Further One-Dimensional Analysis of Long-Rod Penetration of Semi-Infinite Targets

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FURTHER ONE-DIMENSIONAL ANALYSIS OF LONG-ROD PENETRATION OF SEMI-INFINITE TARGETS

APPROVED:

Wallace T. Fowler

George W. Botbyl
This work is dedicated to my wife, Tamara. Her love and support were invaluable.
FURTHER ONE-DIMENSIONAL ANALYSIS OF
LONG-ROD PENETRATION OF
SEMI-INFINITE TARGETS

by

JOHN DANIEL CINNAMON, B.S.

REPORT
Presented to the Faculty of the Graduate School of
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July 29, 1992
ABSTRACT

FURTHER ONE-DIMENSIONAL ANALYSIS OF

LONG-ROD PENETRATION OF

SEMI-INFINITE TARGETS

by

JOHN DANIEL CINNAMON, B.S.

SUPERVISING PROFESSOR: Wallace T. Fowler

A new one-dimensional analysis of long-rod penetration of semi-infinite targets is presented. Models in this field attempt to accurately describe the penetration process of a long rod of material into a semi-infinite construct of target material. This one-dimensional analysis predicts profile hole diameters and penetration depths over a wide range of material combinations and extends the previous analysis in this area to include hypervelocity impacts. This approach utilizes values of material dynamic yield strengths, known impact conditions, and the well established crater volume/kinetic energy relationship to predict crater hole characteristics over an impact velocity range of 1 to 6 km/s. The analysis
presented here includes an initial transient phase and modifies previous estimates for pressure on the penetrator tip at steady-state. The average pressure at steady state was found to be a constant value over the range of impact velocities for a particular shot combination. The specific value for the average pressure was found to be a direct function of target strength. The resulting equations retain computational simplicity by remaining completely algebraic in nature. Correlation with a majority of the readily available experimental data is given. The results are quite good for a one-dimensional model.
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List of Symbols and Abbreviations

a  slope of the crater volume/kinetic energy relationship
A  cross-sectional area of the mushroom of the rigid-plastic penetrator
A₀  initial cross-sectional area of the undeformed penetrator
A₁  cross-sectional area of the penetrator at impact
A₁  cross-sectional area of the penetrator at steady state
b  intercept of the crater volume/kinetic energy relationship
D  original diameter of the undeformed penetrator
e  engineering strain in the mushroom of the penetrator
e₀  engineering strain in the mushroom at impact
e₁  engineering strain in the mushroom at steady state
E₀  kinetic energy of the penetrator at impact
T  current undeformed section length
L  original length of the undeformed penetrator
Pₐ  pressure on the axis of the penetrator tip
P₀  pressure on the axis of the penetrator tip at impact
P₁  pressure on the axis of the penetrator tip at steady state
P  average pressure on the penetrator tip
P₀  average pressure on the penetrator tip at impact
P₁  average pressure on the penetrator tip at steady state
Q  average pressure on the penetrator tip at steady state, independent of v₀
r  radial distance from the axis of the penetrator
R  original undeformed penetrator rod radius
R_t  dynamic yield strength of target
u  current penetration velocity
u_0  penetration velocity at impact
v  current velocity of the undeformed section
v_0  impact velocity
V_c  crater volume of the recovered target
Y_p  dynamic yield strength of penetrator
z  penetration depth
\rho  penetrator density
\mu^2  ratio of target density to penetrator density
1.0 INTRODUCTION

In 1967, Tate [1] and simultaneously Alekseevskii [2], published a one-dimensional theory for the penetration of semi-infinite targets by long rods. Tate [3] published a second paper in 1969. These papers form the basis upon which the accepted theory of one-dimensional rod penetration rests. This theory attempts to accurately describe the penetration process of a long rod of material into a semi-infinite construct of target material. A prediction of the resulting penetration hole characteristics is available from this analysis. By comparing the theoretical predictions to experimental data, we can evaluate the accuracy of these one-dimensional models.

In 1987, Jones, et al. [4] modified the equations of motion for the undeformed rod section by employing a balance of linear impulse and momentum. In subsequent papers [5,6] an improvement in the correlation between theoretical and experimental data was observed. However, this approach relied upon post-test measurements for estimates of the engineering strain in the mushroom of the penetrator tip. The strain was assumed to be constant throughout the process and equal to that measured from the profile hole diameter. By using the modified Bernoulli Equation proposed by Tate, the pressure, penetration velocity, and undeformed section velocity were coupled with the modified equation of motion and an equation for the conservation of mass passing from the undeformed section of the penetrator. This system was solved by integration and the resulting prediction of penetration
depths were compared to experimentally observed values from the recovered targets. The penetrator and target dynamic yield strengths were taken to be constant during the event and estimated from laboratory values for yield strengths at the highest available strain-rates.

These results were satisfying and showed promise, but based analysis of key parameters on post-experiment measurement. Kerber, et al. [7], in 1990, utilized a well established crater volume - kinetic energy relationship, e.g. [8], to remove some of the dependence on post-test measurements. The aim was to produce engineering strain as a by-product of the solution of the model. However, because the penetration process was treated as being dominated by the steady-state, the results did not correlate well over a large range of impact velocities and rod lengths. The correlations with experiment were satisfactory for longer rods and higher impact velocities where steady state penetration can be presumed to dominate the event. The initial and terminal transients were neglected in this and previous analysis. The initial transient would appear to become more important in shorter rods and lower impact velocities.

In 1992, Cinnamon, et al. [9,10] reported a significant increase in accuracy by incorporating an initial transient phase to the penetration process. This analysis was based on observations by Ravid, et al. [11], Gillis, et al. [12], and Jones, et al. [13]. This latest approach yielded good accuracy for a one-dimensional model in the 1 to 3 km/s velocity range. The results depended only on physical parameters determined before testing and the well
established crater volume/kinetic energy relationship. The terminal transient was neglected in this model. The resulting equations were completely algebraic in nature. Of some concern in this approach was that the trends in the predicted penetration depth curves were tending toward significantly high values for impact velocities above 3 km/s. When experimental data for 3 - 6 km/s were evaluated, the model's accuracy indeed deteriorated. In examining the model, it became obvious that the predictions of pressure available from the modified Bernoulli Equation were simply too high. In fact, the correlation discovered by Cinnamon, et al. [9,10] indicated that the pressure profile could be successfully modified in such a way that the average pressure at steady-state predicted by the Bernoulli Equation was reduced by a specific factor (which depended exclusively on target strength) to match experimental penetration depths. The parabolic nature in which the modified Bernoulli Equation predicts the interface pressure at steady state as impact velocity increases was somewhat minimized by this technique.

In 1992, Jones, et al. [14], the pressure distribution on the penetrator tip was modified to attempt to correct for the deteriorating results above 2.5 km/s. With a judicious choice of two dimensionless parameters, slight improvement was observed. However, a physical relationship between these parameters and the model was not discovered. The pressures predicted by the modified Bernoulli Equation were once again reduced to achieve acceptable results. However, this new pressure
distribution contained the same parabolic component as the previous approach.

In this paper, several approaches to the choice of pressure at steady state are evaluated. It became clear that past successful solution forms for the pressure distribution were those that cancelled out the parabolic nature of the modified Bernoulli Equation. As a consequence, when the average pressure at steady state is taken to be a particular value for a specific material combination over all impact velocities and directly related to target strength, a significant increase in model accuracy resulted. The model was successfully extended up to 6 km/s. This paper includes a large body of results reflecting analysis of all readily available experimental data in this field of research, a great majority of which has been compiled by Anderson, et al. [15].

2.0 THEORY

The general concepts of the rod penetration process are detailed in Figure A. The undeformed penetrator is a cylinder of known length and diameter which impacts the target at a nominal normal incidence at a known velocity. The penetrator enters the target and experiences mushrooming in the tip. When the event has concluded, a crater with a measurable diameter and depth remains.
(a) shows the undeformed rod of length $L$ and initial cross-sectional area $A_i$. The shaded portion will be lost to erosion. (b) shows the penetration event. $\ell$ is the undeformed section length and $z$ is the penetration depth.

2.1 Primary Governing Equations

In [4], Jones, et al. proposed a modification to the equation of motion of the undeformed section of a rod penetrator. This equation is

$$\ell \ddot{v} + \ell (v-u) = \frac{-P}{\rho(1+e)}$$

where $\ell$ is the undeformed section length, $v$ is the current undeformed section velocity, $u$ is the penetration velocity, $P$ is the average pressure on the penetrator tip, $\rho$ is the penetrator density, and $e$ is the engineering strain in the
penetrator mushroom. The details behind the development of this equation are contained in [4].

Wilson, et al. in [5] added another key equation to the analysis. The conservation of mass across the plastic interface between the mushroom and the undeformed section of the penetrator was given in the form

\[ e \dot{e} = v - u \]  \hspace{1cm} (2)

The penetrator is assumed to be rigid-plastic during the event. In both (1) and (2), dots over the symbols represent differentiation with respect to time. The engineering strain in the mushroom is compressive and therefore negative. The current value of the engineering strain in the mushroom is defined to be

\[ e = \frac{A_i}{A} - 1 \]  \hspace{1cm} (3)

where \( A_i \) is the initial cross-sectional area of the undeformed penetrator and \( A \) is the current cross-sectional area of the mushroom.

2.2 Previous Pressure Analysis

The average pressure on the penetrator tip, \( P \), can be varied by considering various pressure profiles. In [4] and [5], the pressure was assumed to be uniform across the mushroom face. The intensity of the pressure was assumed to be the solution of the modified Bernoulli equation. This
equation, from [1 - 3], is applied at steady state and relates pressure on the axis of the specimen at the penetrator tip, $p_a$, to the undeformed section speed $v$, the penetration velocity $u$, and material properties of the target and penetrator. The modified Bernoulli equation is

$$p_a = \frac{1}{2} \rho v^2 + R_t = \frac{1}{2} \rho (v - u)^2 + Y_p$$

(4)

where $R_t$ and $Y_p$ are dynamic yield strengths of the target and penetrator at suitably high strain rates respectively, and $\mu^2$ is the ratio of the penetrator to target density. The pressure was taken in [4] and [5] not to vary across the mushroom face.

Gillis, et al., in [12] suggested that the uniform pressure profile was not realistic and proposed a parabolic form which was symmetric about the axis of the specimen and zero on the edge of the mushroom. In this case, the pressure $p$ had the form

$$p = p_a \left(1 - \frac{r^2}{R^2}\right)$$

(5)

where $p_a$ is the current pressure on the rod axis, $R$ is the original rod radius, and $r$ is the radial distance from the axis. The factor $(1+e)$ in the denominator of the pressure term in (1) forces the pressure to act over the deformed mushroom face with area $A$, even though (5) refers to the original rod configuration.

The average pressure $P$ can be computed, in general, by
When $P$ was calculated for (5), the result was $\frac{P_a}{2}$. The predictions for penetration depths were somewhat improved over the previously assumed uniform pressure distribution. The results still suffered from ignoring the initial and final transients in the penetration event. This parabolic pressure distribution effectively reduced the parabolic nature of (4) by a factor of two. Although this improved the model by reducing the effect of (4), the addition of at least an initial transient seemed warranted.

In Cinnamon, et al. [9,10], the pressure distribution was generalized to the form

$$p = p_a (1 - \frac{r^2}{R^2})^n$$

which made the average pressure term become

$$P = \frac{P_a}{(n+1)}$$

This new pressure distribution was successfully employed with an initial transient phase to correlate to a large number of experimental cases. The pressure exponent $n$ was found to be a direct function of target strength. Although this approach provided improved results over
previous approaches, the analysis could not be successfully extended beyond 3 km/s.

Jones, et al. [14] attempted to correct the problem by adding a uniform component to the pressure profile. The distribution had the form

\[ p = q + (pa - q)(1 - \frac{r^2}{R^2})^n \]  (9)

and the average pressure became

\[ P = \frac{nq}{(n+1)} + \frac{pa}{(n+1)} \]  (10)

This new pressure profile improved results slightly and extended them into the hypervelocity range (3 - 6 km/s). However, \( q \) and \( n \) were not successfully correlated to any physical parameters.

In this paper, another approach to the determination of the average pressure will be explored and a successful correlation to target strength will be presented. In addition, the widest possible body of available data will be used in the correlation.

2.3 Transient Penetration Analysis

In Cinnamon, et al. [9,10], an initial transient phase of penetration was added to the model. The transient phase is characterized by impact shock effects and complete mushroom growth which occurs between impact and the beginning of steady state penetration. We assume that the
penetrator impacts the target at a known velocity, \( v_0 \), of sufficient magnitude (i.e. \( v_0 > 1 \text{ km/s} \)) such that the undeformed section cannot sustain any appreciable deceleration (i.e. \( \dot{v} = 0 \)) during the initial transient. This means that \( v = v_0 \) throughout the transient or mushrooming phase of penetration. During the initial transient, the mushroom develops from a cross-sectional area \( A_0 \) at impact to \( A_1 \), when steady state begins. The mushroom retains an area of \( A_1 \) throughout the steady state portion of the penetration process until the end of the event. Ravid, et al. [11] reported that there was little change in penetration velocity \( u \) during the shock/impact stage of the initial transient phase. Motivated by this observation, we assumed that the penetration velocity was approximately constant (i.e. \( u = u_0 \)) throughout the initial transient. Hence, (1) becomes

\[
\ddot{e} \ (v_0 - u_0) = \frac{-P}{\rho (1+e)} \quad (11)
\]

and (2) is modified to

\[
\dot{e} \ = \ v_0 - u_0 \quad (12)
\]

These two equations, (11) and (12), govern the mushrooming of the rod during the initial transient phase of the penetration event which precedes the steady state. At impact, the engineering strain in the mushroom is \( e_0 \). When steady state is reached, the strain becomes \( e_1 \).
By eliminating \( \dot{e} \) between (11) and (12) and solving for \( e \), we arrive at an expression for the engineering strain.

\[
e = \frac{-(v_0 - u_0)^2}{(v_0 - u_0)^2 + \frac{p}{\rho}}
\]  

(13)

This relationship provides us with an explicit formula for the strain in the mushroom as it develops during the transient phase. This equation governs the behavior until the beginning of the steady state penetration phase. The pressure on the axis, \( p_a \), is changing rapidly during mushroom formation. It has a large value, \( p_0 \), at impact and a reduced value, \( p_1 \), at steady state.

When steady state is reached, we assume that the modified Bernoulli Equation, (4), is valid. At the transition point between the transient and steady state portions of the event, (4) can be expressed as

\[
p_1 = \frac{1}{2} \mu^2 \rho u_0^2 + R_t = \frac{1}{2} \rho (v_0 - u_0)^2 + Y_p
\]

(14)

and (13) can be written as

\[
e_1 = \frac{-(v_0 - u_0)^2}{(v_0 - u_0)^2 + \frac{p_1}{\rho}}
\]

(15)

where \( p_1 \) is the average pressure on the penetrator tip at the beginning of steady state.
Equation (14) can then be used to solve for \( u_0 \) in terms of the known quantities \( v_0, \rho, \mu^2, R_t, \) and \( Y_p \). The remaining variable in the system of equations is \( P_1 \). The determination of \( P_1 \) is the primary focus of this paper.

The penetration velocity \( u_0 \) can be found algebraically from (14). The three primary cases are outlined below.

For equal penetrator and target dynamic yield strength (i.e. \( R_t = Y_p \)) and equal densities (i.e. \( \mu^2 = 1 \)), \( u_0 \) reduces to

\[
\frac{1}{2} v_0
\]

For unequal penetrator and target dynamic yield strength (i.e. \( R_t \neq Y_p \)) and equal densities (i.e. \( \mu^2 = 1 \)), \( u_0 \) becomes

\[
\frac{\rho v_0^2 + 2(Y_p - R_t)}{2\rho v_0}
\]

In the general case, \( u_0 \) is given by

\[
\frac{-v_0}{\mu^2 - 1} + \frac{1}{\rho(\mu^2 - 1)} \left[ \rho^2 v_0^2 - 2\rho (\mu^2 - 1) (R_t \cdot Y_p \cdot \frac{1}{2} \rho v_0^2) \right]^{1/2}
\]
2.4 Impact Conditions

The model outlined above also provides some information about conditions at impact. The strain on impact $e_0$ can be calculated from (13) if we know the average pressure on the penetrator tip at impact, $P_0$.

$$e_0 = \frac{-(v_0-u_0)^2}{(v_0-u_0)^2 + P_0}$$

(19)

The impact pressure can be estimated from elementary shock physics, using

$$P_0 = \rho u_s u_0$$

(20)

where $u_s$ is the shock speed in the target. Values for $u_s$ as a function of $u_0$ can be found in shock Hugoniot tables, e.g. [16]. Calculation of $P_0$ from a known value of $p_0$ can typically be accomplished using the same approach as the calculation of $P_1$ from $p_a$.

2.5 Cratering Analysis

The mathematical model for the behavior of the penetrator is a rigid-plastic, instantaneously eroding rod model. As a result, the penetrator enters the target with some impact engineering strain $e_0$ which expands to $e_1$ during the transient. The impact pressure $p_0$ is usually very high relative to the steady state pressure $p_1$. Although this pressure decreases rapidly during mushroom
formation in the transient phase, the values for $p_0$ can be significant. The mushroom diameter grows from the time of impact through the transient phase, and ceases at the beginning of the steady state portion of the event. The shock/impact stage takes place in a period of a few microseconds, see [11].

The instantaneous erosion assumption prevents the model from accounting for any additional erosion of the target - which occurs in actual practice. There is typically appreciable change in target geometry due to penetrator and target material ejection from the crater. As a consequence, the recovered targets will appear to have more cylindrical-type craters than the model would predict. Figure B illustrates the crater predicted by the mathematical model, and Figure C indicates how the actual geometry frequently appears.

![Figure B. Idealized Crater Geometry](image)

$A_0$ and $A_1$ are the cross-sectional areas at impact and at steady state respectively. $z$ is the penetration depth.
Figure C. Actual Crater Geometry
$A_1$ is the observed crater cross-sectional area. The crater is assumed cylindrical with altitude $z$.

2.6 Penetration Analysis

Predicting penetration depths from the above model is made possible through the use of a somewhat empirical approach. For a number of years, researchers have observed a significant correlation between crater volume in the recovered targets and impact kinetic energy, e.g. [8]. This relationship appears to be linear for impact cases of sufficiently, but not excessively, high energy and can be expressed in the form

$$V_c = a E_0 + b \quad (21)$$

where $V_c$ is the crater volume, $E_0$ is the impact kinetic energy and is equal to $\frac{1}{2} \rho A_1 L v_0^2$, and the variables $a$ and $b$
are regression constants determined from the available experimental data. The linear fit is performed for each shot combination.

These crater volume/kinetic energy relationships are computed from data points for a particular shot combination. The reliability of the linear fit is, of course, a function of the number of data points available. Since most experimental tests are quite expensive, frequently the data is sparse and/or somewhat scattered. When the linear fit predicts cratering at zero impact kinetic energy, or in some other way reflects erroneous trends, the usefulness of the particular case is significantly reduced. This phenomenon is typically avoided by a sufficient number of experimental points. The accuracy of the penetration prediction is extremely dependent on the crater volume/kinetic energy relationship arrived at using this technique.

By adopting the cylindrical approximation for the crater geometry discussed above, we can generate predictions for the penetration depths. Because of the ejection of material from the crater, the cross-sectional area of the recovered target hole, $A_1$, will be

$$A_1 = \frac{A_i}{1 + e_1} \quad (22)$$

The crater volume can be expressed as

$$V_c = A_1 z \quad (23)$$
where \( z \) is the penetration depth. The penetration depth can be predicted by applying the crater volume/kinetic energy relationship to \( V_c \). From (21) and (23), \( z \) is given by

\[
z = \frac{1}{A_i} (1+e_1) V_c = \frac{1}{A_i} (1 + e_1) (aE_0 + b) \quad (24)
\]

Thus, the penetration depth can be expressed as an algebraic function of known material properties and impact conditions, and as a function of the average pressure on the penetrator tip at steady state.

### 3.0 CURRENT PRESSURE PROFILE ANALYSIS

In the previous work outlined above, it was noted that one of the primary difficulties in achieving good predictions for the crater characteristics was the manner in which the modified Bernoulli Equation (4) predicted pressure as a function of \( u \) and \( v \). The parabolic nature of (4) tended to over-predict penetration depths for the higher velocity cases (3 - 6 km/s). Earlier successes, particularly in the intermediate velocity range (1 - 3 km/s), were the result of pressure distributions that tended to reduce the effect of (4). In this paper, additional pressure distribution analysis is performed to address this problem.
3.1 Previous Pressure Distributions

An obvious departure point for the attempt to improve the one-dimensional analysis and extend it into the hypervelocity range was to begin with the previous pressure distributions. A great body of additional experimental data became available in [15] that expanded the range of materials, greatly increased the number of shot combinations accessible for analysis, and provided data in the hypervelocity range. The pressure profiles detailed in (7) and (9) were examined for possible application in these new cases.

The distribution in (7) was unable to compensate for the parabolic nature of (4) while maintaining accuracy in the intermediate velocity ranges if n was considered a constant as it was in [9,10]. Attempts to model n as a function of material properties or impact conditions were not successful. In general, it was observed that n needed to increase with impact velocity to essentially cancel the effects of the modified Bernoulli Equation. The distribution in (9) followed the same trend. Although the results improved, the net effect was to choose n and q to counter the dramatic pressure increase dictated by equation (4).

3.2 New Pressure Distribution

To attempt to find a solution to this dilemma, a new, more general, pressure distribution was proposed by Jones [17]. This pressure profile takes the form
This profile is a more complex and versatile one. A great deal of additional control over the shape of the pressure distribution was provided by (25). The average pressure then is given by

\[
P = q + (p_a - q) \left( 1 - \left( \frac{r}{R} \right)^m \right) \tag{25}
\]

where \( \Gamma \) is the well known mathematical gamma function.

As a significant number of cases were examined, it became increasingly clear that in order to achieve the desired trends in the theoretical penetration curves, \( n \) and \( m \) were chosen in such a way as to essentially eliminate the effect of the second term in (26). That is, an average pressure comprised of a single value, \( q \), which was unvarying over the range of impact velocities, achieved the best results. This discovery matched our previous experience with (7) and (9). Apparently, the magnitude and trends in the pressures predicted by the modified Bernoulli Equation were not leading to acceptable results. All previous successes were based on choices that reduced or eliminated the contribution of the velocity dependent axial pressure in (4) to the value of \( P_1 \).
3.3 Revised Average Pressure Approach

With the results described above, another approach was required. It became clear that the previous calculation of the average pressure at steady state, $P_1$, was not acceptable. To simplify the analysis, the connection of $P_1$ to a particular pressure distribution is ignored.

The problem is further simplified by the fact that the desired trends in the penetration depths result from a constant value for $P_1$ over the entire velocity range.

This approach does not imply that the pressure distributions have the same shape for differing velocities, or that $p_a$ is equivalent for all velocities, simply that the value of $P_1$ is constant for a particular shot combination over all impact velocities.

4.0 RESULTS

When $P_1$ was assumed to be some constant average pressure on the penetrator tip independent of $V_0$, defined here to be $Q$, for all impact velocities in a particular shot combination, the results of this model improved tremendously. The penetration depth theoretical curves adopted the trends present in the experimental data (i.e. penetration depths leveling off as velocity increases toward 6 km/s). In addition, the values for $Q$ which yielded the best results correlated strongly to target strength.

In order to establish the most credible and most complete correlation possible, all available data were employed. As a consequence, this report includes a large number of figures. In order for this theory to be applied
effectively, experimental data sets must have included crater diameters. A number of the cases in [15] did not provide this information. In addition, a minimum of two data points were required to construct the crater-volume/kinetic energy relationship. Hence, some other cases could not be evaluated. With those limitations in mind, the author applied all readily available cases to this model and reports the results.

Table 1 summarizes all the cases and provides essential data about each shot combination. The figure numbers referred to can be found in Appendix A. One figure number is assigned to each shot combination. For each case, several figures denoted by the figure number and lower case letters are to be found. The first three figures of each case (corresponding to a, b, and c) are the strain vs. impact velocity, crater volume vs. impact kinetic energy, and penetration depth vs. impact velocity graphs. Any additional graphs will occur in pairs. These additional figures (i.e. d and e, etc.) refer to cases in which the crater volume/kinetic energy relationship was modified by removing certain data points that appear erroneous or that may skew the general trend. When these points are removed, the relationship changes - which modifies the resulting penetration curve. Of course this does not alter the strain curve in any way - avoiding the necessity of including an additional one. The first three figures for each case include all reported data.

The figures report theoretical curves superimposed on discrete experimental points. The upper curve in the
Table 1. Experimental Data Summary

<table>
<thead>
<tr>
<th>Fig. #</th>
<th>Penetrator</th>
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<th>$\rho$ (kg/m$^3$)</th>
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**Notes:**
- MPa: Megapascal
- kg/m³: Kilograms per cubic meter
- RHA: Reduced Hardness Area
- Ref.: Reference number
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strain vs. impact velocity figures represents the estimate for $\varepsilon_0$.

When these cases were evaluated, a certain value for $Q$ could be chosen to match the experimental data. This $Q$ was found to strongly correlate to target strength. Table 2 reports each of the different targets present in the data and their corresponding $Q$ value.

<table>
<thead>
<tr>
<th>Fig. #</th>
<th>Penetrator</th>
<th>$Y_p$ (MPa)</th>
<th>$\rho$ (kg/m$^3$)</th>
<th>$L$ (mm)</th>
<th>$L_D$</th>
<th>Target</th>
<th>$R_t$ (MPa)</th>
<th>$\mu^2\rho$ (kg/m$^3$)</th>
<th>Ref.</th>
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Table 2. Correlation of Q to Target Strength

<table>
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<th>Target Material</th>
<th>$R_t$ (MPa)</th>
<th>Q (GPa)</th>
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<tr>
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<tr>
<td>D17</td>
<td>2500</td>
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With this information, it was immediately evident that some strong relationship existed between Q and target strength. Figure D depicts the values for Q chosen to allow the model to predict penetration depths and crater diameters accurately against $R_t$. A curve is fit through the data to both illustrate the correlation and provide a functional relationship between Q and $R_t$. The best fit is

$$ Q = 3.8 \left( 1 - e^{-0.00135 \ R_t} \right) - 0.8 \quad (27) $$
With $Q$ as a direct function of target strength, this one-dimensional penetration model can be expressed in terms of known material properties and impact conditions. The crater volume/kinetic energy curve is still needed to allow for the calculation of penetration depths, however.

To place this revised model into context with regard to those used in [9,10] and [14], Figures E and F are presented. These figures illustrate the comparative results between these three different approaches to the calculation of the average pressure at steady state.
Figure E. Comparative Case: Strain vs. Impact Velocity
(a) indicates theory from [9,10]
(b) indicates theory from [14]
(c) indicates current model

Figure F. Comparative Case: Penetration Depth vs. Impact Velocity
(a) indicates theory from [9,10]
(b) indicates theory from [14]
(c) indicates current model
It is clear from the above comparative figures (i.e. E & F) that each subsequent theory lowered the average pressure in such a way as to lower the penetration and strain curves. This particular case that appears in Figures E and F was chosen for its higher velocity data and its presentation in [14].

5.0 CONCLUSION

In this paper, a new one-dimensional model for the penetration of semi-infinite targets by long rods was developed. Its basis was the revision of the previous techniques employed to calculate the average pressure at steady state. The modified Bernoulli Equation appears to result in values for $P_1$ that are simply too high for the impact cases in the hypervelocity range.

By correlating a new approach to a great body of data, it was discovered that a single value for the average pressure on the penetrator tip at steady state, $Q$, could successfully represent the pressure at steady state over the impact velocity range of 1 to 6 km/s. This formulation for $Q$ allowed the model to improve its accuracy and its trends at higher velocities. In addition, this value $Q$ was shown to have a strong correlation to target strength. This approach has resulted in a completely algebraic solution that relies only on known test parameters and the well established crater volume/kinetic energy relationship.
Future work will involve an effort to revise or replace the modified Bernoulli Equation's estimate for pressure at steady state. Additional analysis also needs to be conducted to ascertain the form of the pressure distribution that leads to a constant Q over all impact velocities.

The aim of this paper was to extend the one-dimensional penetration analysis into the hypervelocity range. The resulting model offers good accuracy for a one-dimensional description of the penetration event.
Appendix A:
Figure 1a. Strain vs. Impact Velocity

Figure 1b. Crater Volume vs. Impact Kinetic Energy
Figure 1c. Penetration Depth vs. Impact Velocity

Figure 1d. Crater Volume vs. Impact Kinetic Energy
Figure 1e. Penetration Depth vs. Impact Velocity

Figure 2a. Strain vs. Impact Velocity
Figure 2b. Crater Volume vs. Impact Kinetic Energy

Figure 2c. Penetration Depth vs. Impact Velocity
Figure 3a. Strain vs. Impact Velocity

Figure 3b. Crater Volume vs. Impact Kinetic Energy
Figure 3c. Penetration Depth vs. Impact Velocity

Figure 4a. Strain vs. Impact Velocity
Figure 4b. Crater Volume vs. Impact Kinetic Energy

Figure 4c. Penetration Depth vs. Impact Velocity
Figure 5a. Strain vs. Impact Velocity

Figure 5b. Crater Volume vs. Impact Kinetic Energy
Figure 5c. Penetration Depth vs. Impact Velocity

Figure 6a. Strain vs. Impact Velocity
Figure 6b. Crater Volume vs. Impact Kinetic Energy

Figure 6c. Penetration Depth vs. Impact Velocity
Figure 6d. Crater Volume vs. Impact Kinetic Energy

Figure 6e. Penetration Depth vs. Impact Velocity
Figure 7a. Strain vs. Impact Velocity

Figure 7b. Crater Volume vs. Impact Kinetic Energy
Figure 7c. Penetration Depth vs. Impact Velocity

Figure 7d. Crater Volume vs. Impact Kinetic Energy
Figure 7e. Penetration Depth vs. Impact Velocity

Figure 8a. Strain vs. Impact Velocity
Figure 8b. Crater Volume vs. Impact Kinetic Energy

Figure 8c. Penetration Depth vs. Impact Velocity
Figure 8d. Crater Volume vs. Impact Kinetic Energy

Figure 8e. Penetration Depth vs. Impact Velocity
Figure 9a. Strain vs. Impact Velocity

Figure 9b. Crater Volume vs. Impact Kinetic Energy
Figure 9c. Penetration Depth vs. Impact Velocity

Figure 9d. Crater Volume vs. Impact Kinetic Energy
Figure 9e. Penetration Depth vs. Impact Velocity

Figure 10a. Strain vs. Impact Velocity
Figure 10b. Crater Volume vs. Impact Kinetic Energy

Figure 10c. Penetration Depth vs. Impact Velocity
Figure 11a. Strain vs. Impact Velocity

Figure 11b. Crater Volume vs. Impact Kinetic Energy
Figure 11c. Penetration Depth vs. Impact Velocity

Figure 12a. Strain vs. Impact Velocity
Figure 12b. Crater Volume vs. Impact Kinetic Energy

Figure 12c. Penetration Depth vs. Impact Velocity
Fig. 13a. Strain vs. Impact Velocity

Fig. 13b. Crater Volume vs. Impact Kinetic Energy
Figure 13c. Penetration Depth vs. Impact Velocity

Figure 14a. Strain vs. Impact Velocity
Figure 14b. Crater Volume vs. Impact Kinetic Energy

Figure 14c. Penetration Depth vs. Impact Velocity
Figure 15a. Strain vs. Impact Velocity

Figure 15b. Crater Volume vs. Impact Kinetic Energy
Figure 15c. Penetration Depth vs. Impact Velocity

Figure 16a. Strain vs. Impact Velocity
Figure 16b. Crater Volume vs. Impact Kinetic Energy

Figure 16c. Penetration Depth vs. Impact Velocity
Figure 17a. Strain vs. Impact Velocity

Figure 17b. Crater Volume vs. Impact Kinetic Energy
Figure 17c. Penetration Depth vs. Impact Velocity

Figure 18a. Strain vs. Impact Velocity
Figure 18b. Crater Volume vs. Impact Kinetic Energy

Figure 18c. Penetration Depth vs. Impact Velocity
Figure 19a. Strain vs. Impact Velocity

Figure 19b. Crater Volume vs. Impact Kinetic Energy
Figure 19c. Penetration Depth vs. Impact Velocity

Figure 20a. Strain vs. Impact Velocity
Figure 20b. Crater Volume vs. Impact Kinetic Energy

Figure 20c. Penetration Depth vs. Impact Velocity
Figure 21a. Strain vs. Impact Velocity

Figure 21b. Crater Volume vs. Impact Kinetic Energy
Figure 21c. Penetration Depth vs. Impact Velocity

Figure 22a. Strain vs. Impact Velocity
Figure 22b. Crater Volume vs. Impact Kinetic Energy

Figure 22c. Penetration Depth vs. Impact Velocity
Figure 23a. Strain vs. Impact Velocity

Figure 23b. Crater Volume vs. Impact Kinetic Energy
Figure 23c. Penetration Depth vs. Impact Velocity

Figure 24a. Strain vs. Impact Velocity
Figure 24b. Crater Volume vs. Impact Kinetic Energy

Figure 24c. Penetration Depth vs. Impact Velocity
Figure 24d. Crater Volume vs. Impact Kinetic Energy

Figure 24e. Penetration Depth vs. Impact Velocity
Figure 25a. Strain vs. Impact Velocity

Figure 25b. Crater Volume vs. Impact Kinetic Energy
Figure 25c. Penetration Depth vs. Impact Velocity

Figure 25d. Crater Volume vs. Impact Kinetic Energy
Figure 25a. Penetration Depth vs. Impact Velocity

Figure 26a. Strain vs. Impact Velocity
Figure 26b. Crater Volume vs. Impact Kinetic Energy

Figure 26c. Penetration Depth vs. Impact Velocity
Figure 27a. Strain vs. Impact Velocity

Figure 27b. Crater Volume vs. Impact Kinetic Energy
Figure 27c. Penetration Depth vs. Impact Velocity

Figure 27d. Crater Volume vs. Impact Kinetic Energy
Figure 27e. Penetration Depth vs. Impact Velocity

Figure 28a. Strain vs. Impact Velocity
Figure 28b. Crater Volume vs. Impact Kinetic Energy

Figure 28c. Penetration Depth vs. Impact Velocity
Figure 28d. Crater Volume vs. Impact Kinetic Energy

Figure 28e. Penetration Depth vs. Impact Velocity
Figure 29a. Strain vs. Impact Velocity

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Figure 29b. Crater Volume vs. Impact Kinetic Energy
Figure 29c. Penetration Depth vs. Impact Velocity

Figure 30a. Strain vs. Impact Velocity
Figure 30b. Crater Volume vs. Impact Kinetic Energy

Figure 30c. Penetration Depth vs. Impact Velocity
Figure 30d. Crater Volume vs. Impact Kinetic Energy

Figure 30e. Penetration Depth vs. Impact Velocity
Figure 31a. Strain vs. Impact Velocity

Figure 31b. Crater Volume vs. Impact Kinetic Energy
Figure 31c. Penetration Depth vs. Impact Velocity

Figure 32a. Strain vs. Impact Velocity
Figure 32b. Crater Volume vs. Impact Kinetic Energy

Figure 32c. Penetration Depth vs. Impact Velocity
Figure 32d. Crater Volume vs. Impact Kinetic Energy

Figure 32e. Penetration Depth vs. Impact Velocity
Figure 33a. Strain vs. Impact Velocity

Figure 33b. Crater Volume vs. Impact Kinetic Energy
Figure 33c. Penetration Depth vs. Impact Velocity

Figure 34a. Strain vs. Impact Velocity
Figure 34b. Crater Volume vs. Impact Kinetic Energy

Figure 34c. Penetration Depth vs. Impact Velocity
Figure 35a. Strain vs. Impact Velocity

Figure 35b. Crater Volume vs. Impact Kinetic Energy
Figure 35c. Penetration Depth vs. Impact Velocity

Figure 35d. Crater Volume vs. Impact Kinetic Energy
Figure 35e. Penetration Depth vs. Impact Velocity

Figure 36a. Strain vs. Impact Velocity
Figure 36b. Crater Volume vs. Impact Kinetic Energy

Figure 36c. Penetration Depth vs. Impact Velocity
Figure 37a. Strain vs. Impact Velocity

Figure 37b. Crater Volume vs. Impact Kinetic Energy
Figure 37c. Penetration Depth vs. Impact Velocity

Figure 37d. Crater Volume vs. Impact Kinetic Energy
Figure 37e. Penetration Depth vs. Impact Velocity

Figure 38a. Strain vs. Impact Velocity
Figure 38b. Crater Volume vs. Impact Kinetic Energy

Figure 38c. Penetration Depth vs. Impact Velocity
Figure 38d. Crater Volume vs. Impact Kinetic Energy

Figure 38e. Penetration Depth vs. Impact Velocity
Figure 39a. Strain vs. Impact Velocity

Figure 39b. Crater Volume vs. Impact Kinetic Energy
Figure 39c. Penetration Depth vs. Impact Velocity

Figure 39d. Crater Volume vs. Impact Kinetic Energy
Figure 39e. Penetration Depth vs. Impact Velocity

Figure 40a. Strain vs. Impact Velocity
Figure 40b. Crater Volume vs. Impact Kinetic Energy

Figure 40c. Penetration Depth vs. Impact Velocity
Figure 40d. Crater Volume vs. Impact Kinetic Energy

Figure 40e. Penetration Depth vs. Impact Velocity
Figure 41a. Strain vs. Impact Velocity

Figure 41b. Crater Volume vs. Impact Kinetic Energy
Figure 41c. Penetration Depth vs. Impact Velocity

Figure 41d. Crater Volume vs. Impact Kinetic Energy
Figure 41e. Penetration Depth vs. Impact Velocity

Figure 42a. Strain vs. Impact Velocity
Figure 42b. Crater Volume vs. Impact Kinetic Energy

Figure 42c. Penetration Depth vs. Impact Velocity
Figure 42d. Crater Volume vs. Impact Kinetic Energy

Figure 42e. Penetration Depth vs. Impact Velocity
Figure 43a. Strain vs. Impact Velocity

Figure 43b. Crater Volume vs. Impact Kinetic Energy
Figure 43c. Penetration Depth vs. Impact Velocity

Figure 44a. Strain vs. Impact Velocity
Figure 44b. Crater Volume vs. Impact Kinetic Energy

Figure 44c. Penetration Depth vs. Impact Velocity
Figure 44d. Crater Volume vs. Impact Kinetic Energy

Figure 44e. Penetration Depth vs. Impact Velocity
Figure 45a. Strain vs. Impact Velocity

Figure 45b. Crater Volume vs. Impact Kinetic Energy
Figure 45c. Penetration Depth vs. Impact Velocity

Figure 45d. Crater Volume vs. Impact Kinetic Energy
Figure 45e. Penetration Depth vs. Impact Velocity

Figure 45f. Crater Volume vs. Impact Kinetic Energy
Figure 45g. Penetration Depth vs. Impact Velocity

Figure 46a. Strain vs. Impact Velocity
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Figure 46c. Penetration Depth vs. Impact Velocity
Figure 46d. Crater Volume vs. Impact Kinetic Energy

Figure 46e. Penetration Depth vs. Impact Velocity
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Figure 47c. Penetration Depth vs. Impact Velocity

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Figure 48b. Crater Volume vs. Impact Kinetic Energy

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Figure 49c. Penetration Depth vs. Impact Velocity

Figure 49d. Crater Volume vs. Impact Kinetic Energy
Figure 49e. Penetration Depth vs. Impact Velocity

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Figure 53c. Penetration Depth vs. Impact Velocity

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Figure 54d. Crater Volume vs. Impact Kinetic Energy

Figure 54e. Penetration Depth vs. Impact Velocity
Figure 55a. Strain vs. Impact Velocity

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Figure 56d. Crater Volume vs. Impact Kinetic Energy

Figure 56e. Penetration Depth vs. Impact Velocity
Figure 57a. Strain vs. Impact Velocity

Figure 57b. Crater Volume vs. Impact Kinetic Energy
Figure 57c. Penetration Depth vs. Impact Velocity

Figure 58a. Strain vs. Impact Velocity
Figure 58b. Crater Volume vs. Impact Kinetic Energy

Figure 58c. Penetration Depth vs. Impact Velocity
Figure 59a. Strain vs. Impact Velocity

Figure 59b. Crater Volume vs. Impact Kinetic Energy
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Figure 60a. Strain vs. Impact Velocity
Figure 60b. Crater Volume vs. Impact Kinetic Energy

Figure 60c. Penetration Depth vs. Impact Velocity
Figure 61a. Strain vs. Impact Velocity

Figure 61b. Crater Volume vs. Impact Kinetic Energy
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Figure 62a. Strain vs. Impact Velocity
Figure 62b. Crater Volume vs. Impact Kinetic Energy

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Figure 62d. Crater Volume vs. Impact Kinetic Energy

Figure 62e. Penetration Depth vs. Impact Velocity
Figure 63a. Strain vs. Impact Velocity

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Figure 64a. Strain vs. Impact Velocity
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Figure 66a. Strain vs. Impact Velocity
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Figure 66e. Penetration Depth vs. Impact Velocity
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6.0 REFERENCES


18. [3] as reported in [15]


20. [5], also reported in [15]

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24. reported in [9]


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