-449**

**POLYMER NANOCOMPOSITES WITH  
PRESCRIBED MORPHOLOGY:  
GOING BEYOND NANOPARTICLE-  
FILLED POLYMERS (PREPRINT)**



**Richard A. Vaia  
John F. Maguire**

**OCTOBER 2006**

**Approved for public release; distribution unlimited.**

**STINFO COPY**

**This is a work of the U.S. Government and is not subject to copyright protection in the United States.**

**MATERIALS AND MANUFACTURING DIRECTORATE  
AIR FORCE RESEARCH LABORATORY  
AIR FORCE MATERIEL COMMAND  
WRIGHT-PATTERSON AIR FORCE BASE, OH 45433-7750**

## NOTICE AND SIGNATURE PAGE

Using Government drawings, specifications, or other data included in this document for any purpose other than Government procurement does not in any way obligate the U.S. Government. The fact that the Government formulated or supplied the drawings, specifications, or other data does not license the holder or any other person or corporation; or convey any rights or permission to manufacture, use, or sell any patented invention that may relate to them.

This report was cleared for public release by the Air Force Research Laboratory Wright Site (AFRL/WS) Public Affairs Office and is available to the general public, including foreign nationals. Copies may be obtained from the Defense Technical Information Center (DTIC) (<http://www.dtic.mil>).

AFRL-ML-WP-TP-2007-449 HAS BEEN REVIEWED AND IS APPROVED FOR PUBLICATION IN ACCORDANCE WITH ASSIGNED DISTRIBUTION STATEMENT.

\*//Signature//

RICHARD VAIA, Project Engineer  
Polymer Branch  
Nonmetallic Materials Division

//Signature//

JOHN F. MAGUIRE, Chief  
Polymer Branch  
Nonmetallic Materials Division

//Signature//

SHASHI K. SHARMA, Acting Deputy Chief  
Nonmetallic Materials Division  
Materials and Manufacturing Directorate

This report is published in the interest of scientific and technical information exchange, and its publication does not constitute the Government's approval or disapproval of its ideas or findings.

\*Disseminated copies will show “//Signature//” stamped or typed above the signature blocks.

# REPORT DOCUMENTATION PAGE

*Form Approved*  
OMB No. 0704-0188

The public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing the collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden, to Department of Defense, Washington Headquarters Services, Directorate for Information Operations and Reports (0704-0188), 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302. Respondents should be aware that notwithstanding any other provision of law, no person shall be subject to any penalty for failing to comply with a collection of information if it does not display a currently valid OMB control number. **PLEASE DO NOT RETURN YOUR FORM TO THE ABOVE ADDRESS.**

<b>1. REPORT DATE (DD-MM-YY)</b> October 2006		<b>2. REPORT TYPE</b> Journal Article Preprint		<b>3. DATES COVERED (From - To)</b>	
<b>4. TITLE AND SUBTITLE</b> POLYMER NANOCOMPOSITES WITH PRESCRIBED MORPHOLOGY: GOING BEYOND NANOPARTICLE-FILLED POLYMERS (PREPRINT)				<b>5a. CONTRACT NUMBER</b> In-house	
				<b>5b. GRANT NUMBER</b>	
				<b>5c. PROGRAM ELEMENT NUMBER</b> 62102F	
<b>6. AUTHOR(S)</b> Richard A. Vaia and John F. Maguire				<b>5d. PROJECT NUMBER</b> 4347	
				<b>5e. TASK NUMBER</b> RG	
				<b>5f. WORK UNIT NUMBER</b> M01R1000	
<b>7. PERFORMING ORGANIZATION NAME(S) AND ADDRESS(ES)</b> Polymer Branch (AFRL/MLBP) Nonmetallic Materials Division Materials and Manufacturing Directorate Air Force Research Laboratory, Air Force Materiel Command Wright-Patterson Air Force Base, OH 45433-7750				<b>8. PERFORMING ORGANIZATION REPORT NUMBER</b> AFRL-ML-WP-TP-2007-449	
<b>9. SPONSORING/MONITORING AGENCY NAME(S) AND ADDRESS(ES)</b> Materials and Manufacturing Directorate Air Force Research Laboratory Air Force Materiel Command Wright-Patterson AFB, OH 45433-7750				<b>10. SPONSORING/MONITORING AGENCY ACRONYM(S)</b> AFRL-ML-WP	
				<b>11. SPONSORING/MONITORING AGENCY REPORT NUMBER(S)</b> AFRL-ML-WP-TP-2007-449	
<b>12. DISTRIBUTION/AVAILABILITY STATEMENT</b> Approved for public release; distribution unlimited.					
<b>13. SUPPLEMENTARY NOTES</b> Journal article submitted to Polymer Science and Technology, Vol. 17, No. 5 2006. This is a work of the U.S. Government and is not subject to copyright protection in the United States. PAO Case Number: AFRL/WS 06-2651, 08 Nov 2006. Paper contains color content.					
<b>14. ABSTRACT</b> Polymer nanocomposites (PNCs), i.e., nanoparticles (spheres, rods, and plates) dispersed in a polymer matrix, have garnered substantial academic and industrial interest since their inception, ca. 1990. With respect to the neat matrix, nanoparticle dispersion has been shown to enhance physical (e.g., barrier, erosion resistance, and reduced flammability), thermomechanical (e.g., heat distortion temperature, thermal expansion coefficient, and stiffness), and processing (e.g., surface finish and melt strength) characteristics. Beyond maximization of the nanoparticle dispersion, however, the morphology of these materials is many times uncontrolled, yielding isotropic nanofilled systems, not necessarily spatially "engineered, designed and tailored" materials. To impact high-technology applications requiring unique electrical, thermal, and optical properties, manufacturing techniques enabling control of the nanoparticle arrangement and distribution must be developed. This paper will examine the status of approaches for directing the hierarchical morphology of nanoparticle dispersions in three dimensions, and beyond uniaxial alignment, using examples from the literature to highlight the potential and issues. Ultimately, two general approaches to this challenge are emerging, namely, external-in (directed patterning of nanoparticle dispersions) and internal-out (mesophase assembly of nanoparticles).					
<b>15. SUBJECT TERMS</b> Nanoparticle, polymer nanocomposites (PNCs), thermomechanical, external-in, internal-out					
<b>16. SECURITY CLASSIFICATION OF:</b>			<b>17. LIMITATION OF ABSTRACT:</b> SAR	<b>18. NUMBER OF PAGES</b> 22	<b>19a. NAME OF RESPONSIBLE PERSON (Monitor)</b> Richard Vaia
<b>a. REPORT</b> Unclassified	<b>b. ABSTRACT</b> Unclassified	<b>c. THIS PAGE</b> Unclassified			

# Polymer Nanocomposites with Prescribed Morphology: Going beyond Nanoparticle-Filled Polymers

Richard A. Vaia\* and John F. Maguire

*Materials and Manufacturing Directorate, Air Force Research Laboratory, Wright-Patterson  
Air Force Base, Dayton, Ohio 45433-7750*

*Received November 12, 2006. Revised Manuscript Received February 16, 2007*

Polymer nanocomposites (PNCs), i.e. nanoparticles (spheres, rods, and plates) dispersed in a polymer matrix, have garnered substantial academic and industrial interest since their inception, ca. 1990. With respect to the neat matrix, nanoparticle dispersion has been shown to enhance physical (e.g., barrier, erosion resistance, and reduced flammability), thermomechanical (e.g., heat distortion temperature, thermal expansion coefficient, and stiffness), and processing (e.g., surface finish and melt strength) characteristics. Beyond maximization of the nanoparticle dispersion, however, the morphology of these materials is many times uncontrolled, yielding isotropic nanofilled systems, not necessarily spatially “engineered, designed and tailored” materials. To impact high-technology applications requiring unique electrical, thermal, and optical properties, manufacturing techniques enabling control of the nanoparticle arrangement and distribution must be developed. This paper will examine the status of approaches for directing the hierarchical morphology of nanoparticle dispersions in three dimensions, and beyond uniaxial alignment, using examples from the literature to highlight the potential and issues. Ultimately, two general approaches to this challenge are emerging, namely, external-in (directed patterning of nanoparticle dispersions) and internal-out (mesophase assembly of nanoparticles).

## Introduction

The large variety of plastics available on the market today is the result of blending, that is, combining various polymers or adding micrometer-scale or larger fillers, such as minerals, ceramics, and metals (or even air). Over the past decade, the utility of inorganic nanoparticles as additives to enhance polymer performance has been established and now provides additional opportunities for many diverse commercial applications. Low-volume additions (1–10%) of isotropic nanoparticles, such as titania, alumina, and silver, and anisotropic nanoparticles, such as layered silicates (nanoclays) or carbon nanotubes, provide property enhancements with respect to the neat resin that are comparable to that achieved by conventional loadings (15–40%) of traditional micrometer-scale inorganic fillers. The lower loadings facilitate processing and reduce component weight. Most important though is the unique value-added properties and property combinations that are not normally possible with traditional fillers, such as reduced permeability, optical clarity, self-passivation, and flammability, oxidation, and ablation resistance. Beyond maximizing nanoparticle dispersion, however, the morphology of these materials is many times uncontrolled, yielding isotropic nanofilled systems, not necessarily spatially “engineered, designed and tailored” composite materials.

This review article will endeavor to “look-over-the-horizon” at the challenges and opportunities in providing the tool box to direct polymer nanocomposite morphology *in the bulk*, that is, deliver “nanocomposites-by-design”.

Among the many challenges as polymer nanocomposites (PNCs) move beyond commodity plastic applications, precise morphology control is paramount. Random arrangements of nanoparticles will not provide optimized electrical, thermal, or optical performance for many potential high-technology applications, such as dielectric underfills for electronic packaging, printed flexible electronics, engineered aerospace structural components, reconfigurable conductive adhesives, and optical gratings to just mention a few.

Specifically, we will focus on the methodologies to control the arrangement and distribution of dispersed, preformed nanoparticles beyond uniaxial alignment ( $D_{\infty}$  and  $C_{\infty}$ ). After a brief summary of the current status of PNCs (Background) and a discussion emphasizing the opportunities afforded by the ability to control nanoparticle hierarchy (Going Beyond Filled Systems), two general approaches to this challenge are explored, namely: external-in (Directed Patterning of Nanoparticle Dispersions) and internal-out (Mesophase Assembly of NanoParticles). Selected examples from the literature are used to highlight the potential and challenges; interested readers are encouraged to further explore the literature surrounding these examples for additional information. We want to emphasize that these examples are by no means intended to be inclusive. A collection of efforts that demonstrate the potential (explicitly or implicitly) is rapidly growing at the intersection of chemistry, physics, materials science, and biomaterials, including those highlighted here, as well as others, such as backfilling of sacrificial templates or in-situ nanoparticle formation within a mesophase.

\* Corresponding author. E-mail: richard.vaia@wpafb.af.mil.

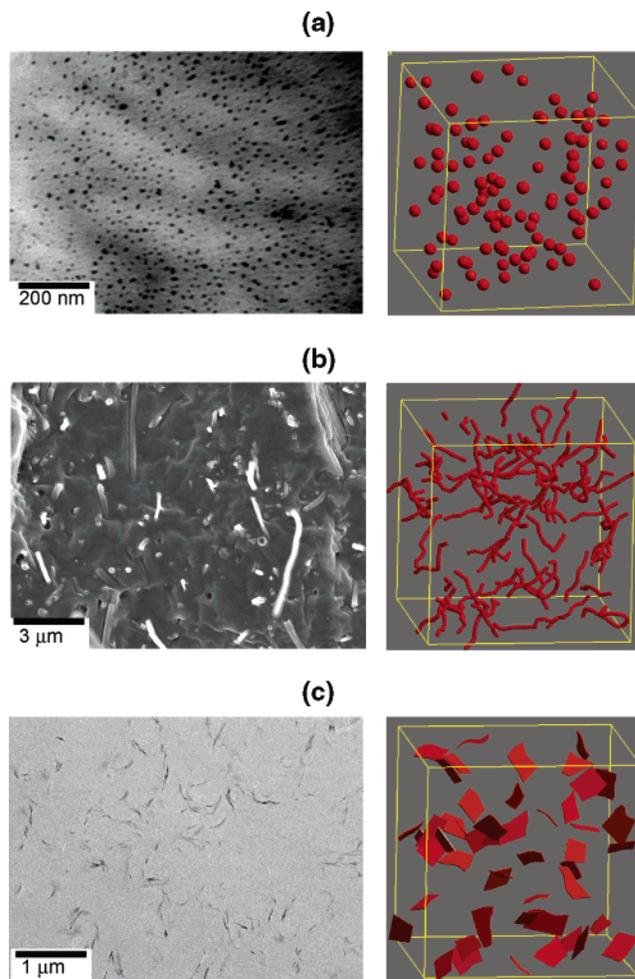
## Background

Polymeric nanocomposites have been an area of intense industrial and academic research for the past 15 years. No matter the measure—articles, patents, or research and development funding—worldwide efforts in PNCs have been growing exponentially. For example, the total number of hits for “polymer” and “nanocomposite” on SciFinder (Chemical Abstract Service (CAS) of the American Chemical Society) from 1988 to 2005 is >9400, where the yearly number has approximately doubled every 2 years since 1992.<sup>1</sup> Recent market surveys have estimated global consumption of PNCs at tens of millions of pounds (~\$250M), with a potential annual average growth rate of 24% to almost 100 million pounds in 2011 at a value exceeding \$500–800M.<sup>2–4</sup> Major revenues are forecast from large commercial opportunities, such as automobile, coatings, and packaging, where lower cost, higher performance resins would improve durability and design flexibility while lowering unit price. In light of global polymer production, which from oil alone exceeds 200–450 billion pounds annually, nanoparticle additions to plastics affords one of the commercially largest and diverse near-term applications of nanotechnology.

Since the first reports in the early 1990s<sup>5–10</sup> the term “polymer nanocomposite” has evolved to refer to a multi-component system, where the major constituent is a polymer or blend thereof and the minor constituent exhibits a length scale below 100 nm. As such, the term is sometimes used as a synonym for inorganic–organic hybrids, molecular composites, or to encompass mature commercial products, such as filled polymers with carbon black or fumed silica. The numerous reports of large property changes with very small (<5 vol %) addition of nanoparticles have fueled the view that nanoparticle addition to polymers delivers huge dividends.

Given the extensive variety of nanoparticles now commercially accessible (clays, carbon nanotubes, quantum dots, metals, silica, titania, zirconia, and various oxides, etc.), the potential combinations of polymers and nanoparticles, and thus the tailorability of the property suite, is essentially endless. The diversity in scientific investigation, technology advancement, processing innovations, and product development is staggering. A significant number of excellent review papers (e.g., clays<sup>11–18</sup> and carbon nanotubes<sup>17–21</sup>) and books<sup>22–25</sup> are available that chronicle and summarize the status of various nanoparticle–polymer combinations and the broad scientific and technological challenges still to be overcome.

Arguably, the goal for the vast majority of these investigations is to achieve increased thermomechanical performance through dispersion at the single-particle level. The resulting PNCs are treated much as an isotropic, filled polymer. Thus, from the historic perspective, nanocomposites today are really nanoparticle-filled plastics, Figure 1.<sup>26–28</sup> The use of the moniker “composites”, though, invokes strong parallels to traditional continuous fiber-reinforced composites (CFRC) and the ability to spatially engineer, design, and tailor materials performance for a given application. The payoff of CFRC manufacturing technologies is exemplified by the incredible material advancements that enable current aero-

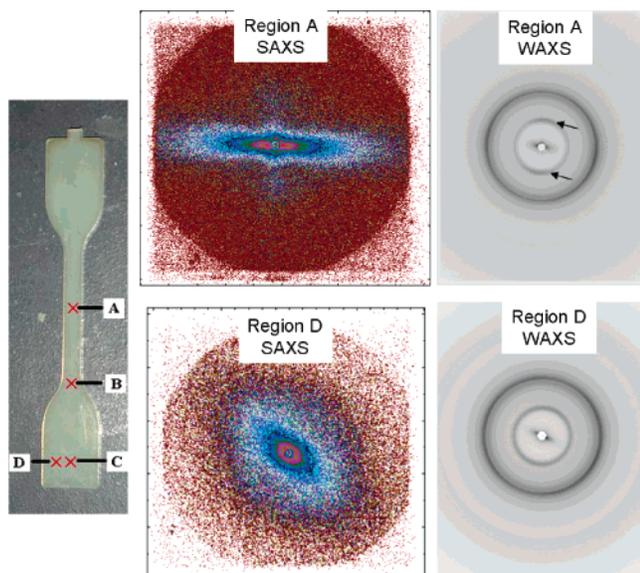


**Figure 1.** Representative polymer nanocomposite morphologies exhibiting random dispersion of spherical (0D), rodlike (1D), and platelike (2D) nanofillers. (a) Spherical: CdSe/ZnS quantum dots with carboxylic acid surface functionality (Evident Technologies) dispersed at  $5 \times 10^{15}$  particles/cm<sup>3</sup> in a recombinant silk-elastin protein (MW, 70 000).<sup>26</sup> (b) Rod: 5 wt % carbon nanofibers (Applied Science, Inc) dispersed in thermoplastic polyurethane.<sup>27</sup> and (c) Plate: 3 wt % organically modified montmorillonite (Cloisite 30A) in Epon 862/W cured epoxy.<sup>28</sup>

space systems, both military and civilian. Current processing and fabrication approaches for PNCs fall well short of this fabrication and design capability.

## Going beyond Filled Systems

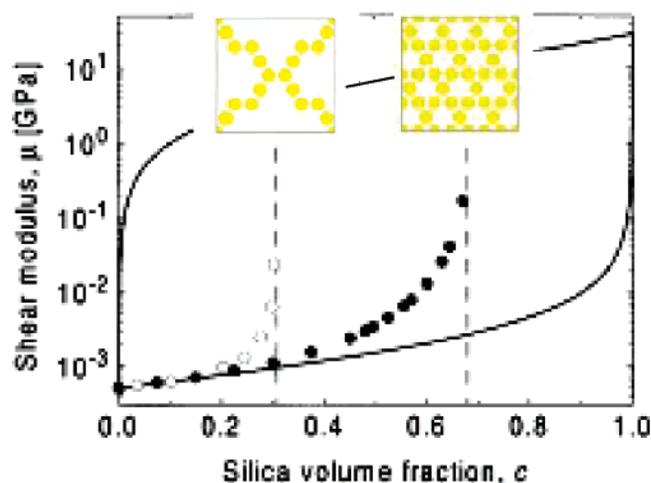
Performance enhancements of polymer nanocomposites capitalize on advantages afforded by up to a 3 orders of magnitude spatial refinement of morphology relative to traditional micrometer-scale filled polymers and composites. This contrasts nanotechnology in electronics, optics, and data storage, where the nanoscale provides access to new physical processes based on quantum phenomenon. Many discussions considering the implications and physical manifestation of this refinement of PNC morphology can be found in the literature.<sup>29–33</sup> For polymers, many bulk properties are related to the size of the polymer chain, which is characterized by the radius of gyration,  $R_g$  (~2–20 nm). The dominate length scale of the morphology becomes critically important as the dimensions of particle and polymer, as well as the interfacial curvature and interparticle distance, become comparable. At



**Figure 2.** Example of the impact of processing on macroscopic distribution of local morphologies in PNCs. Small-angle and wide-angle X-ray scattering (SAXS and WAXS) patterns from injection-molded nylon 6/montmorillonite nanocomposites, reflecting alignment of local structure.<sup>16</sup> The neck region (A in the optical picture of the sample) contains highly aligned, well-dispersed montmorillonite layers, as indicated by the oriented streak in SAXS. A weak peak on the meridian of the WAXS patterns arises from well-oriented  $\gamma$ -phase crystal lamellae (see arrows). Within the base region (region D), the montmorillonite layers reflect the fountain flow pattern of the injection molding. The polymer crystallites are not highly aligned, however, as indicated by the more uniform azimuthal scattering in the WAXS pattern from region D.

this point, the propensity of interface and cooperativity between particles dominate macroscopic properties. Strong fundamental parallels can be drawn with efforts on thermal-mechanical characteristics of ultrathin films of polymers<sup>34,35</sup> or of fluids confined within nanopores.<sup>36</sup> Schadler and co-workers recently provided strong experimental evidence for these parallels through the depression of glass transition temperature in silica-filled polystyrene.<sup>31</sup>

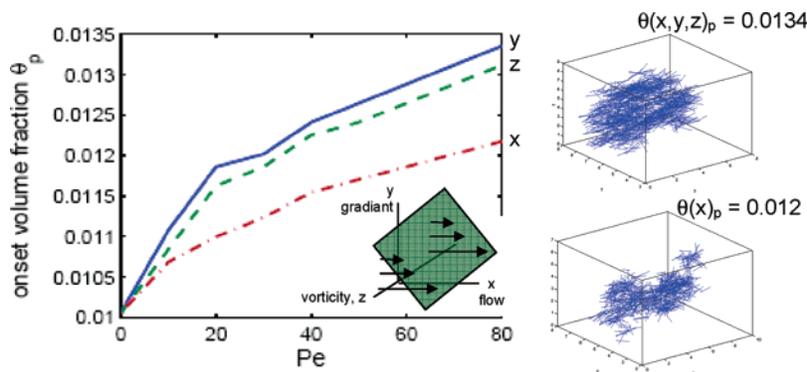
For the final material performance, though, the extent, spatial arrangement, and ordering of the constituents is as important, if not more so, than simply the degree of refinement of the morphology or the role of the interface region on polymer conformation and dynamics. For example, morphology hierarchy is well-acknowledged in the processing (injection molding, casting, etc.) of semicrystalline polymers and polymer blends where macroscopic variations, such as skin-core and fountain-flow patterns, arise from specific processing geometry and impact final part design and performance. Comparable processing issues also ultimately determine the utility of many high-performance liquid crystalline polymers. For current nanoparticle-filled products, the hierarchy of morphology is even more important than comparable micrometer-scale filled polymers due to the extreme aspect ratio of many nanoparticles, such as exfoliated clays and carbon nanotubes. Kojima and co-workers were one of the first to document and discuss the impact of processing conditions on the fine structure of PNCs.<sup>37,38</sup> They demonstrated that the relative orientation of montmorillonite layers and the concomitant impact on nylon 6 crystallites varied with distance from the sample surface. Figure 2 summarizes X-ray diffraction studies of an injection-molded



**Figure 3.** Numerical results obtained by Gusev and Rozman<sup>40</sup> for honeycomb and weblike packing arrays of silica particles ( $E = 70$  GPa,  $\nu = 0.2$ ) in an elastomer matrix ( $K = 1$  GPa,  $G = 0.0005$  GPa), demonstrating the impact of controlled morphology on enhancing mechanical properties of a composite. The two solid lines give the predictions of the Hashin-Shtrikman variational bounds. The vertical dashed lines show the maximum packing density that can be achieved assuming weblike and honeycomb packing arrays of identical cylindrical fibers. Reprinted with permission from ref 40. Copyright 1999 Elsevier.

bar of Nylon 6–montmorillonite nanocomposite showing that the orientation of montmorillonite and polymer crystallites reflects the filling pattern of the mold. Similar observations are now common in the PNC literature. More recently, Curliss noted the substantial role of the fiber mat in globally templating the alignment of montmorillonite layers during vacuum-assisted resin-transfer molding (VARTM).<sup>39</sup> Comparison of bulk PNC epoxy to reinforced composite with PNC epoxy matrix indicated that this templated local alignment along the carbon fiber axis was beneficial. Unanticipated improvements of transverse and axial strength in the final fiber-reinforced structure were attributed to the alignment of the montmorillonite along the carbon-fiber surface. Recognition of the critical importance of quantifying and ascertaining the contribution of connectivity of NPs, the resultant structure of this “NP network,” if it exhibits fractal or self-similar features, and the associated properties of this network, such as compliance, relaxation times, junction strength, and spanning linkage flexibility, are increasing.

Modeling efforts to establish structure–performance correlations further support the need for more refined processing techniques. Reports predict huge dividends in mechanical, barrier, and electrical performance if processing could prescribe precise spatial arrangement of nanoparticles. For example, Gusev and Rozman showed that comparable shear moduli could be obtained at only half the volume fraction of particles if a weblike morphology could be generated rather than random or hexagonal arrangement, Figure 3.<sup>40</sup> Additional work by Gusev and co-workers on barrier properties<sup>41</sup> and the coefficient of thermal expansion (CTE),<sup>42</sup> and by Boyce et al.<sup>43</sup> and Buxton and Balasz<sup>4</sup> on mechanical reinforcement further point to the importance of nanoparticle arrangement in achieving maximum effect at minimum filler loading. Using continuum mechanics, Liu and Kumar examined the elastic constants of single-wall carbon nanotube



**Figure 4.** “Percolation phase diagram” of anisotropic percolation thresholds showing the relation between the onset volume fraction  $\theta_p$  (percolation threshold) and normalized shear rate  $Pe$ , for monodisperse rods of aspect ratio 50.<sup>46</sup> The orientational probability distribution function (PDF) of the rods were derived from Doi–Hess theory for flowing rigid-rod macromolecules in a viscous solvent. Representative critical spanning percolation cluster at  $Pe = 80$  corresponding to (top) 3D percolation ( $\theta_p = 0.0134$ ) with percolating paths spanning the gradient ( $y$ ), vorticity ( $z$ ), and flow ( $x$ ) direction and (bottom) 1D percolation ( $\theta_p = 0.012$ ) with percolating path only spanning the flow ( $x$ ) direction.

(SWNT) ropes and fibers and showed that shear moduli of the fiber drops precipitously as the width of the uniaxial orientation function of the SWNT increases; upward of 2 orders of magnitude for 10–20% disorder at larger fiber diameters.<sup>45</sup> This sensitivity to deviations from perfect order, and its implications to fiber spinning, parallels that known for high-performance rigid-rod polymer fibers. Very recently, Forest and co-workers considered anisotropic geometric percolation of high-aspect-ratio rod ensembles dispersed in a viscous solvent and subjected to controlled rheological flows.<sup>46</sup> Results indicated that the spanning dimension of the percolating cluster can be controlled using a combination of volume fraction and shear rate. Percolation phase diagrams result, in which shear rate produces transitions between 3D, 2D, and 1D percolating clusters, as well as the loss of percolation at a given volume fraction, Figure 4. Numerous applications, including smart materials, piezo- and pyroresistive sensors and actuators, are enabled by the production of films with anisotropic conductivity arising from controlling the structure of the percolation network.

What is the status then of processing techniques that will provide a feedback to, and a demonstration of, these insights? It seems clear that more than uniaxial control of rods and plates is necessary.

Ultimately, two general approaches to this challenge, paralleling nanofabrication concepts, are emerging, namely, external-in (top-down) and internal-out (bottom-up), Figure 5. For external-in, directed patterning of nanoparticle dispersions (DPND) relies on the creation, by an external means, of a multidimensional morphology directing potential, such as a spatially varying field or susceptibility within the material. This transforms a random distribution to a prescribed, ordered construction. For internal-out, mesophase assembly of nanoparticle (MANP) relies on the ability to tailor interparticle interactions, both particle–particle and particle–matrix, to result in thermodynamically stable (and defined) mesophases. The possibilities for higher order structures are bolstered by the continuing successes in the reproducible production of nanospheres, fibers, plates, and other geometries with manufacturing tolerances that approach the current norm of micrometer-scale fibers, colloids, and films (<1%).<sup>47–52</sup> These will provide nanoscale building blocks approaching the precision of molecules, macromol-

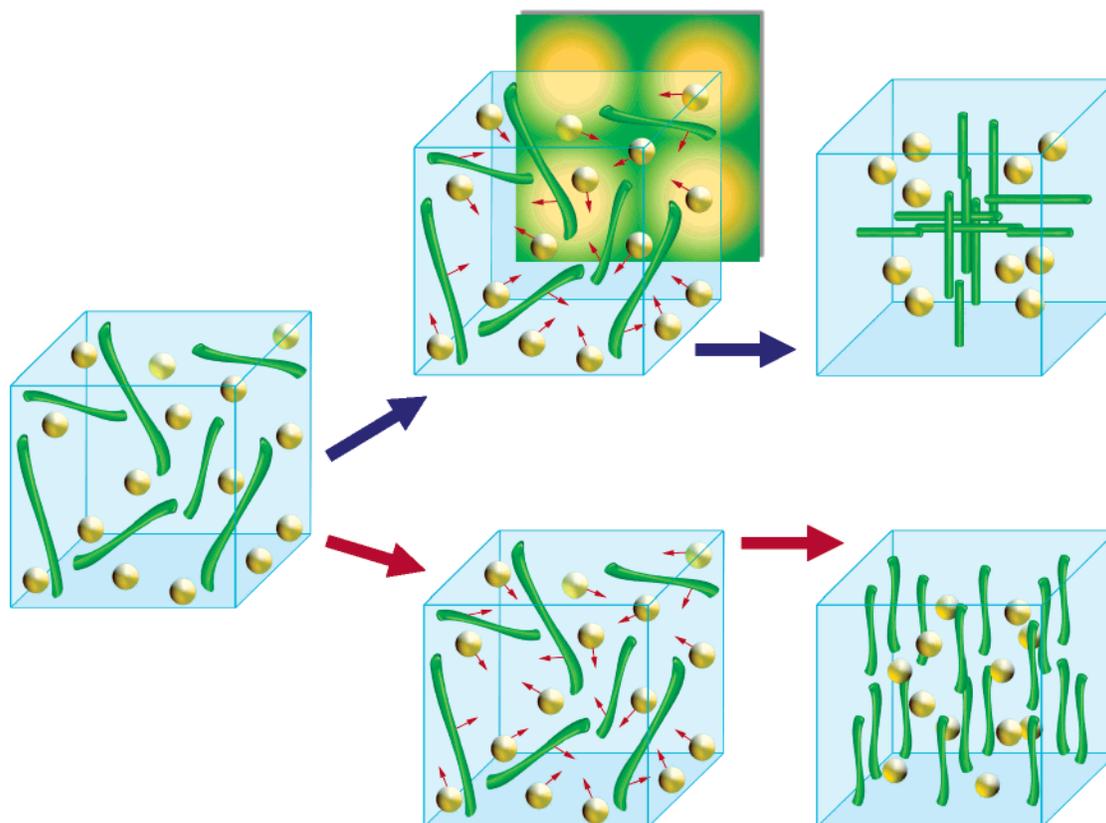
ecules, and biomacromolecules and thus access to the associated complex phase space. Note that these concepts are not mutually exclusive, but can be combined, e.g., directed patterning of mesophase assembled NPs.

### Directed Patterning of Nanoparticle Dispersions

The ability to uniaxially align nanoelements, both plates and tubes, using external forces and gradients has been extensively demonstrated. Approaches include sedimentation,<sup>53</sup> spin coating,<sup>54</sup> mechanical deformation (fiber spinning,<sup>55–58</sup> film blowing,<sup>59</sup> injection molding,<sup>60</sup> and shear<sup>61,62</sup>), magnetic fields,<sup>63,64</sup> and electrical gradients.<sup>65</sup> One-dimensional, out-of-plane periodicity has also been shown by sequential deposition approaches such as electrostatic or hydrogen bond mediated self-assembly.<sup>66,67</sup>

Much less work has been directed toward establishing robust processing techniques that enable broad tunability of two- or three-dimensional structures within the bulk. Creation of multidimensional structures is challenged by how to controllably generate a multidimensional, morphology-directing potential within the bulk material. For example, if the control input is mechanical deformation, morphology manipulation will depend not only on the magnitude and gradient of an anisotropic external deformation (all components of the stress tensor) but also on the local distribution of the stress field and on the coupling to, and interplay of, the nanoparticle’s shape and mechanical response (buckling, fracture, etc.), surface energies, and viscoelastic properties, etc. Extensive efforts on the refinement of two-component polymer blends exemplify the complexity of this specific case.<sup>68</sup>

The following sections highlight four areas which provide an overview of current capabilities. The initial three sections (Mechanical Deformation, Electric and Magnetic Fields, and Optical Fields) discuss the coupling between applied external fields and nanoparticle orientation, distribution and interparticle interactions. The last section (Multi-component Interfacial Systems) considers the possibilities afforded by interfacial segregation of nanoparticles in an immiscible blend, and thus templating the nanoparticle distribution to the interfacial regions of the immiscible blend, and subsequent macroscopic manipulation of the structure



**Figure 5.** Schematic representation of the two general approaches to control nanoparticle distribution and arrangement beyond random order: external-in (top) and internal-out (bottom). For external-in, directed patterning of nanoparticle dispersions (DPND) relies on the creation, by an external means, of a multidimensional morphology directing potential (top center: graded background contours), such as a spatially varying field (depicted above) or susceptibility within the material. Imposing this complex field on the sample (top center) transforms a random distribution to a prescribed, ordered construction due to mass flow of nanoparticles (arrows) to minimize the potential energy of the system within the externally applied field. For internal-out, mesophase assembly of nanoparticles (MANP) relies on the ability to tailor interparticle interactions, both particle–particle and particle–matrix, to result in a thermodynamically stable (and defined) mesophase. The phase behavior of the system may be modulated (bottom center) by uniform changes in the systems intrinsic (pressure or temperature) or extrinsic (number density or entropy) thermodynamic parameters. In contrast to the externally patterned potential applied for DPND, changes in particle organization in MANP occur in response to a uniform change in the system’s environment.

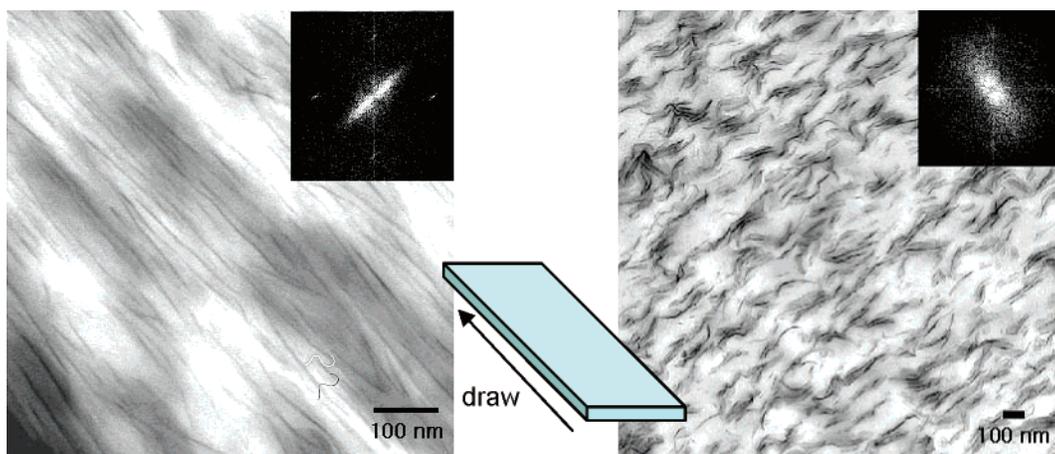
via distortions and refinements of the underlying multiphase morphology.

**Mechanical Deformation.** For mechanical deformation, uniaxial or in-plane biaxial arrangement of the nanoparticle and polymer crystallites is common, as noted by the aforementioned approaches. Multidimensional control of the morphology by complex strain fields is rarely reported. However, some reports of using controlled strain inputs to introduce local distortions of the nanoparticle can be found. For example, beginning from quenched biaxially extruded films of nylon 6/montmorillonite nanocomposites, Park and co-workers observed within zone-drawn films that the montmorillonite layers buckle perpendicularly to the draw direction, analogous to failure of a uniaxially strained sheet of paper.<sup>69</sup> The failure mode appeared to occur for a collection of parallel aluminosilicate layers (2–4) and exhibited a mean spatial frequency, Figure 6. This behavior has analogies to shear banding and deformation of uniaxially aligned lamellar and cylindrical block copolymes.<sup>70–72</sup>

In general, nanoparticle alignment, as well as morphology refinement, depends on the type of imposed flow, whether extensional, shear, or mixed.<sup>73</sup> This is analogous to other complex fluids including polymers. Fundamentally, the complexity of the local response to imposed stress or strain not only depends on the interfacial strength, particle size,

and mechanical characteristics of the constituents but also on elastic instabilities, such as buckling, of the high aspect ratio nanoparticles, as well as the cooperative response of nanoparticles coupled through overlapping local strain fields. The extent of the latter is many times reflected in a lower volume fraction for mechanical percolation ( $\epsilon \rightarrow 0$ ) than electrical percolation. Systematic exploration of this complex processing space to produce hierarchical structures by inducing particle orientation and particle deformation, as well as transferring insights from shear thinning (pseudo-plastic) and shear thickening (dilatant) fluids, is still in its infancy.

**Electric and Magnetic Fields.** The morphological origins of stiffening of rheological fluids, such as electro- (ER) and magneto- (MR), provide a perspective on the possibilities for externally applied, spatially patterned, electrical, and magnetic fields for complex PNC fabrication. Under the electric field, particulate additives, such as oxides or metals, reversibly form fibrous structures within the nonconducting fluid. These controlled aggregates are parallel to the applied field and can increase the viscosity of the ER fluid by a factor of up to  $10^5$ .<sup>74</sup> Similar morphologies are found for magnetorheological fluids where particle–particle interactions, which arise from induced magnetic dipoles, lead to chaining and subsequently long-range, periodic ordering of the particle chains parallel to the lines of magnetic flux.<sup>75</sup> Gradients of

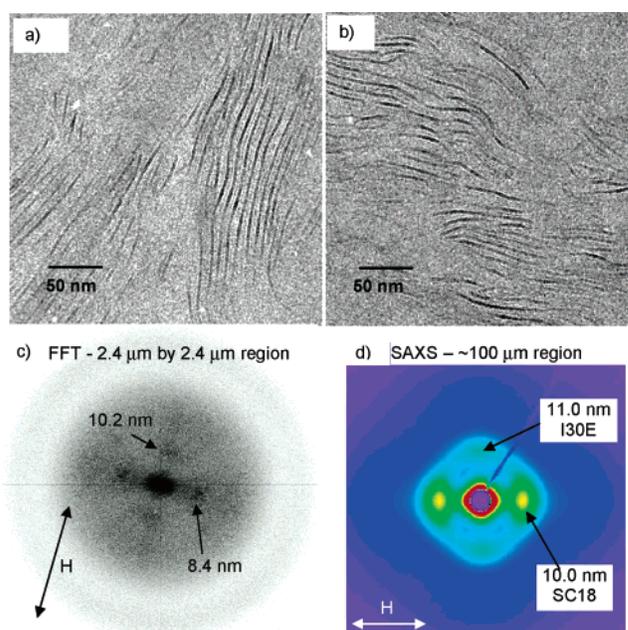


**Figure 6.** Transmission electron micrographs from Park et al.<sup>69</sup> of a biaxially protruded and then drawn film of nylon 6/montmorillonite nanocomposite perpendicular (left) and parallel (right) to the draw direction demonstrating buckling of the nanolayers. Reprinted with permission from ref 69. Copyright 2005 American Chemical Society.

magnetic or electric fields also generate lateral forces that can be used to pattern nanoparticles. For example, electrophoresis approaches have been extensively developed in the biological community for nucleic acid and protein purification<sup>76</sup> and di-electrophoresis for nanofabrication to align particles between electrodes or trap nanoparticles in specific regions.<sup>77</sup> Also, local instabilities at the interface between two fluids can be enhanced with applied electrical or magnetic fields, providing a means to control the growth of certain Fourier components of the spatial frequency of the composition modulation.<sup>78,79</sup>

Even with these proven successes though, direct application to large-scale PNC production is challenging. Generating reconfigurable gradients in three dimensions with micrometer (or sub-micrometer) features requires complex tooling with reconfigurable, patterned electrodes. As particle size decreases, the fields necessary to overcome thermal randomization increase considerably,<sup>80,81</sup> and may approach the breakdown strength of the matrix. This drastically limits sample thickness. Also anisotropic particles (rods and plates) drastically increase the complexity (and predictability) of field response with respect to the more commonly used spherical shapes. Finally, for thermoplastics, high viscosity results in long mass transport times, even over short sub-micrometer distances. Nevertheless, adaptations of these concepts to PNCs, especially thermosets where order can be “stabilized” via postprocessing polymerization, are now providing intriguing possibilities to move beyond uniaxial alignment.

As an example, Koerner and co-workers<sup>63</sup> exploited the orthogonal magnetic susceptibility of montmorillonites (MMTs) from different mineral deposits to fabricate a three-dimensional morphology composed of orthogonal alignment of aluminosilicate layers using a uniaxial external magnetic field. Depending on the source, MMTs exhibit remnant magnetization arising from antiferro- and ferrimagnetic impurities and align with layers parallel or perpendicular to the field, Figure 7. Within a few minutes, application of static magnetic fields (1.2 or 11.7 T) induces stable alignment of organically modified MMT within an epoxy resin at room temperature. Structural relaxation in the absence of the field is orders of magnitude slower, enabling the alignment to be



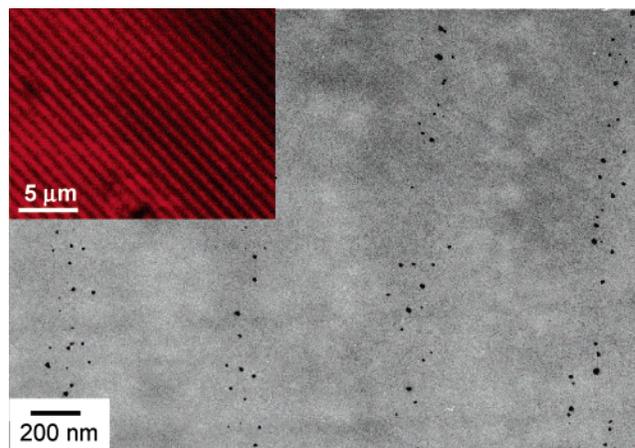
**Figure 7.** Magnetic alignment of an equal mixture of two octadecyl ammonium modified montmorillonites (SC18 and NC18) in epoxy (6 wt % total).<sup>63</sup> High-resolution transmission electron micrograph of regions of the nanocomposite with (a) SC18 layers parallel to and (b) NC18 layers perpendicular to the applied magnetic field. (c) Digital Fourier transformation of a transmission electron micrograph of a  $2.4 \mu\text{m} \times 2.4 \mu\text{m}$  area showing correlation peaks arising from the independently aligned montmorillonites. (d) Small-angle X-ray pattern showing a four-point pattern caused by orthogonally aligned MMT tactoids.

captured during the subsequent cure. This could offer a pathway to a viable manufacturing technology. Thermal mechanical measurements demonstrate that this morphology manipulation impacts the coefficient of thermal expansion (CTE), decreasing CTE by an additional 20–30% in the direction of maximum MMT alignment with respect to that of the isotropic PNC. This example demonstrates that through the use of two or more nanoparticles with different susceptibilities to the externally applied field, the complexities associated with applying non-uniaxial fields to the sample could be overcome. In support of this concept, aligned arrays along the electric field direction of binary blends of silver and silver oxide nanoparticles ( $\sim 8.5$  nm) in poly(methyl methacrylate) film has recently been demonstrated,<sup>82</sup> provid-

ing further glimpses at the possibilities for complex internal structures when nanoparticle blends are used. Other reports discussing CNTs in magnetic fields<sup>83</sup> and electric fields<sup>84</sup> are available. Comparable results have also been reported for electric field assisted alignment of MMT in epoxy.<sup>65</sup> As with dual frequency liquid crystals, frequency-dependent differences in the direction and magnitude of the induced dipole on anisotropic nanoparticles may afford benign tuning of the distribution;<sup>85</sup> however, this has yet to be demonstrated.

In addition to directional reinforcement to maximize thermal-mechanical stability, these processing approaches provide avenues to create polymer nanocomposites with uniquely tailored dielectric ( $\epsilon = \epsilon' + i\epsilon''$ ) properties for sensors as well as electronics and power applications. The addition of nanoparticles have been shown to increase dielectric breakdown strength and durability of polymers under extreme electric fields, opening new possibilities for ultrahigh-energy density capacitors with microsecond discharge rates and ultradurable electrical insulation materials,<sup>86,87</sup> especially in consideration of the availability of anisotropic oxide nanocrystals with inherently large dielectric constants, such as ZnO and BaTiO<sub>3</sub>. Also, these approaches may provide the necessary control to create aligned yet internally interrupted metal nanostrands within an insulating matrix. Saha and Chakravorty have argued that these nanostructures exhibit very low internal depolarization field and thus exhibit a Gor'kov-Eliashberg effect, yielding  $\epsilon$  from  $10^5 - 10^{10}$ .<sup>88,89</sup>

**Optical Fields.** Holographic photopolymerization has also been shown as a viable nanoparticle directing tool, creating periodic, multidimensional arrangements of gold nanoparticles, colloids, layered silicates, TiO<sub>2</sub>, ZnO<sub>2</sub>, and CdSe.<sup>90–92</sup> Interference of two or more coherent laser beams within a photoreactive monomer syrup results in a periodic intensity distribution that initiates a self-similar periodic, polymerization process. A higher intensity results in a locally faster polymerization rate. This distribution in polymerization rate causes a spatially periodic distribution of high molecular weight polymer, which changes over time. The temporal evolution of the polymer can locally induce phase transitions, such as demixing, or result in mass flow, which concentrates (or traps) nanoparticles in regions of low (or high) intensity depending on particle size and surface reactivity. Figure 8 summarizes one such fabricated structure generated using two-beam interference lithography—planes of gold nanoparticles oriented perpendicular to the film surface. More complex structures paralleling those demonstrated for holographic polymer dispersed liquid crystals<sup>93,94</sup> should be possible with four- and six-beam holography, such planar trigonal and orthogonal lattices with parallel rods, or three-dimensional cubic or orthorhombic P structures. These more complex structures are of intense interest for technologies utilizing photonic crystals (PCs) and photonic band gap materials, such as switchable filters, laser cavities, superprisms, omnidirectional reflectors, and platforms for optical-based chemical sensing. Single-step holographic routes to PCs that enable precise arrangement of optoelectronic nanoparticles, such as quantum dots or plasmonic metallic particles, with electrooptic constituents, such as liquid

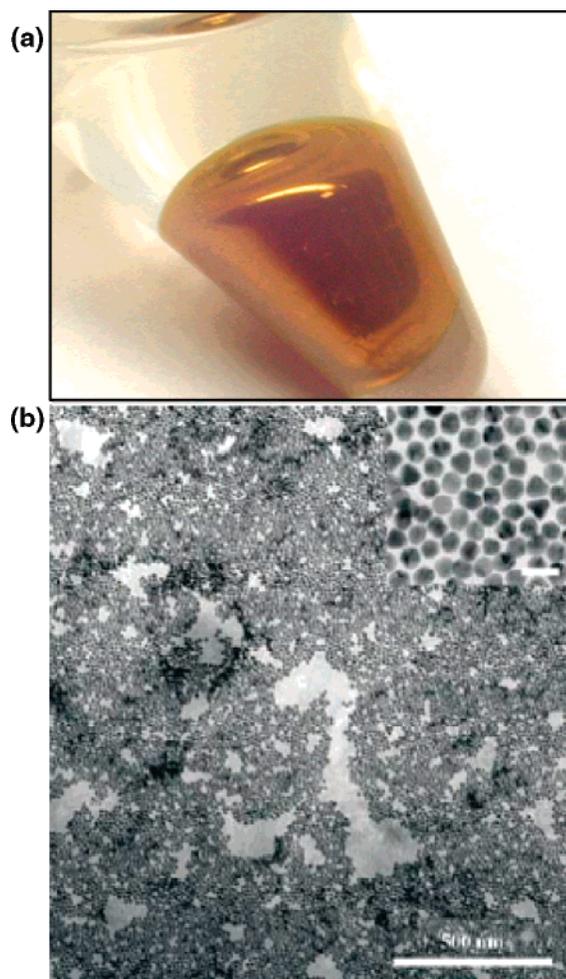


**Figure 8.** Holographic photopolymerization patterning of gold nanoparticles in an acrylate resin.<sup>90</sup> The transmission electron micrograph shows a cross-section of planes of gold nanoparticles on an  $\sim 450$  nm period that are oriented perpendicular to the substrate. The planes of Au nanoparticles span the  $10 \mu\text{m}$  film thickness. The inset shows an optical confocal microscopy image of the surface of the transmission grating formed by these patterned gold nanoparticles.

crystals, would provide exciting alternatives to multistep template and infiltration fabrication concepts. Initial efforts along these lines are beginning to be reported.<sup>95,96</sup> Many other applications predicated on the ability to simultaneously control periodic structures and complex nano- and microstructures can be envisioned, ranging from tailored membranes for separation and catalysis to unique tissue scaffolds, vibration isolation mounts, and phononic crystals.

**Multicomponent Interfacial Systems.** Immiscibility is common in many multicomponent systems, both polymers and nanoparticles. The interface between the immiscible phases affords opportunity for structural organization of nanoparticles. Conceptually, this derives from a common approach to compatibilization of immiscible blends—the addition of a small quantity of a third component that is miscible with both phases (cosolvent), has an interfacial energy with each of the two constituents that is less than that between the constituents individually, or is a structured amphiphile whose sections are miscible with each component respectively.<sup>97</sup> The additive is designed such that segregation to the interface is energetically favorable. The process leads to reduction of the interfacial tension, facilitating refinement of the phase domains; stabilization of the morphology against high stress and strain processing (e.g., in injection molding); and/or enhancement of adhesion between the phases in the solid. The distribution of interfacial regions within the final materials depends on process history and thus provides a means of tailoring the macroscopic arrangement of nanoparticles.

As an example, layered silicates have been discussed as compatibilizers for immiscible polymer blends, and thus demonstrating the potential to sequester at the interface.<sup>98,99</sup> Ahn et al.<sup>98</sup> observed that when a very small amount of organoclay (1%) is added to an immiscible polyethylene–polybutylene terephthalate blend, the organoclay is located at the interface and the organoclay tactoids are disrupted into thinner tactoids of some tens of nanometers. The presence of organoclay at the interface hydrodynamically stabilizes the blend morphology by suppressing the coalescence of the



**Figure 9.** (a) Photograph of the self-assembled Au nanoparticles at the water/toluene interface in a plastic Eppendorf tube from Mohwald et al.<sup>104</sup> The tube has been tilted such that the colored area corresponds to the water/toluene interface. (b) TEM image (scale bar, 500 nm) of the monolayer of 12 nm Au nanoparticles formed at the water/toluene interface. Inset: high-magnification TEM picture (scale bar, 25 nm). Reprinted with permission from ref 104. Copyright 2004 Wiley-VCH.

droplets and also makes the two-phase polymer blend morphology more thermally stable. Sequestering montmorillonite at the interface in immiscible polystyrene/polypropylene blends has also been observed.<sup>100</sup>

Pickering emulsions (emulsions stabilized by solid particles) using nanoparticles also offer interesting possibilities. Pickering emulsions are encountered in various natural and industrial processes such as crude oil recovery, oil separation, cosmetic preparation, and wastewater treatment. An initially high interfacial energy between two phases, such as oil and water, can be decreased by the assembly of the particles at the interface.<sup>101,102</sup> The decrease in surface energy favors the formation of a monolayer of nanoparticles located at the interface, Figure 9. Additionally, the relatively low viscosity of the two phases implies the process is highly dynamic, which enables errors (defects) to be corrected rapidly. Russell and co-workers<sup>103</sup> estimated that the stabilization energy,  $\Delta E$ , decreases as  $r^2$ , where  $r$  is the radius of the particle. For a water/oil mixture,  $\Delta E \approx 1 \times 10^5 k_B T$  for 100 nm particle and  $\Delta E \approx 1 \times 10^1 k_B T$  for 1 nm particle. Using CdSe nanoparticles of two different radii (2.7 and 4.6 nm), they not only showed CdSe segregation to the interface of an oil–

water droplet but a preferential adsorption of the larger particle. Mohwald and co-workers<sup>104</sup> have created close-packed nanoparticle layers and nanoalloys using Au, Ag, or  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>. Robust and water dispersible colloidosomes with shells predominantly composed of monolayers of liquid-like, close-packed magnetite nanoparticles has also been created through the gelation of an aqueous phase with agarose after assembly of magnetite nanoparticles at the interface.<sup>105</sup> Ashby and Binks<sup>106</sup> reported Pickering emulsions using laponite. Additionally, *N,N*-dimethylformamide–water interface has been used to assemble carbon nanotube–metal nanoparticle composite materials;<sup>107</sup> and aqueous acrylamide (AAm) or 2-hydroxyethyl methacrylate (HEMA) in cyclohexane stabilized by hydrophobic Cloisite 20A (MMT20) lead to polymer–clay nanocomposite latex particles.<sup>108</sup>

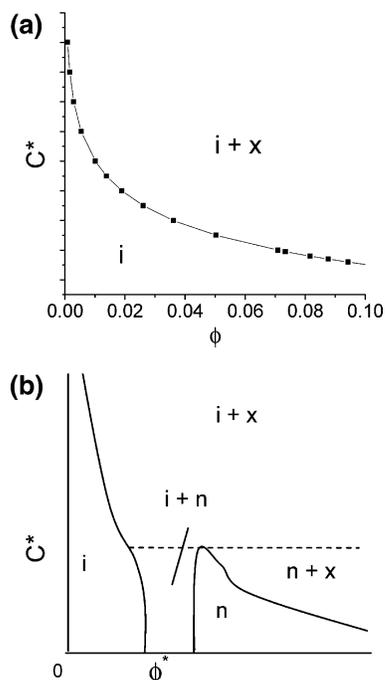
The ability to arbitrarily tune the surface properties of the nanoparticle increases the range of potential two-phase systems and morphologies amenable to interfacial sequestration. Recent efforts on mixed surfactant passivation have demonstrated this potential. For spherical particles, mixtures of polymers<sup>109</sup> or organic molecules<sup>110</sup> as passivation yielded nanoparticles with amphiphilic character. These Janus-like nanoparticles can spontaneously assemble into higher-order organic–inorganic structures, such as cylindrical micelles.<sup>111</sup> In many regards, this conceptually parallels substantial synthetic work on multiarm miko (heteroarm) polymers.<sup>112,113</sup>

In general multicomponent interfacial systems share many common features with nanoparticle sequestration in mesophases discussed below, differing mostly in the composition and morphology of the polymer media, the density of the interfacial area, and the number density of nanoparticles.

### Mesophase Assembly of Nanoparticles

Although in principle directed patterning of nanoparticle dispersions provides access to any arbitrary geometry, manufacturing complexities establish practical limits on the thickness and volume of the material impacted. These include the excessive time for pattern development in highly viscous media such as polymers, complex external tooling with vibration isolation to maintain structural precision on the sub-micrometer scale for nonbatch manufacturing, and, in some instances (electrical and optical), uniformly achieving the energy density requirements throughout thicker ( $> 100 \mu\text{m}$ ) samples. To complement direct patterning of nanoparticle dispersion (DPND) approaches, the inherent organization derived from a thermodynamically well-defined phase provides an alternative approach to the fabrication challenge.

The phase behavior of hard-body particles has been extensively developed and serves as a starting point from which the properties of nanoparticle–polymer combinations can be considered.<sup>114–116</sup> These idealized systems are defined when the constituent particles interact exclusively through short-ranged repulsive forces at the point of particle contact, with no dissipative behavior. For angular particulates whose shape deviates significantly from spherical, entropically driven excluded volume effects give rise to a wide range of complex ordering and phase behavior including the isotropic–nematic phase transition in liquid crystalline systems.<sup>117</sup> The increased local order within the nematic phase increases the



**Figure 10.** Simplified schematic representations of the phase diagrams for both spherical (a) and rodlike particles (b) of volume fraction,  $\phi$ , dispersed in a polymer of reduced concentration,  $c^*$ . Spherical particles at low concentration may be highly dispersed with little agglomeration (i). As the concentration rises depletion forces become increasingly important and give rise to agglomeration and phase separation ( $i + x$ ). As seen in b, the situation for rodlike particles is considerably more complex in that even very small concentrations of rods can give phase separation, extended two-phase coexistence regions, and critical phenomena. Orientational correlations result in the formation of a nematic phase (n), which for various volume fraction of rods may be in equilibrium with the lower volume fraction isotropic phase (i) as well as the higher volume fraction aggregate structures (x).

number of translational and rotational eigenstates within the system giving rise to a spontaneous entropically driven transition with negative free energy. The phase behavior of such systems, including rods and discotics (discs), are critically important to a broad range of technologies such as liquid crystals, high-performance polymers, and biomolecules. Temperature dependence and phase behavior of such systems have been treated theoretically via the early formulations of Flory and Abe<sup>118</sup> and Onsager,<sup>119</sup> and more recently by de Gennes and Proust.<sup>120</sup>

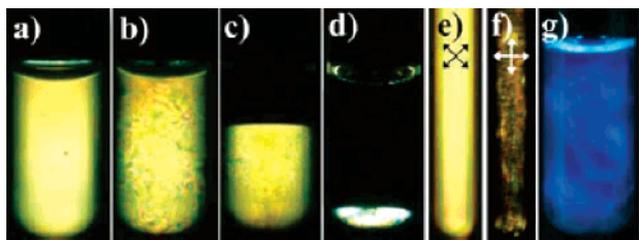
These general topological considerations imply nanoparticle–polymer (or reactive monomer) blends conceptually will show a great structural richness with the possibility to form a variety of exotic phases. Not only do the usual enthalpic and entropic interactions contribute to the chemical potentials, but now the topological constraints imposed by size and geometry and the added complexity of surface interactions in what are thermodynamically small systems need to be taken into account. Subtle changes in any of these parameters may tip the balance, causing major perturbation on the locus of the curve of equal chemical potential, Figure 10. For a realistic nanoparticle–fluid system, the extent of the immiscibility window between a low- and high-volume fraction nanoparticle phase will depend on numerous inter-related contributions to the free energy in addition to the configurational entropy of the constituents. These include intermolecular interactions as well as the relative impact of

mixing on the internal degrees of freedom of the constituents (e.g., chain conformation). Additionally, the experimentally accessible region in this phase space will depend on the magnitude of the intensive variables (temperature, pressure, and chemical potential, etc.), polydispersity of the constituents, and kinetic considerations such as gelation.

The possibilities afforded by discrete phase behavior of polymer nanocomposite systems may be systematically considered on the basis of how the polymer and nanoparticle are combined to form the constituents of the mixture. To contrast the morphologies created by traditional nanoparticle dispersion, we'll discuss three general types: *nanoparticle–macromolecular systems* comprised of nanoparticles within a complex fluid, such as linear polymers, block copolymers, dendrimers, or liquid crystals; *structured block nanoparticles* that are analogous to block copolymers or macromolecular stars, but where one “block” or the core is the nanoparticle and the other block or arms is a polymer; and *nanoparticle–nanoparticle systems* containing mixtures of nanoparticles of various sizes and shapes with organic or macromolecular coronas.

**Nanoparticle–Macromolecular Systems.** The examination of the phase behavior of nanoparticle dispersion in small molecule or macromolecule medium is rapidly expanding. Approaches range from ordered phases arising from the combination of functional nanoparticles and linear polymer to the use of the phase structure of a block copolymer to template nanoparticle organization.

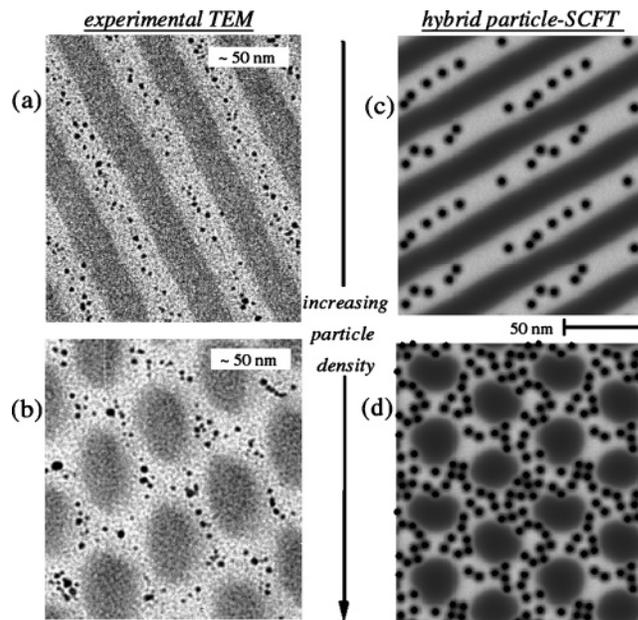
The similarity between the size of the nanoparticle and the chain results in coupling between the configurations of the constituents (mixing entropy), the internal degrees of conformational freedom of the polymer, and intermolecular and interparticle interactions. This complicates the unambiguous development of models to predict nanoparticle dispersion and phase behavior. Nevertheless, the importance of various factors is beginning to be more widely appreciated. Hooper and Schweizer<sup>121,122</sup> as well as Zhong et al.<sup>123</sup> argue that for athermal mixtures that there exists oscillatory depletion forces between two hard spheres due to monomer-level packing correlations and the key geometric variable determining nanoparticle aggregation in flexible, high molecular weight polymers is the ratio of particle to monomer diameter. Thus, purely athermal mixtures of polymers and nanoparticles may be macroscopically phase-separated at equilibrium, requiring attractive polymer–particle interactions to achieve miscibility. Recently, simple thermodynamic arguments based on a Flory formulation of the free energy provided a framework to Mackay and co-workers to explain the miscible–immiscible boundary for a C<sub>60</sub>–polystyrene system.<sup>124</sup> The formulation included distortion of the polymer chain configuration by nanoparticles and a relative increase in accessible surface area of the nanoparticle due to packing considerations within the nanoparticle aggregate. These factors lead to a maximization of the miscibility window for an optimal nanoparticle size. Finally, Ginzburg and Balasz combined self-consistent field theory (SCFT) with density functional theory to calculate the equilibrium behavior of polymer–clay mixtures, demonstrating the potential to form isotropic, nematic, smectic (lamellar), columnar, crystal, and



**Figure 11.** (a–g) Phase behavior of aqueous suspensions of  $\text{H}_3\text{Sb}_3\text{P}_2\text{O}_{14}$  single-layers from Gabriel et al.<sup>128</sup> observed between cross-polarizers (the isotropic phase appears dark): (a) 2 mL of birefringent gel phase ( $\phi = 1.98\%$ ; topological defects are so dense that the texture appears homogeneous at the scale of the photograph); (b) 2 mL of birefringent fluid phase ( $\phi = 0.93\%$ ); (c) 2 mL of a biphasic sample with an overall volume fraction  $\phi = 0.65\%$ ; (d) 2 mL of a biphasic sample with an overall volume fraction  $\phi = 0.03\%$ ; (e and f) magnetically aligned sample observed in a 5 mm NMR tube that has been immersed 10 min in a 18.7 T field at 50 °C, in two different orientations compared to the polarizer/analyzer system represented by arrows ( $\phi = 0.75\%$ ); (g) sample iridescence ( $\phi = 0.75\%$ ) observed in natural light and due to light scattering by the  $\text{H}_3\text{Sb}_3\text{P}_2\text{O}_{14}$  layers stacked with a period of 225 nm. Reprinted with permission from ref 128. Copyright 2001 Nature Publishing Group.

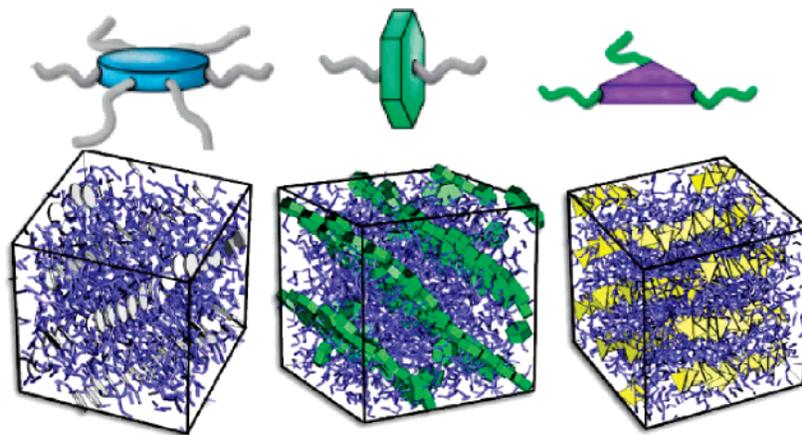
plastic solid (house-of-cards), as well as a two-phase (immiscible), mixtures.<sup>125</sup> They also examined the impact of surface architecture (tethering density, molecular weight, and composition) of polymers end-tethered to the clay surface.<sup>126</sup> Most of these predictions though have yet to be systematically investigated in experimental montmorillonite–thermoplastic systems.

Experimentally, however, the occurrence of complex mesophases of nanoparticle dispersions is well-established, although most systems examined to date are in a small molecule medium. For example, experimental phase diagrams have been reported for so-called lyotropic, inorganic liquid crystals.<sup>127</sup> Mesophase formation has been reported for molecular nanowires, tubes, ribbons, and rods ( $\text{Li}_2\text{Mo}_6\text{-Se}_6$ , imogolite,  $\text{Nb}_2\text{PS}_{10}$ ,  $\text{V}_2\text{O}_5$ , boehmite ( $\gamma\text{-AlOOH}$ ), akaganeite ( $\beta\text{-FeOOH}$ ), goethite ( $\alpha\text{-FeOOH}$ )) and nanoplatelets and disks (smectic clays,  $\text{H}_3\text{Sb}_3\text{P}_2\text{O}_{14}$ , gibbsite ( $\text{Al}(\text{OH})_3$ ), and  $\text{Ni}(\text{OH})_2$ ), Figure 11.<sup>128</sup> Experimental phase diagrams also take into consideration limits due to gelation and aggregation.<sup>129</sup> Alivisatos and co-workers have reported similar behavior for CdSe quantum rods in anhydrous cyclohexane, with anticipated applications for electrooptic devices.<sup>130</sup> Finally, taking inspiration from dryjet wet-spinning of liquid crystalline rigid-rod polymers, such as poly(*p*-phenylene benzobisoxazole) in poly(phosphoric acid), Davis et al. have demonstrated the two-phase “Flory” chimney between isotropic and nematic phases for solutions of purified single-wall carbon nanotubes in superacids, such as oleum.<sup>131</sup> As with the rigid-rod polymers, fiber-spinning from the nematic phase results in a higher degree of molecular order along the fiber axis. In all of these cases, the mesophase could be captured in a solid monolith by replacing the fluid media with a chemically reactive media, although investigations along these lines are minimal. In general, these morphologies are mostly nematic, however, lacking order beyond uniaxial. Nevertheless, these observations establish a parallel to liquid crystal technology, point to the potential for higher order mesophases, and afford yet to be explored opportunities for processing and properties of PNCs.



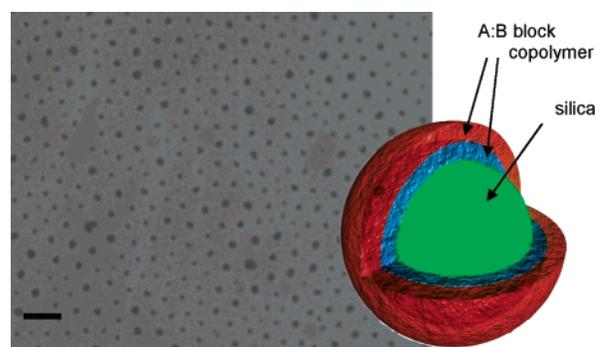
**Figure 12.** Transmission electron microscopy images from Fredrickson et al.<sup>136</sup> of  $f_{\text{PS}} = 0.50$  PS-P2VP diblock nanocomposite containing PS-functionalized Au particles with particle volume fraction (a) = 0.10 and (b) = 0.35. Total particle size including the PS shell is 2.6 times that of the Au cores seen as black dots in the TEM images. Hybrid particle field simulation results (right column) show monomer volume fractions representing PS (light), P2VP (dark), and particles (black). Simulation parameters:  $f_{\text{PS}} = 0.5$ , Flory parameter ( $\chi$ ) for PS-P2VP diblock  $\chi = 0.16$ ,  $\lambda = 0.16$ , and particle area fraction = 0.04 (c) and = 0.18 (d). The particle configurations shown are representative of those obtained on the basis of several independent runs for a given nanoparticle density. Reprinted with permission from ref 136. Copyright 2006 The American Physical Society.

Paralleling nanoparticle dispersions within a homopolymer or organic media, sequestration of the nanoparticle within one region of a single-component complex fluid, such as a block copolymer, has been demonstrated as a relatively facile route to many complex morphologies.<sup>132–135</sup> Fredrickson and co-workers have utilized a self-consistent field theory in which particle coordinates and chemical potential field variables are simultaneously updated in the simulation.<sup>136</sup> The fluid model can contain polymers of arbitrary chemical and architectural complexity, along with particles of all shapes, sizes, and surface treatments. The initial simulations of polystyrene-poly(2-vinylpyridine) (PS-P2VP) containing PS-functionalized Au nanoparticles paralleled experimental data, Figure 12.<sup>132</sup> Additionally, increasing Au-particle density drove the structure from the hexagonal to lamella phase; however, the SCFT results (two-dimensional) underestimated the particle density at the morphological transitions. Many experimental analogues to the Au–PS-P2VP system are available including CdSe nanoparticles in PS-P2VP<sup>133</sup> and montmorillonites in polystyrene-polybutadiene-polystyrene (SBS) triblock.<sup>134</sup> Overall, the architecture and composition of the corona, as well as particle size, provide fine control of the location of the nanoparticle within the preferred domain.<sup>135</sup> By tuning the nanoparticle so that they prefer the interface between the chains of the block, Kramer and co-workers have been able to drive phase transition of the diblock from lamellar to bicontinuous.<sup>137</sup>



**Figure 13.** Predicted mesophases of single-component melts of structured block nanoparticles by Glotzer et al.<sup>140</sup> (Left) Hexagonal columnar (cylinder) phase formed from edge-tethered disks. The disks are tilted with respect to the interfacial normal. (Middle) Lamellar phase formed from face-tethered hexagonal disks. The disks pack hexagonally within the sheet. (Right) Hexagonal columnar phase formed where tethers are attached to the three vertices of a triangular plate. The triangles form a twist about the cylindrical axis.

**Structured Block Nanoparticles.** Contrasting two-component systems where one is the nanoparticle and the other is the macromolecule, structured block nanoparticles are monolithic structures where the nanoparticle and macromolecule are coupled in a controlled and reproducible fashion. The phase behavior of these structured single components is exemplified by block copolymers<sup>138</sup> or liquid crystal polymers.<sup>139</sup> For example, depending on the architecture (A:B length ratio) and composition (intermolecular interactions) of the two chains in an AB diblock copolymer, microphase separation leads to many morphologies, including spheres, cylinders, lamellae, and bicontinuous structures (double-diamond, perforated-lamellar). Recently, Glotzer and co-workers examined the phase behavior of amphiphiles consisting of nanoparticles of various shapes with a finite number and position of surface-tethered polymers using Brownian dynamics of coarse grain particle and linear bead-spring chains (FENE potential).<sup>140</sup> Figure 13 summarizes some of the intriguing possibilities for symmetrically functionalized building blocks. Currently, the synthetic control to make these “building blocks” with sufficient precision is still being developed. However, significant strides toward the controlled synthesis<sup>141–143</sup> and purification<sup>144</sup> of multivalent nanoparticles, which would enable spatially and numerically defined coupling of polymer chains to the nanoparticle, are continually being reported. For example, Banin and colleagues<sup>141</sup> have demonstrated the single-phase synthesis of hybrid metal–semiconductor nanoparticles, such as symmetric two-sided Au-tipped CdSe nanorods (nanodumbbells (NDBs)) or asymmetric one-sided Au-tipped rods. Dumbbell-like Au–Fe<sub>3</sub>O<sub>4</sub> Janus nanoparticles were synthesized using decomposition of Fe(CO)<sub>5</sub> on the surface of the Au nanoparticles followed by oxidation in 1-octadecene solvent.<sup>142</sup> Giannelis and co-workers have begun developing synthetic approaches to create monolithic nanoparticle liquids through judicious selection of the organic corona.<sup>145</sup> Roan demonstrated that end-grafted immiscible homopolymers can confer multivalence to nanoparticles, resulting in soft nanopolyhedra with structures identical to those found in small clusters of colloidal microspheres.<sup>146</sup> Krishnamoorti and Matyjaszewski have begun exploring the phase behavior of

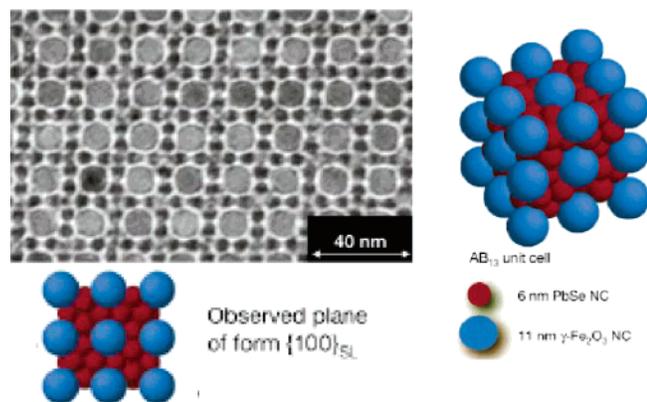


**Figure 14.** Transmission electron micrograph of a two-dimensional assembly of silica nanoparticles with surface-grafted poly(methyl methacrylate)-poly(*n*-butyl acrylate) (PMMA-PBA) block copolymer (50:50 PBA: PMMA, MW ~ 80 K, inner block of PBA).<sup>147</sup> Scale bar is 100 nm.

nanoparticles with macromolecular corona’s consisting of surface-grafted, strongly segregating diblock copolymers, Figure 14.<sup>147</sup> In parallel with the synthetic tools, others, such as Landman and Luedtke, are beginning to explore the implications of the hard–soft periodicity of these assemblies on their thermal and mechanical properties.<sup>148</sup>

**Nanoparticle–Nanoparticle Systems.** Conceptually following blends of nanoparticles with linear chains, combinations of different structured nanoparticles will provide additional routes to even more complex morphologies reminiscent of atomic crystal structures. The possibilities are exemplified by prior work on colloidal crystals and nanoparticle superstructures.

It is noteworthy that, even for a simple example of binary mixtures of hard spheres, complex superlattice structures are predicted to occur for specific size ratios and have been observed for classic micrometer-scale colloids.<sup>149–152</sup> Equilibrium phases with stoichiometry of AB<sub>2</sub> and AB<sub>13</sub> and with upward of 112 spheres per unit cell can form due solely to entropic considerations, that is, due to maximization of the nanoparticle packing density. Murray and co-workers have recently demonstrated comparable structures for binary nanoparticle blends, Figure 15.<sup>153,154</sup> These recent efforts have also indicated that not only does electrical charges on sterically stabilized nanoparticles determine the superlattice

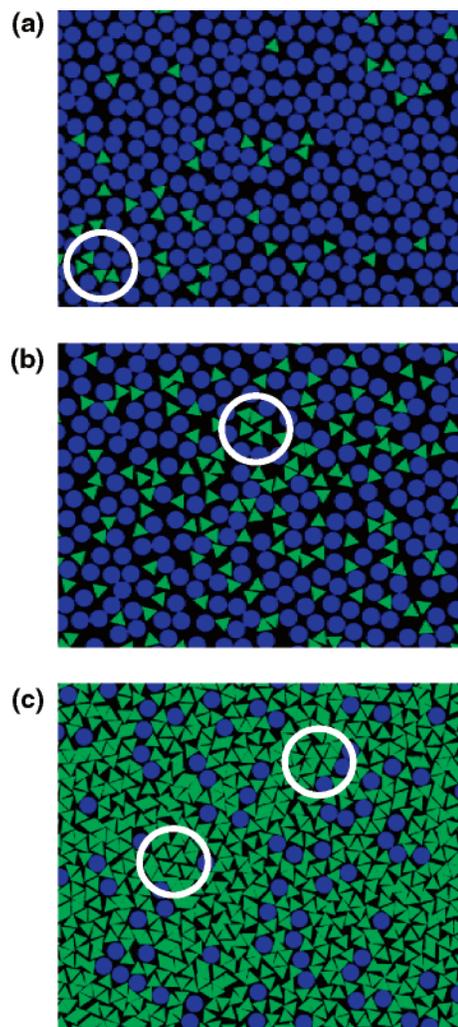


**Figure 15.** Transmission electron micrograph and sketches of  $AB_{13}$  superlattices (isostructural with intermetallic phase  $NaZn_{13}$ ; space group, 226) of 11 nm  $\gamma$ - $Fe_2O_3$  and 6 nm PbSe nanocrystals from Murray et al.<sup>153</sup>  $AB_{13}$  unit cell depicted is built up of eight cubic subunits.

stoichiometry; additional contributions from van der Waals, steric, and dipolar forces also influence the final order and thus can lead to an even greater array of superlattices.

Not all potential structured nanoparticles are spherical however. Nonspherical interparticle potential can easily be envisioned for multivalent nanoparticles derived from surface-tethered polymers.<sup>51,146</sup> As an example, in a series of elegant experiments, Whitesides et al.<sup>155,156</sup> have demonstrated clearly the importance of geometrical considerations in the development of long-range structure and have suggested how such structures might be used to engineer nanomaterials. Fraden and co-workers have extensively examined the phase behavior of various shaped viruses including rods with proteins and polymers.<sup>157,158</sup> Recently de Wild et al.<sup>159</sup> have used molecules with triangular symmetry (sub-phthalocyanine) to produce a number of long-range structures on silver surfaces. Recent theoretical and modeling of two-dimensional angular particulate systems<sup>160–164</sup> has demonstrated the richness of phase behavior that can occur even in the most simple systems. As previously noted, in all of this work an important question arises, namely, to what degree are the observed phases a consequence of the purely topological and space filling attributes of the particle and to what degree might they be a function of chemical interaction?<sup>165</sup>

The process of self-assembly of these complex structures can generally be divided into two broad categories: static and dynamic.<sup>155</sup> Static self-assembly is said to occur when the components are in a global or quasi-equilibrium state. The examples discussed above are representative of static self-assembly. While an energy burden may occur to enter the static state, energy dissipation is not required to maintain this state. Static equilibrium is most often seen in crystalline assemblies of proteins, colloids, or mesoparticles. In contrast, dynamic self-assembly requires a continual infusion of energy to maintain its quasi-equilibrium state. For instance cellular mitosis<sup>166</sup> and circadian rhythms<sup>167</sup> may be viewed as a form of dynamic self-assembly. In this regard it is intriguing to note that “simple” diffusion, which is governed by a nonlinear differential equation, can give rise to complex patterns when two components with differing diffusion rates exist within the same system.<sup>168</sup> These concepts manifest in Turing patterns<sup>169</sup> and Belousov–Zhabotinsky (BZ) reactions



**Figure 16.** Molecular dynamic simulations of mixtures of hard smooth impenetrable triangles and circles in two dimensions.<sup>175</sup> Even when the mole fraction of the triangles is low (a,  $\chi = 0.15$ ), it is noteworthy that the triangles are not randomly distributed on the local scale. As the concentration increases (b,  $\chi = 0.375$ ) into the transition region, this tendency to local clustering increases (example circled). At the highest concentration (c,  $\chi = 0.80$ ) the triangular fluid forms a liquid quasi-crystalline phase (analogous to Penrose tiling, examples circled) that imposes a nonrandom local order on the minority circular particles.

and are thought to be the root of many structures and singling functions in biology.<sup>170</sup>

Overall, there has been comparatively little theoretical and modeling work on the self-assembly dynamics of hard space-filling bodies of angular, nonspherical geometry. These dynamics calculations are exceedingly intensive in terms of computer time, and it is perhaps for this reason that such studies have only recently been extended to hard space-filling bodies of more complex shape.<sup>164</sup> Frenkel and Maguire<sup>171</sup> and Magda et al.<sup>172</sup> have reported the transport properties of a fluid of infinitely thin hard-line segments. More recently Huthman et al. has used their approach to simulate granular cooling of hard needles,<sup>173</sup> while Yoshimura and Mukoyama<sup>174</sup> have made detailed studies of isolated binary chattering collisions between rods. For example, the hard triangle system undergoes what appears to be a second-order phase transition at a packing fraction of  $\sim 0.67$  and exhibits a clear tendency to form long-lived hexagonal clusters. Similarly in mixtures of circles and triangles, Figure 16,<sup>175</sup> there is

evidence of clustering of triangular particles at low mole fraction ( $\sim 0.1$ ) that increase as the system approaches the transition region (0.6).

### Conclusion

Efforts to date on PNCs with random dispersions of nanoparticles (nanoparticles: spheres, rods, or plates) have proven that overcoming traditionally antagonistic combinations of properties, while maintaining the manufacturing and processing flexibility inherent to polymers and resins, is possible. For the future, the explosion of functional NPs promises to expand these possibilities and enable plastics with never-before-realized property suites. Our understanding of the property–morphology correlations in these systems are arguably still in its infancy, but is progressing rapidly, identifying the relative contributions of interface properties, NP distribution, NP–NP correlations, and NP–matrix interactions to the final property suite. These investigations are showing that the chemistry and physics of PNCs parallel many macromolecular and colloidal systems, not just that of filled polymers and polymer blends. Thus, when compared to current PNCs, the richness of structures, morphologies, and properties exemplified by these related macromolecular and colloidal areas underscores the nascent state and untapped possibilities awaiting future development of PNCs with tailored morphology. The addition of increased morphology control and its associated property implications, including anisotropy in physical characteristics on the macroscale, will broaden the potential impact of PNCs beyond current commodity applications, such as automotive, packaging, and coatings. Possibilities include active components of devices, such as MEMs, chemical–biological sensors, photovoltaics, and ultrahigh-energy density capacitors, as well as being the enabling material for numerous high-technology applications, including reconfigurable antennas, thermal control materials, lightning strike, and electrostatic discharge protection systems and passive optical filters and coatings.

Of the numerous challenges to achieve these goals, three stand out with respect to their fundamental impact and pervasiveness: phase behavior, constitutive relationships, and NP interface modification.

Theory, modeling, and experimental validation of the phase behavior of nanoparticle-complex fluid systems underpins all concepts to add order to the random assembly. Quantification of NP–matrix and NP–NP interactions and its impact on structure drives selection of NP modification. Differentiation between thermodynamically dispersed and aggregated states with respect to extrinsic and intrinsic variables will ultimately determine the long-term stability, manufacturability, and the necessity for additional concepts to “trap” potentially desired metastable states. Current investigations noted above are beginning to lay down the framework for spherical NPs. However, the complexities and required increase of precision necessitated by NP anisotropy (rods and plates) and multiple intermolecular and interparticle interactions of comparable magnitude but opposite polarity (i.e., attractive and repulsive) minimizes the direct applicability (and some argue relevance) of current constructs to

systems of commercial interest. Bridging the diverse spatial and temporal scales embodied in these nanoscale systems is a major future challenge.

The second major challenge is the development of structure–property constitutive relationships to establish reasonable bonds on the expected property improvements. Without these correlations, exploration of new structures is based on intuition. More importantly, though, these correlations are paramount to guide processing design and enable the system engineer to evaluate trade-offs and benefits afforded by the new material offerings. The majority of current efforts on random PNCs are based on applying previously validated constitutive relationships, for example Halpin–Tsai or Mori–Tanaka for mechanical behavior–morphology correlations. The success of these relationships in describing property trends is not consistent across the spectrum of PNCs. Quantitative investigation of why they are successful or not, especially when the PNC under consideration invalidates many of the base assumptions underlying the constitutive relationship, is hampered by inadequate systematic control of the NP dispersion, morphology, and interface. The refinement of processing approaches such as those discussed above will have a profound impact on this critical challenge.

Finally, realization of any morphology, whether random or ordered, is predicated on the ability to disperse and manipulate the nanoparticle, whether a sphere, rod, or plate. Previous efforts have definitely shown that a single universal dispersion approach is not realistic. A suite of robust, efficient, scalable, and cost-effective methodologies that are tailorable to various systems under consideration is necessary. Innovations to this end can be found throughout the literature, ranging from noncovalent to covalent; small molecular to macromolecular; chemical to mechanical; and single-phase to multiple-phase/emulsions. The current struggle is predicated on the reality of the uniqueness of a concept to a particular NP–polymer system, and the requirement to refine and formulate, especially in light of scale-up, cost, and processability criteria. For example, if an interfacial modifier is necessary and its role is simply to enable dispersion, cost considerations imply a minimal amount should be used, but this is at odds with the more than 200 m<sup>2</sup>/g surface area of NPs and limits the applicability of many macromolecular concepts. Furthermore, as NP size and polydispersity decrease, they form ordered low-dimensional crystals and assemblies with substantial fraction of direct NP–NP surface interactions relative to a disordered aggregate or agglomerate. These interactions stabilize the NP assembly, lowering its free energy and increasing the challenge of disrupting the structure. These systems are especially challenging to disperse as exemplified by single-wall carbon nanotubes (SWNTs), which arrange in the form of ropes in contrast to multiwall carbon nanotubes (MWNTs) or carbon nanofibers (CNFs) which for low-density aggregates are of individual tubes. Beyond dispersion, however, is the objective of incorporating functionality directly at the interface to further utilize the extensive interfacial area afforded by the NP. For example, incorporating field (e.g., light (RF–UV–vis), magnetic, or electrical)—triggered conformation changes or

thermally responsive macromolecules on the NP surface would provide PNCs with switchable property suites, i.e., tunable, responsive and adaptive materials. The utilization of such systems would range from smart tires whose damping characteristics change on the basis of relative humidity to control-release tissue scaffolds. Married with increased morphology control and thus the ability to tune a property in one-direction would provide even greater possibilities, such as smart membranes, filters, and morphing structures.

In conclusion, the development of a plurality of approaches to control nanoparticle order will inevitably provide solutions to future (and current) manufacturing needs. However, it is difficult to imagine large-scale production of PNC components based on many of the discussed techniques. Ultimately, complete hierarchical control of the morphology will not rely on a single approach but, as with current manufacturing technologies, will depend on the separation of fine- and coarse-scale control to integrate function and form. Highly ordered nanostructures in the form of sub-micrometer coatings, fibers, and films will be integrated sequentially into larger components, such as commonly done in multilamellae film fabrication for food and beverage packaging, in high-performance optical coatings for windows and lenses, and in fiber-weaving of multifunctional composite structures, just to name a few. Thus, true “polymer nanocomposites” when integrated into material systems could epitomize the “performance enabled by nanoinside” mantra.

**Acknowledgment.** The authors are grateful to the Air Force Office of Scientific Research and the Air Force Research Laboratory’s Materials and Manufacturing Director for financial support. We would also like to extend our thanks to T. Bunning, Y. Cohen, R. Krishnamoorti, H. Koerner, and L. Drummy for insightful discussions and for providing figures. Finally, R.A.V. is indebted to the Materials Research Laboratory at the University of California, Santa Barbara, as well as E. Kramer, G. Fredrickson, and C. Hawker for their outstanding hospitality during the development of this manuscript.

## References

- Winey, K., Vaia, R., Eds. *Polymer Nanocomposites*; MRS Bulletin; Materials Research Society: Pittsburgh, PA, 2007.
- McWilliams, A. Nanocomposites, Nanoparticles, Nanoclays, and Nanotubes (NANO21C), BCC Research, Norwalk, CT, 2006.
- Polymer Nanocomposites Create Exciting Opportunities in the Plastics Industry, Principia Partners, Jersey City, NJ, 2005.
- Thayer, A. M. Firms Find a New Field of Dreams. *Chemical & Engineering News*; American Chemical Society: Washington, DC, 2000 (Oct 16).
- Fukushima, Y.; Inagaki, S. *J. Inclusion Phenom.* **1987**, *5*, 473–482.
- Fukushima, Y.; Okada, A.; Kawasumi, M.; Kurauchi, T.; Kamigaito, O. *Clay Miner.* **1988**, *23*, 27.
- Usuki, A.; Kawasumi, M.; Kojima, Y.; Okada, A.; Kurauchi, T.; Kamigaito, O. *J. Mater. Res.* **1993**, *8*, 1174.
- Kojima, Y.; Usuki, A.; Kawasumi, M.; Okada, A.; Fukushima, Y.; Kurauchi, T.; Kamigaito, O. *J. Mater. Res.* **1993**, *8*, 1185.
- Vaia, R. A.; Ishii, H.; Giannelis, E. P. *Chem. Mater.* **1993**, *5*, 1694.
- Lan, T.; Pinnavaia, T. J. *Chem. Mater.* **1994**, *6*, 2216.
- Giannelis, E. P. *Adv. Mater.* **1996**, *8* (1), 29–35.
- Alexandre, M.; Dubois, P. *Mater. Sci. Eng. Rep.* **2000**, *R28*, 1–65.
- Ray, S. S.; Okamoto, M. *Prog. Polym. Sci.* **2003**, *2*, 1539.
- Okamoto, M. Polymer/Clay Nanocomposites. In *Encyclopedia of Nanoscience and Nanotechnology*; Nalwa, H. S., Ed.; American Scientific: Stevenson Ranch, CA, 2004; Vol. 8, p 1.
- Thostenson, E., T.; Li, C.; Chou, T.-W. *Compos. Sci. Technol.* **2005**, *65*, 491–516.
- Drummy, L. F.; Koerner, H.; Farmer, B. L.; Vaia, R. A. *Advanced Morphology Characterization of Clay-Based Polymer Nanocomposites*; CMS Workshop Lecture Series, Vol. 14; Clay Minerals Society: 2006.
- Tjong, S. C. *Mater. Sci. Eng. R* **2006**, *53*, 73.
- Hussain, F.; Hojjati, M.; Okamoto, M.; Gorga, R. E. *J. Compos. Mater.* **2006**, *40* (17), 1511.
- Xie, X.-L.; Maia, Y.-W.; Zhou, X.-P. *Mater. Sci. Eng. R* **2005**, *49*, 89–112.
- Moniruzzaman, M.; Winey, K. I. *Macromolecules* **2006**, *39*, 5194.
- Pinnavaia, T. J.; Beall, G. W. *J. Polymer-Clay Nanocomposites*; John Wiley & Sons: New York 2001.
- Krishnamoorti, R.; Vaia, R. A., Eds. *Polymer Nanocomposites: Synthesis, Characterization, and Modeling*; ACS Symposium Series; American Chemical Society: Washington, DC, 2001 (Dec 6).
- Ray, S. S.; Bousmina, M. *Polymer Nanocomposites and Their Applications*; American Scientific: Stevenson Ranch, CA, 2006 (Oct 31).
- Mai, Y.-W.; Yu, Z.-Z., Eds. *Polymer Nanocomposites*; CRC: Boca Raton, FL, 2006 (Mar 15).
- Morgan, A. B., Wilkie, C. A., Eds. *Flame Retardant Polymer Nanocomposites*; John Wiley & Sons: New York, in press.
- Drummy, L.; Naik, R. Private communication.
- Koerner, H.; Vaia, R. A.; Liu, W.; Alexander, M.; Mirau, P. *Polymer* **2005**, *46*, 4405.
- Koerner, H.; Misra, D.; Tan, A.; Drummy, L.; Mirau, P.; Vaia, R. *Polymer* **2006**, *47*, 3426.
- Vaia, R. A.; Giannelis, E. P. *MRS Bull.* **2001**, *26* (5), 394.
- Vaia, R. A.; Wagner, H. D. *Mater. Today* **2004**, *32*.
- Bansal, A.; Yang, H.; Li, C.; Cho, K.; Benicewicz, B. C.; Kumar, S. K.; Schadler, L. S. *Nat. Mater.* **2005**, *4* (9), 693.
- Krishnamoorti, R.; Vaia, R. A.; Giannelis, E. P. *Chem. Mater.* **1996**, *9*, 1728.
- Starr, F. W.; Schroeder, T. B.; Glotzer, S. C. *Macromolecules* **2002**, *35*, 4481.
- Jones, R. A. L.; Richards, R. W. *Polymers at Surfaces and Interfaces*; Cambridge University Press: Cambridge, U.K., 1999.
- Priestley, R. D.; Ellison, C. J.; Broadbelt, L. J.; Torkelson, J. M. *Science* **2005**, *309*, 456.
- Alcoutlabi, M.; McKenna, G. B. *J. Phys. Condens. Matter* **2005**, *17*, R467.
- Kojima, Y.; Usuki, A.; Kawasumi, M.; Okada, A.; Kurauchi, T.; Kamigaito, O.; Kaji, K. *J. Polym. Sci., Part B: Polym. Phys.* **1994**, *32*, 625.
- Kojima, Y.; Usuki, A.; Kawasumi, M.; Okada, A.; Kurauchi, T.; Kamigaito, O.; Kaji, K. *J. Polym. Sci., Part B: Polym. Phys.* **1995**, *33*, 1039.
- Rice, B. P.; Chen, C.; Cloos, L.; Curliss, D. *SAMPE J.* **2001**, *37*, 7.
- Gusev, A. A.; Rozman, M. G. *Comput. Theor. Polym. Sci.* **1999**, *9*, 335.
- Gusev, A. A.; Lusti, H. R. *Adv. Mater.* **2001**, *13*, 1641.
- Gusev, A. A.; Lusti, H. R.; Hine, P. J. *Adv. Eng. Mater.* **2002**, *4* (12), 927.
- Sheng, N.; Boyce, M. C.; Parks, D. M.; Rutledge, G. C.; Abes, J. I.; Cohen, R. E. *Polymer* **2004**, *45*, 487.
- Buxton, G. A.; Balazs, A. C. *J. Chem. Phys.* **2002**, *117* (16), 7649.
- Liu, T.; Kumar, S. *Nano Lett.* **2003**, *5*, 647.
- Zheng, X.; Forest, M. G.; Vaia, R.; Arlen, M.; Zhou, R. Shear-guided anisotropic geometrical percolation in nano-rod ensembles. *Adv. Mater.* **2007**, in press.
- Masala, O.; Seshadri, R. A. *Rev. Mater. Res.* **2004**, *34*, 41.
- Cushing, B. L.; Kolesnichenko, V. L.; O’Connor, C. J. *Nanoparticles Chem. Rev.* **2004**, *104*, 3893.
- Zheng, N.; Fan, J.; Stucky, G. D. *J. Am. Chem. Soc.* **2006**, *128*, 6550.
- Jun, Y.; Seo, J.; Oh, S. J.; Cheon, J. *Coord. Chem. Rev.* **2005**, *249*, 1766.
- Naik, R. R.; Jones, S. E.; Murray, C. J.; McAuliffe, J. C.; Vaia, R. A. M. O. *Adv. Funct. Mater.* **2004**, *14* (1), 25–30.
- Murphy, C. J.; Sau, T. K.; Gole, A. M.; Orendorff, C. J.; Gao, J.; Gou, L.; Honyadi, S. E.; Li, T. *J. Phys. Chem. B* **2005**, *109*, 13857.
- Grillet, A.-C.; Brunel, S.; Chevalier, Y.; Usoni, S.; Ansanay-Alex, V.; Allemand, J. *Polym. Int.* **2004**, *53*, 569.
- Fong, H.; Vaia, R. A.; Sanders, J. H.; Lincoln, D.; Vreugdenhil, A. J.; Liu, W.; Bultman, J.; Chen, C. *Chem. Mater.* **2001**, *13*, 4123.
- Fong, H.; Liu, W.; Wang, C.-S.; Vaia, R. A. *Polymer* **2002**, *43*, 775.
- Pavlikova, S.; Thomann, R.; Reichert, P.; Mulhaupt, R.; Marcincin, A.; Borsig, E. *J. Appl. Polym. Sci.* **2003**, *89*, 604.
- Vigloo, B.; Pénicaud, A.; Coulon, C.; Sauder, C.; Pailler, R.; Journet, C.; Bernier, P.; Poulin, P. *Science* **2003**, *290*, 1331.
- Kumar, S.; Dang, T. D.; Arnold, F. E.; Bhattacharyya, A. R.; Min, B. G.; Zhang, Xiefei Vaia, R. A.; Park, C.; Adams, W. W.; Hauge, R. H.; Smalley, R. E.; Ramesh, S.; Willis, P. A. *Macromolecules* **2002**, *35*, 9039.
- Dirama, T. E.; Goettler, L. A. *Mater. Manuf. Processes* **2006**, *21*, 199.
- Cho, J. W.; Paul, D. R. *Polymer* **2001**, *42*, 1083.
- Krishnamoorti, R.; Yurekli, K. *Curr. Opin. Colloid Interface Sci.* **2001**, *6*, 464.
- Schmidt, G.; Nakatani, A. I.; Butler, P. D.; Karim, A.; Han, C. C. *Macromolecules* **2000**, *33*, 7219.
- Koerner, H.; Hampton, E.; Dean, D.; Turgut, Z.; Drummy, L.; Mirau, P.; Vaia, R. *Chem. Mater.* **2005**, *17*, 1990.
- Casavant, M. J.; Walters, D. A.; Schmidt, J. J.; Smalley, R. E. *J. Appl. Phys.* **2003**, *93* (4), 2153.
- Koerner, H.; Jacobs, J. D.; Tomlin, D. W.; Busbee, J.; Vaia, R. *Adv. Mater.* **2004**, *16*, 297.
- Lvov, Y.; Ariga, K.; Ichinose, I.; Kunitake, T. *Langmuir* **1996**, *12*, 3038.
- Olek, M.; Ostrander, J.; Jurga, S.; Möhwald, H.; Kotov, N.; Kempa, K.; Giersig, M. *Nano Lett.* **2004**, *4*, 1889.

- (68) Tucker, C. L., III; Moldenaers, P. *Annu. Rev. Fluid Mech.* **2002**, *34*, 177.
- (69) Park, S.-Y.; Cho, Y.-H.; Vaia, R. A. *Macromolecules* **2005**, *38*, 1729.
- (70) Cohen, Y.; Albalak, R. J.; Dair, B. J.; Capel, M. S.; Thomas, E. L. *Macromolecules* **2000**, *33*, 6502.
- (71) Qiao, L.; Winey, K. I. *Macromolecules* **2000**, *33*, 851.
- (72) Fairclough, J. P. A.; Salou, C. L. O.; Ryan, A. J.; Hamley, I. W.; Daniel, C.; Helsenby, W. I.; Hall, C.; Lewis, R. A.; Gleeson, A. J.; Diakun, G. P.; Mant, G. R. *Polymer* **2000**, *41* (7), 2577.
- (73) O'Brien, V. T.; Mackay, M. E. *J. Rheol.* **2002**, *46* (3), 557.
- (74) Halsey, T. C. *Science* **1992**, *258*, 761.
- (75) Goncalves, F. D.; Koo, J.-H.; Mehdi, A. *Shock Vib. Dig.* **2006**, *38* (3), 203.
- (76) Westermeier, R. *Electrophoresis in Practice: A Guide to Methods and Applications of DNA and Protein Separations*, 4th ed.; John Wiley & Sons: 2005 (Jan 5).
- (77) Hughes, M. P. *Nanotechnology* **2000**, *11*, 124.
- (78) Morariu, M. D.; Voicu, N. E.; Schäffer, E.; Lin, Z.; Russell, T. P.; Steiner, U. *Nat. Mater.* **2003**, *2*, 48.
- (79) Kim, D.; Lu, W. *Phys. Rev. B* **2006**, *73*, 035206.
- (80) Gascoyne, P. R. C.; Vykoukal, J. *Electrophoresis* **2002**, *23*, 1973.
- (81) Green, N. G.; Ramos, A.; Morgan, H. *J. Phys. D: Appl. Phys.* **2000**, *33*, 632.
- (82) Pal, B. N.; Basu, S.; Chakravorty, D. *J. Appl. Phys.* **2005**, *98*, 084306.
- (83) Casavant, M. J.; Walters, D. A.; Schmidt, J. J.; Smalley, R. E. *J. Appl. Phys.* **2003**, *93*, 2153.
- (84) Park, C.; Wilkinson, J.; Banda, S.; Ounaies, Z.; Wise, K.; Sauti, G.; Lillehei, P. T.; Harrison, J. *J. Polym. Sci., Part B: Polym. Phys.* **2006**, *44*, 1751.
- (85) Koerner, H.; Shiota, A.; Ober, C. K.; Bunning, T. *Science* **1996**, *272*, 252.
- (86) Cao, Y.; Irwin, P. C.; Younsi, K. *IEEE Trans. Dielectr. Electr. Insul.* **2004**, *11* (5), 797.
- (87) Roy, M.; Nelson, J. K.; MacCrone, R. K.; Schadler, L. S.; Rees, C. W.; Keefe, R.; Zenger, W. *IEEE Trans. Dielectr. Electr. Insul.* **2005**, *12* (4), 629.
- (88) Saha, S. K. *Phys. Rev. B* **2004**, *69*, 125416.
- (89) Saha, S. K.; Chakravorty, D. *Appl. Phys. Lett.* **2006**, *89*, 043117.
- (90) Vaia, R. A.; Dennis, C. L.; Natarajan, L. V.; Tondiglia, V. P.; Tomlin, D. W.; Bunning, T. J. *Adv. Mater.* **2001**, *13*, 1570.
- (91) Sanchez, C.; Escoto, M. J.; van Heesch, C.; Basitaansen, C. W. M.; Broer, D. J.; Loos, J.; Nossbaumer, R. *Adv. Funct. Mater.* **2005**, *15*, 1623.
- (92) Suzuki, N.; Tomita, Y.; Kojima, T. *Appl. Phys. Lett.* **2002**, *81*, 4121.
- (93) Jakubiak, R.; Tondiglia, V. P.; Natarajan, L. V.; Sutherland, R. L.; Lloyd, P.; Bunning, T. J.; Vaia, R. A. *Adv. Mater.* **2005**, *17* (23), 2807.
- (94) Tondiglia, V. P.; Natarajan, L. V.; Sutherland, R. L.; Tomlin, D.; Bunning, T. J. *Adv. Mater.* **2002**, *14* (3), 187.
- (95) Jakubiak, R.; Brown, D. P.; Vatansever, F.; Tondiglia, V. P.; Natarajan, L. V.; Tomlin, D. W.; Bunning, T. J.; Vaia, R. A. Holographic Photopolymerization for Fabrication of Electrically Switchable Inorganic-Organic Hybrid Photonic Structures. In *Proceedings of SPIE—The International Society for Optical Engineering*, Vol. 4991; Organic Photonic Materials and Devices V; Grote, J. G., Kaino, T., Eds.; SPIE: Bellingham, WA, 2003; pp 89–97.
- (96) Jakubiak, R.; Tondiglia, V. P.; Natarajan, L. V.; Sutherland, R. L.; Lloyd, P.; Bunning, T. J.; Vaia, R. A. *Adv. Mater.* **2005**, *17* (23), 2807.
- (97) Lipatov, Y. S. *Prog. Polym. Sci.* **2002**, *27*, 1721.
- (98) Hong, J. S.; Namkung, H.; Ahn, K. H.; Lee, S. J.; Kim, C. *Polymer* **2006**, *47*, 3967.
- (99) Gelfer, M. Y.; Song, H. H.; Liu, L.; Hsiao, B. S.; Chu, B.; Rafailovich, M. *J. Polym. Sci., Part B: Polym. Phys.* **2003**, *41*, 44.
- (100) Ray, S. S.; Pouliot, S.; Bousmina, M.; Utracki, L. A. *Polymer* **2004**, *45*, 8403.
- (101) Wolfgang, H. *Angew. Chem., Int. Ed.* **2005**, *44*, 5172.
- (102) Binks, B. P.; Clint, J. H. *Langmuir* **2002**, *18*, 1270.
- (103) Lin, Y.; Skaff, H.; Ermick, T.; Dinsmore, A. D.; Russell, T. P. *Science* **2003**, *299*, 226.
- (104) Duan, H.; Wang, D.; Kurth, D. G.; Mohwald, H. *Angew. Chem.* **2004**, *116*, 5757.
- (105) Duan, H.; Wang, D.; Sobal, N. S.; Giersig, M.; Kurth, D. G.; Molhwal, H. *Nano Lett.* **2005**, *5*, 949.
- (106) Ashby, N. P.; Binks, B. P. *Phys. Chem. Chem. Phys.* **2000**, *2*, 5640.
- (107) Lee, K. Y.; Kim, M.; Hahn, J.; Suh, J. S.; Lee, I.; Kim, K.; Han, S. W. *Langmuir* **2006**, *22*, 1817.
- (108) Voorn, D. J.; Ming, W.; van Herk, A. M. *Macromolecules* **2006**, *39*, 2137.
- (109) Zubarev, E. R.; Xu, J.; Sayyad, A.; Gibson, J. D. *J. Am. Chem. Soc.* **2006**, *128*, 4958.
- (110) Glogowski, E.; He, J.; Russell, T. P.; Emrick, T. *Chem. Commun. (Cambridge)* **2005**, 4050.
- (111) Zubarev, E.; Xu, J.; Sayyad, A.; Gibson, J. P. *J. Am. Chem. Soc.* **2006**, *128*, 15098.
- (112) Tsitsilianis, C.; Voulgaris, D.; Stepanek, M.; Podhajecka, K.; Prochazka, K.; Tuzar, Z.; Brown, W. *Langmuir* **2000**, *16*, 6868.
- (113) Hadjichristidis, N.; Pispas, S.; Pitsikalis, M.; Iatrou, H.; Vlahos, C. *Adv. Polym. Sci.* **1999**, *142*, 71.
- (114) Barker, J.; Henderson, D. *J. Chem. Phys.* **1967**, *47* (8), 2856.
- (115) Gray, C.; Gubbins, K. *Theory of Molecular Fluids*, Vol. 1; Oxford University Press: Oxford, U.K., 1984; Chapter 4.
- (116) Rowlinson, J.; Swinton, F. *Liquids and Liquid Mixtures*, 3rd ed.; Butterworth: London, 1982.
- (117) Wang, Q. *J. Chem. Phys.* **2002**, *116*, 9120.
- (118) Flory, P. J.; Abe, A. *Macromolecules* **1978**, *11*, 1119.
- (119) Onsager, L. *Ann. N. Y. Acad. Sci.* **1949**, *51*, 627.
- (120) de Gennes, P. G.; Proust, J. *The Physics of Liquid Crystals*, 2nd ed.; Oxford University Press: Oxford, U.K., 1993.
- (121) Hooper, J. B.; Schweizer, K. S. *Macromolecules* **2006**, *39*, 5133.
- (122) Chen, Y.-L.; Schweizer, K. S. *J. Phys. Chem. B* **2004**, *108*, 6687.
- (123) Zhao, L.; Li, Y.-G.; Zhong, C.; Mi, J. *J. Chem. Phys.* **2006**, *124*, 144913.
- (124) Mackay, M. E.; Tuteja, A.; Duxbury, P. M.; Hawker, C. J.; Van Horn, B.; Guan, Z.; Chen, G.; Krishnan, R. S. *Science* **2006**, *311*, 1740.
- (125) Ginzburg, V. V.; Balazs, A. C. *Adv. Mater.* **2000**, *12* (23), 1805.
- (126) Ginzburg, V. V.; Singh, C.; Balazs, A. C. *Macromolecules* **2000**, *33*, 1089.
- (127) Gabriel, J.-C. P.; Davidson, P. *Top. Curr. Chem.* **2003**, *226*, 119.
- (128) Gabriel, J. C. P.; Camerel, F.; Lemaire, B. J.; Desvaux, H.; Davidson, P.; Batail, P. *Nature* **2001**, *413*, 504.
- (129) Davidson, P.; Batail, P.; Gabriel, J. C. P.; Livage, J.; Sanchez, C.; Bourdeaux, C. *Prog. Polym. Sci.* **1997**, *22*, 913.
- (130) Li, L.; Marjanska, M.; Park, G. H. J.; Pines, A.; Alivisatos, P. A. *J. Chem. Phys.* **2004**, *120*, 1149.
- (131) Davis, V. A.; Ericson, L. M.; Nicholas, A.; Parra-Vasquez, G.; Fan, H.; Wang, Y.; Prieto, V.; Longoria, J. A.; Ramesh, S.; Saini, R. K.; Kittrell, C.; Billups, W. E.; Adams, W. W.; Hauge, R. H.; Smalley, R. E.; Pasquali, M. *Macromolecules* **2004**, *37*, 154.
- (132) Kim, B. J.; Chiu, J. J.; Yi, G.-R.; Pine, D. J.; Kramer, E. J. *Adv. Mater.* **2005**, *17*, 2618.
- (133) Lin, Y.; Bocker, A.; He, J.; Sill, K.; Xiang, H.; Abetz, C.; Li, X.; Wang, J.; Emrick, T.; Long, S.; Wang, Q.; Balazs, A.; Russell, T. P. *Nature* **2005**, *434*, 55.
- (134) Ha, Y.-H.; Kwon, Y.; Breiner, T.; Chan, E. P.; Tzianetopoulou, T.; Cohen, R. E.; Boyce, M. C.; Thomas, E. L. *Macromolecules* **2005**, *38*, 5170.
- (135) Kim, B.; Bang, J.; Hawker, C. J. E.; Kramer, J. *Macromolecules* **2006**, *39*, 4108.
- (136) Sides, S. W.; Kim, B. J.; Kramer, E. J.; Fredrickson, G. H. *Phys. Rev. Lett.* **2006**, *96*, 250601.
- (137) Kim, B.; Hawker, C. J.; Fredrickson, G. H.; Kramer, E. J. Manuscript in preparation.
- (138) Matsen, M. W.; Bates, F. S. *J. Chem. Phys.* **1997**, *106*, 2436.
- (139) Donald, A. M.; Windle, A. H.; Hanna, S. *Liquid Crystalline Polymers*, 2nd ed.; Cambridge Solid State Science S; Cambridge University Press: Cambridge, U.K., 2006 (Jun 26).
- (140) Glotzer, S. C.; Horsch, M. A.; Iacovella, C. R.; Zhang, Z.; Chan, E. R.; Zhang, X. *Curr. Opin. Colloid Interface Sci.* **2005**, *10*, 287.
- (141) Mokari, T.; Szturm, C. G.; Salant, A.; Rabani, E.; Banin, U. *Nat. Mater.* **2005**, *4*, 855.
- (142) Yu, H.; Chen, M.; Rice, P. M.; Wang, S. X.; White, R. L.; Sun, S. *Nano Lett.* **2005**, *5*, 379.
- (143) Hurst, S. J.; Payne, E. K.; Qin, L.; Mirkin, C. A. *Angew. Chem., Int. Ed.* **2006**, *45*, 2672.
- (144) Levy, R.; Wang, Z.; Duchesne, L.; Doty, R. C.; Cooper, A. I.; Brust, M.; Fernig, D. G. *ChemBioChem* **2006**, *7*, 592.
- (145) Bourlino, A. B.; Herrera, R.; Chalkias, N.; Jiang, D. D.; Zhang, Q.; Archer, L. A.; Giannelis, E. P. *Adv. Mater.* **2005**, *17*, 234.
- (146) Roan, J.-R. *Phys. Rev. Lett.* **2006**, *96*, 248301.
- (147) Krishnamoorti; Matyjaszewski, K. Private communication, 2006.
- (148) Landman, U.; Luedtke, W. D. *Faraday Discuss.* **2004**, *125*, 1.
- (149) Bartlett, P.; Ottewill, R. R.; Pusey, P. N. *Phys. Rev. Lett.* **1992**, *68*, 3801.
- (150) Trizac, E.; Eldridge, M. D.; Madden, P. A. *Mol. Phys.* **1997**, *90* (4), 675.
- (151) Murray, M. J.; Sanders, J. V. *Philos. Mag. A* **1980**, *42*, 721.
- (152) Eldridge, M. D.; Madden, P. A.; Frenkel, D. *Nature* **1993**, *365*, 35.
- (153) Redl, F. X.; Cho, K.-S.; Murray, C. B.; O'Brien, S. *Nature* **2006**, *439*, 55.
- (154) Shevchenko, E. V.; Talapin, D. V.; Kotov, N. A.; O'Brien, S.; Murray, C. B. *Nature* **2003**, *423*, 968.
- (155) Whitesides, G.; Gryzbowski, B. *Science* **2002**, *295*, 2418.
- (156) Mao, C.; Thalladi, V.; Wolfe, D.; Whitesides, S.; Whitesides, G. *J. Am. Chem. Soc.* **2002**, *124*, 14508.
- (157) Fraden, S. Phase Transitions in Colloidal Suspensions of Virus Particles. In *Observation, Prediction, and Simulation of Phase Transitions in Complex Fluids*; Baus, M., Rull, L. F., Ryckaert, J. P., Eds.; NATO ASI Series C, Vol. 460; Kluwer Academic: Dordrecht, The Netherlands, 1995; pp 113–164.
- (158) Dogic, Z.; Frenkel, D.; Fraden, S. *Phys. Rev. E* **2000**, *62* (3), 3925.
- (159) de Wild, M.; Berner, S.; Sizuki, H.; Yanagi, H.; Schlettwein, D.; Ivan, S.; Baratoff, A.; Güntherodt, H.-J. *ChemPhysChem* **2002**, *10*, 881.
- (160) Wojciechowski, K.; Frenkel, D. *Comput. Methods Sci. Technol.* **2004**, *10* (2), 235.
- (161) Schilling, T.; Pronk, S.; Mulder, B.; Frenkel, D. *Phys. Rev. E* **2005**, *71*, 036138.
- (162) Bates, M.; Frenkel, D. *J. Chem. Phys.* **2000**, *112* (22), 10034.
- (163) Martinez-Raton, Y.; Velasco, E.; Mederos, L. *Phys. Rev. E* **2005**, *72*, 031703.
- (164) Benedict, M.; Maguire, J. F. *Phys. Rev. B* **2004**, *70*, 174112.
- (165) Herzfeld, J. *Acc. Chem. Res.* **1996**, *29*, 33.
- (166) Shapiro, J. *Annu. Rev. Microbiol.* **1998**, *52*, 81.

- (167) Sync, S. *Strogatz: How Order Emerges from Chaos in the Universe, Nature, and Daily Life*; Hyperion: New York, 2003.
- (168) Shinbrot, T.; Muzzio, F. *Nature* **2001**, *410*, 251.
- (169) Yang, L.; Epstein, I. R. *Phys. Rev. Lett.* **2003**, *90* (17), 178303.
- (170) Sick, S.; Reinker, S.; Timmer, J.; Schlake, T. *Science* **2006**, *314*, 1447.
- (171) Frenkel, D.; Maguire, J. *Mol. Phys.* **1983**, *49* (3), 503.
- (172) Magda, J.; Davies, H.; Tirrell, M. *J. Chem. Phys.* 1986, *85*, 6674.
- (173) Aspelmeier, T.; Huthmann, M.; Zeppelius, A. Free Cooling of Particles with Rotational Degrees of Freedom. In *Granular Gases*; Poschel, T., Herding, S., Eds.; Springer: Berlin, 2001.
- (174) Mukoyama, Y.; Yoshimura, A. *J. Phys A: Math. Gen.* **1997**, *30*, 6667.
- (175) Benedict, M.; Elliott, J.; Maguire, J. F. Personal communication.

CM062693+