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DYNAMIC CRACK CURVING - A PHOTOELASTIC EVALUATION

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## DYNAMIC CRACK CURVING - A PHOTOELASTIC EVALUATION

by

M. Ramulu\* and A. S. Kobayashi\*\*

A dynamic crack curving criterion, which is valid under combined modes I and II or mode I loading and which is based on either the maximum circumferential stress or minimum strain energy density factor at a reference distance of  $r_0 = \frac{1}{128\pi} \left( \frac{K_I}{\sigma_{OX}} \right)^2 V^2$  ( $c, c_1, c_2$ ) crack deformation, is developed. Directional stability of a mode I crack propagation is attained when  $r_0 > r_c$ , where  $r_c$  for Homalite-100 was determined from dynamic photoelastic experiments. In the presence of mode II crack deformation, positive remote stress component, i.e.,  $\sigma_{OX} > 0$  and negative remote stress component, i.e.,  $\sigma_{OX} < 0$ , was found to enhance and suppress crack curving, respectively.

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The objective of the present study is to derive a dynamic crack curving criterion applicable to both mode I and combined modes I and II crack tip deformation. To this goal, dynamic extension of two static crack curving criteria, that is the maximum circumferential stress criterion and the minimum strain energy density criterion at a critical distance  $r_c$ , was considered. The developed theoretical relations were evaluated numerically and the influence of  $\sigma_{ox}$  and crack velocity on crack curving direction were deduced. Crack curving angles predicted by the two dynamic crack curving criteria were then compared with experimental results, obtained from past dynamic photoelastic investigation.

#### DYNAMIC CRACK CURVING CRITERIA

##### Elasto-dynamic Crack Tip Stress Field

The dynamic crack curving criteria, are derived from the near field, mixed mode elasto-dynamic state of stress associated with a crack tip propagating at constant velocity. This dynamic state of stress is given by Freund (17,18) in terms of local rectangular and polar co-ordinates of  $(x,y)$  and  $(r,\theta)$ , respectively, with origin at the crack tip, and the mode I and II dynamic stress intensity factors,  $K_I$  and  $K_{II}^*$ , respectively. The authors [19] have added to Freund's near field, dynamic state of stress the second order term of  $\sigma_{ox}$  which is acting parallel to the direction of crack extension. This dynamic singular crack tip stress field under mixed mode loading for small  $\theta$  values differs from the corresponding static stress field in that the largest principal singular tensile stress acts parallel to the x-axis, a fact which not only contributes to crack curving but also to dynamic crack branching. Furthermore, this region

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\*The superscript "dyn" to identify dynamic stress intensity factor will not be used in this paper, since all quantities refer to dynamic values.

ahead of the running crack where  $|\sigma_{xx}| > \sigma_{yy}$  increases with increases in crack speed and  $\sigma_{ox}$  even under pure mode-II crack tip deformation [19]. This inevitable involvement of  $\sigma_{ox}$  forms the basis of incorporating  $\sigma_{ox}$  in the dynamic crack curving criteria presented in this paper.

#### Maximum Circumferential Stress Theory

The angle,  $\theta_c$ , at which circumferential stress,  $\sigma_{\theta\theta}$ , is maximum, can be obtained from the following,

$$\frac{\partial \sigma_{\theta\theta}}{\partial \theta} = 0 \quad \sigma_{\theta\theta} > 0 \quad (1)$$

where the added  $\sigma_{\theta\theta} > 0$  is to assure fracture under tensile state of stress. Equation (1), when evaluated in conjunction with a pure mode I dynamic crack tip state of stress will yield a transcendental relation between the critical values of  $\theta$  and  $r$ . Furthermore, by setting  $\theta = 0$  in Equation (1), we obtain

$$r_0 = \frac{1}{128\pi} \left\{ \left( \frac{K_I}{\sigma_{ox}} \right) V(c, c_1, c_2) \right\}^2 \quad (2a)$$

$$\text{where } V(c, c_1, c_2) = \left[ B_I(c) \left\{ -(1+S_2^2)(2-3S_1^2) - \frac{4S_1S_2}{1+S_2^2} (14+3S_2^2) - 16S_1(S_1-S_2) + 16(1+S_2^2) \right\} \right] \quad (2b)$$

$$B_I(c) = \left\{ \frac{(1+S_2^2)}{4S_1S_2 - (1+S_2^2)^2} \right\} \quad (2c)$$

$$S_1^2 = \left[ 1 - \frac{c^2}{c_1^2} \right] \quad ; \quad S_2^2 = \left[ 1 - \frac{c^2}{c_2^2} \right] \quad (2d)$$

and  $c$ ,  $c_1$  and  $c_2$  are the crack velocity, dilatational wave velocity, and distortional wave velocity, respectively. It can be easily shown that for zero crack velocity or  $c = 0$ , Equation (2a) reduces to Streit and Finnie's solution

$$[14] \text{ of } r_0 = \frac{9}{128\pi} \left( \frac{K_I}{\sigma_{ox}} \right)^2$$

Figure 1 shows the velocity effect on  $r_0$  which is plotted in a non-dimensional form of  $\left[ \sqrt{16} \frac{\sigma_{ox}}{K_I} \right]$  where the dynamic  $r_0$  is always less than the corresponding static  $r_0$  at experimentally observed crack velocities of  $0 < c \leq 0.33$  and is insensitive to the sign of  $\sigma_{ox}$ . The terminal crack velocity of  $c/c_1 = 0.325$ , in Figure 1 where  $r_0 = 0$  coincides with the terminal velocity predicted by Yoffe [15].

#### Minimum Strain Energy Density Theory

According to this theory, the crack will extend to the location of the minimum strain energy density factor,  $S_{min}$ , or

$$\frac{\partial S}{\partial \theta} = 0 \text{ at } \theta = \theta_c \quad (3)$$

The intensity of the strain energy density,  $S$ , for the state of plane strain can be written as

$$S = \frac{1}{2} \frac{(1+\nu)}{E} \left[ (1-\nu)(\sigma_{xx}^2 + \sigma_{yy}^2) - 2\nu(\sigma_{xx} \cdot \sigma_{yy}) + 2\sigma_{xy}^2 \right] \quad (4)$$

where  $E$  and  $\nu$  are the modulus of elasticity and Poisson's ratio, respectively. Substituting the dynamic mixed mode crack tip stresses into Equation (4) and then into Equation (3) yields

$$\left\{ [(1-\nu)\sigma_{xx} - \nu\sigma_{yy}] \frac{\partial \sigma_{xx}}{\partial \theta} + [(1-\nu)\sigma_{yy} - \nu\sigma_{xx}] \frac{\partial \sigma_{yy}}{\partial \theta} + 2\sigma_{xy} \frac{\partial \sigma_{xy}}{\partial \theta} \right\} + \sqrt{16} \sigma_{ox} \left\{ (1-\nu) \frac{\partial \sigma_{xx}}{\partial \theta} - \nu \frac{\partial \sigma_{yy}}{\partial \theta} \right\} = 0 \quad (5)$$

By setting Poisson's ratio  $\nu = 1/3$ ,  $\sigma_{ox} = 0$  as a crack velocity of  $c \rightarrow 0$  in Equations (5), the static angular predictions in Reference [13] are recovered. When a non-vanishing second order term of  $\sigma_{ox}$  is considered in Equation (5)

yields four  $\theta_e$  values, a pair for  $S_{\max}$  and another pair of  $S_{\min}$  for given values of  $c$ ,  $K_{II}/K_I$ ,  $r_0$  and  $\sigma_{ox}$ . Only the negative root of  $\theta_e$  corresponding to positive  $K_{II}/K_I$  and the positive root of  $\theta_e$  for negative  $K_{II}/K_I$ , to the tensile loading are of interest [13]. Numerical values of these  $\theta_e$  will be discussed in the following section.

Actual evaluation of Equation (3) will show that curving of a straight crack propagating at the lower velocity can be considered only by incorporating the nonsingular term of  $\sigma_{ox}$  in the minimum strain energy density criteria. Such possibility of crack curving without  $K_{II}$  values and under the minimum strain energy criterion has not been considered by others.

#### Comparison of Maximum $\sigma_{ox}$ and Minimum S Criterion

Figure 2 shows the predicted crack curving angles for crack velocities,  $c/c_1$ , from 0 to 0.25 by maximum stress and the minimum strain energy density criteria when  $\sigma_{ox} = 0$ . Without the second order term, both criteria predicted the same crack curving angles for much of the crack velocity range. Although the crack curving angle at higher crack velocities are significant for lower crack velocities of  $c/c_1 \leq 0.15$ , the predicted crack curving angle, which is referred to as fracture angle from hereon, is almost constant and is in close agreement with corresponding static fracture angles.

The effects of the non-singular term of  $\sigma_{ox}$  and reference radius  $r_0$  in predicting the fracture angle by both maximum  $\sigma_{\theta\theta}$  and minimum S theories at various crack velocities are shown in Figure 3 for  $\nu = 1/3$ , and  $K_{II}/K_I = -0.1$ , and  $\sigma_{ox}/K_I = -1.0$  and  $1.0$ . Note that fracture angle for negative  $\sigma_{ox}$ , are much smaller than those with positive  $\sigma_{ox}$ . Also, larger  $r_0$  results in larger changes in the fracture angle. For larger values of  $r_0$ , the differences in predicted fracture angles due to maximum circumferential stress theory and minimum strain

energy density theory are larger at higher crack velocities. This importance of  $r_0$  value in characterizing the direction of the fracture angles is discussed in Reference [14].

## EXPERIMENTAL VERIFICATION

### Dynamic Isochromatics:

For a single, pure mode-I or combined modes I and II crack propagating at a constant velocity, the dynamic crack tip isochromatic patterns together with the predicted path are shown in Figure 4. Changes in the remote stress,  $\sigma_{ox}$ , results in backward or forward tilting of the dynamic isochromatics. For a given  $\sigma_{ox}$ , the change in the sign of  $K_{II}$ , results in a mirror image of the isochromatics. Detailed discussion of the changes in dynamic isochromatics with variations in  $K_{II}/K_I$  and  $\sigma_{ox}/K_I$  can be found in Reference [19].

### Data Reduction Procedure

Experimentally determined dynamic isochromatics surrounding a running crack often exhibits moderate unsymmetry. Such photoelastic patterns were heretofore considered experimental abnormalities and were ignored by averaging the unsymmetric patterns during the data reduction process. Careful postmortem inspection of the fracture specimens, however, show that the higher magnitudes of  $\sigma_{ox}$  of isochromatics and slightly unsymmetric isochromatics are often associated with slightly curved crack patterns. With the development of a data reduction procedure [19] for evaluating dynamic  $K_{II}$  together with  $K_I$  and  $\sigma_{ox}$  values, it became possible to investigate the above criteria by extracting  $K_I$  and  $K_{II}$  and  $\sigma_{ox}$  from the previously recorded dynamic isochromatics surrounding running crack tips of curved cracks. An optimization method developed by the authors based on the overdeterministic least square procedure was also used to extract the dynamic three parameters  $K_I$ ,  $K_{II}$  and  $\sigma_{ox}$  from the recorded dynamic photoelastic pattern surrounding a running crack [19,20].

The dynamic crack curving criteria developed for pure mode-I loading conditions require accurate determination of  $K_I$  and  $\sqrt{\sigma_{ox}}$ . Accuracy of the data reduction procedure used in this investigation was verified by using the above data reduction procedure to calculate  $K_I$  and  $\sqrt{\sigma_{ox}}$  from previously generated isochromatics generated by three parameters of  $K_I$ ,  $\sqrt{\sigma_{ox}}$ , and  $A_3$  with  $K_{II} = 0$  [21]. The recovered two dynamic parameters  $K_I$  and  $\sqrt{\sigma_{ox}}$  agreed within  $\pm 0.5\%$  and  $\pm 5\%$ , respectively, with the given results. This series of numerical experiments showed that the two parameter characterization procedure involving  $K_I$  and  $\sqrt{\sigma_{ox}}$  should describe reasonably well the stress field in the vicinity of a running crack tip.

The crack curving angle was measured along the crack path by averaging the measured crack curving angle on front and back surfaces of the fractured specimen since the crack surfaces of some of the curved cracks were not perpendicular to the specimen surfaces. The maximum variation between the front and back crack curving angles was about 3 degrees for severely curved cracks. Similar differences in out-of-phase crack curving were also observed by Williams et al., in their PMMA specimens [2].

### Results

Figure 5 shows three frames out of a 16 frame dynamic photoelastic record of a curving crack in a Homalite-100 dynamic tear test (DTT) specimen of 9.58 mm (3/8 in) thick, 88.9 x 400 mm (3 1/2 x 15 in). This beam with a blunt initial crack of 6.4 mm (7/32 in) in length was impact loaded by a drop weight of 1.48 kg (3.25 lb) [22]. The crack emanated from the blunt saw-cut crack and propagated through much of the height of the beam prior to curving near the region of impact loading. Further details of the experimental setup, crack velocity measurements and dynamic calibration of the Homalite-100 material used are found in

Reference [22]. Figure 6 shows  $K_I$ ,  $K_{II}$ ,  $\overline{\sigma}_{ox}$  and  $r_o$  which is computed by Equation (2), obtained from the dynamic photoelastic pattern preceding and immediately after the crack curving in Figure 5.  $K_{II}$  is negligible in comparison to  $K_I$  and at the point of instability and pronounced fluctuation in  $\overline{\sigma}_{ox}$  is noted. After crack curving  $K_{II}$  and  $\overline{\sigma}_{ox}$  increased while  $K_I$  and crack velocity dropped rapidly.  $r_o$  was close to 1.5 mm throughout crack propagation and reached a minimum value of  $r_c = 1$  mm during the onset of crack curving.

Figure 7 shows a slightly curved crack and the associated  $K_I$ ,  $K_{II}$ ,  $\overline{\sigma}_{ox}$  and  $r_o$  in a fracturing 9.53 mm (3/8 in) thick, 254 x 254 mm (10 x 10 in) single-edge-notch (SEN) Homalite-100 specimen [23]. Gradual increase and decrease of  $K_I$  and a very small  $K_{II}$  with a rapid fluctuation of  $\overline{\sigma}_{ox}$  and  $r_o$  are noted. Three SEN results were evaluated where  $K_I$  reached a maximum value,  $K_{II}$  was negligible and  $\overline{\sigma}_{ox}$  was increasing prior to crack curving. At the onset of instability, a sudden drop in  $K_I$  and larger  $\overline{\sigma}_{ox}$  with  $K_{II} = 0$  are observed.  $r_o$  dropped sharply to an average value of 1.5 mm at the point of instability. This minimum  $r_o$  value will be referred to  $r_c$  which will be found to be a material parameter associated with dynamic crack curving. The small negative  $K_{II}$ , which appeared immediately after crack instability, resulted in a positive angle of crack curving. This result is not only in agreement with the analytically predicted angles in Figure 3 but is also in agreement with similar observation in crack curving under stable crack growth conditions [24]. The rapid oscillations of  $r_o$  in all the three SEN specimens appeared to be related to the rapid but opposing oscillations in  $\overline{\sigma}_{ox}$ .

Figure 8 shows a curved crack and the associated,  $K_I$ ,  $K_{II}$ ,  $\overline{\sigma}_{ox}$  and  $r_o$  in a Homalite-100, wedge-loaded, rectangular double cantilever beam (WL-RDCB) specimen of 9.6 mm (3/8 in) thick and 76.2 x 152.4 mm (3 x 6 in) with a blunt initial crack of length 2.36 mm (0.093 in). Experimental details of this series

of tests can be found in Reference [25]. Fluctuations in dynamic fracture parameters  $K_I$ ,  $K_{II}$ ,  $\overline{\sigma_{ox}}$  and  $r_o$  is noted all along the curved crack path. The crack curved continuously without any kinks and is a characteristic fracture path of DCB specimens.

Figure 9 shows five frames out of a 16-frame dynamic photoelastic record of a curving crack in a 9.53 mm (3/8 in) thick, 254 x 254 mm (10 x 10 in) Homalite-100 single edge notch (SEN) specimen loaded under fixed gripped tension. The crack emanated from a small precrack 150  $\mu$ sec after impact by a flat-nosed projectile. The severe stress wave reflections in this specimen caused the crack to curve continuously in a zig-zag manner. Details of this experiment can be found in Reference [26]. Figure 10 shows the corresponding  $K_I$ ,  $K_{II}$ ,  $\overline{\sigma_{ox}}$  and  $r_o$  variations associated with the unsymmetric dynamic isochromatics in this test. Severe stress wave loading caused the crack to curve immediately after propagation and  $r_c$  is about 1.35 mm at this crack kinking. Throughout crack propagation,  $\overline{\sigma_{ox}}$  changed signs and is related to the zig-zagged crack path.

Fracture angles of curved cracks measured in nine dynamic photoelasticity tests and the corresponding fracture angles computed by the maximum  $\overline{\sigma_{\theta\theta}}$  and minimum  $S$  theories are summarized in the Table 1. Remarkable agreements in experimentally measured and numerically computed results by both the theories, using an experimentally measured  $r_c \approx 1.3$  mm for Homalite-100 are noted. Crack curving in our experiments for mode I, crack propagation ranged between  $\pm 25^\circ$  to a minimum of  $2^\circ$  for severe to moderate curving.

#### DISCUSSIONS

The closed form elasticity solution for a circular arc crack under uniform stress field provides a simple check on the accuracy of using the near field solution of a straight crack in the results cited above. The static solution given by Panasyuk and Brezhnitskiy [27] in the vicinity of a circular arc crack

with an included angle  $2\alpha$  differ with straight crack solution only by a multiplication factor of

$$K_I^{\text{curved}} = K_I^{\text{straight}} \cos\alpha/2 / (1 + \sin^2\alpha/2) \quad (6a)$$

$$K_{II}^{\text{curved}} = K_{II}^{\text{straight}} \sin\alpha/2 / (1 + \sin^2\alpha/2) \quad (6b)$$

$$\bar{\sigma}_{ox}^{\text{curved}} = \bar{\sigma}_{ox}^{\text{straight}} \sin^2\alpha/2 / (1 + \sin\alpha/2) \quad (6c).$$

where the superscripts "straight" and "curved" refer to crack tip parameters associated with a straight and curved crack, respectively. As an estimate of possible errors involved in using a straight crack solution to evaluate the fracture parameters of a curved crack were determined by least square fitting the above exact static solution of a curved crack and the corresponding solution for a straight crack to the two extreme curved cracks associated with the last data points in Figures 6 and 8. The resultant  $K_I$ ,  $K_{II}$  and  $\bar{\sigma}_{ox}$  of the straight crack solutions are within 10%, 28% and 6%, respectively of the corresponding solutions for circular arc cracks of  $\alpha = 25$  and  $28^\circ$ . Thus, possible error introduced by using a second order dynamic crack tip state of stress of a straight crack in place of a curved crack should be negligible for most of the curved crack problems of  $\alpha = 5 \sim 10^\circ$  in this investigation.

The developed dynamic crack curving criterion shows the large  $\bar{\sigma}_{ox}$  contributes to crack instability and is in agreement with Benbow and Roesler's conclusion involving static experiments [28]. Cotterell [29-31] referring to Williams analysis [32], showed that the crack path will be unstable when  $\bar{\sigma}_{ox}$  is positive. The above static crack stability criterion [28-31] correlates well with the experimental results of DCB specimens but cannot explain dynamic crack curvings

in fracture specimens of SEN, CT, and DTT where  $\overline{\sigma}_{ox}$  is negative. The proposed criterion for the directional stability of a propagating crack is independent of the sign of the  $\overline{\sigma}_{ox}$ , and is thus applicable to all crack curving data considered in this paper.

As shown in Figure 3 the influence of non-singular stress is more pronounced for moderate values of  $r_0$  irrespective of the sign of  $K_{II}/K_I$ . This result re-emphasizes the importance of the non-singular stress term  $\overline{\sigma}_{ox}$ , which, when neglected, can lead to inaccurate results as observed by Tirosh [33].

Considering the fact that dynamic photoelasticity experiments cited in this paper were conducted by four different investigators over a period of ten years with different shipments of Homalite-100, the consistent results of  $r_c \approx 1.3$  mm is noticeable. In a critical review on  $r_c$  associated with the minimum S criterion of crack curving, Theocaris and Andriopoulos (14) also determined experimentally  $r_c = 1.3$  mm (0.05") for polymethylmethacrylate.

Finally, the crack curving criterion by Karihaloo et al. [12] requires that  $K_{II}$  be known immediately before and after crack curving. The lack of sensitivity in this analysis precluded precise variations of the very small  $K_{II}$  before or after crack curving and thus this crack curving could not be checked.

## CONCLUSIONS

1. A dynamic crack curving criterion based on the directional stability of a running crack under pure mode-I loading is developed.
2. Dynamic fracture angle under pure mode I and mixed mode I and II conditions can be predicted by using either the maximum circumferential stress or the minimum strain energy density theories with the non-singular stress term  $\overline{\sigma}_{ox}$ .
3. Positive  $\overline{\sigma}_{ox}$  always enhances the crack curving and negative  $\overline{\sigma}_{ox}$  reduces the fracture angle irrespective of the sign of  $K_{II}/K_I$ .

4. Experimental results with and without  $K_{II}$  proved that  $r_c$  is a material constant. The critical value of Homalite-100 is  $r_c = 1.3 \text{ mm (0.05 in)}$ .

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TABLE I  
SUMMARY OF EXPERIMENTAL AND  
THEORETICAL RESULTS

Total Number of Experiments:	9
Type of Fracture Specimen:	DTT, SEN, WL-RDCB
Number of Data Points:	81
Crack Velocity, $c/c_1$ :	0.03 to 0.21
$K_I$ (MPa m)	0.50 to 1.59
$K_{II}/K_I$	-0.22 to 0.18
$\sigma_x/K_I$	-2.89 to 4.04
Experimental Fracture Angle Associated with Crack Curving	$-20^\circ$ to $26^\circ$
Theoretical Prediction of Fracture Angle	$-20^\circ$ to $25^\circ$
$r_c$ (mm)	1.0 to 1.5

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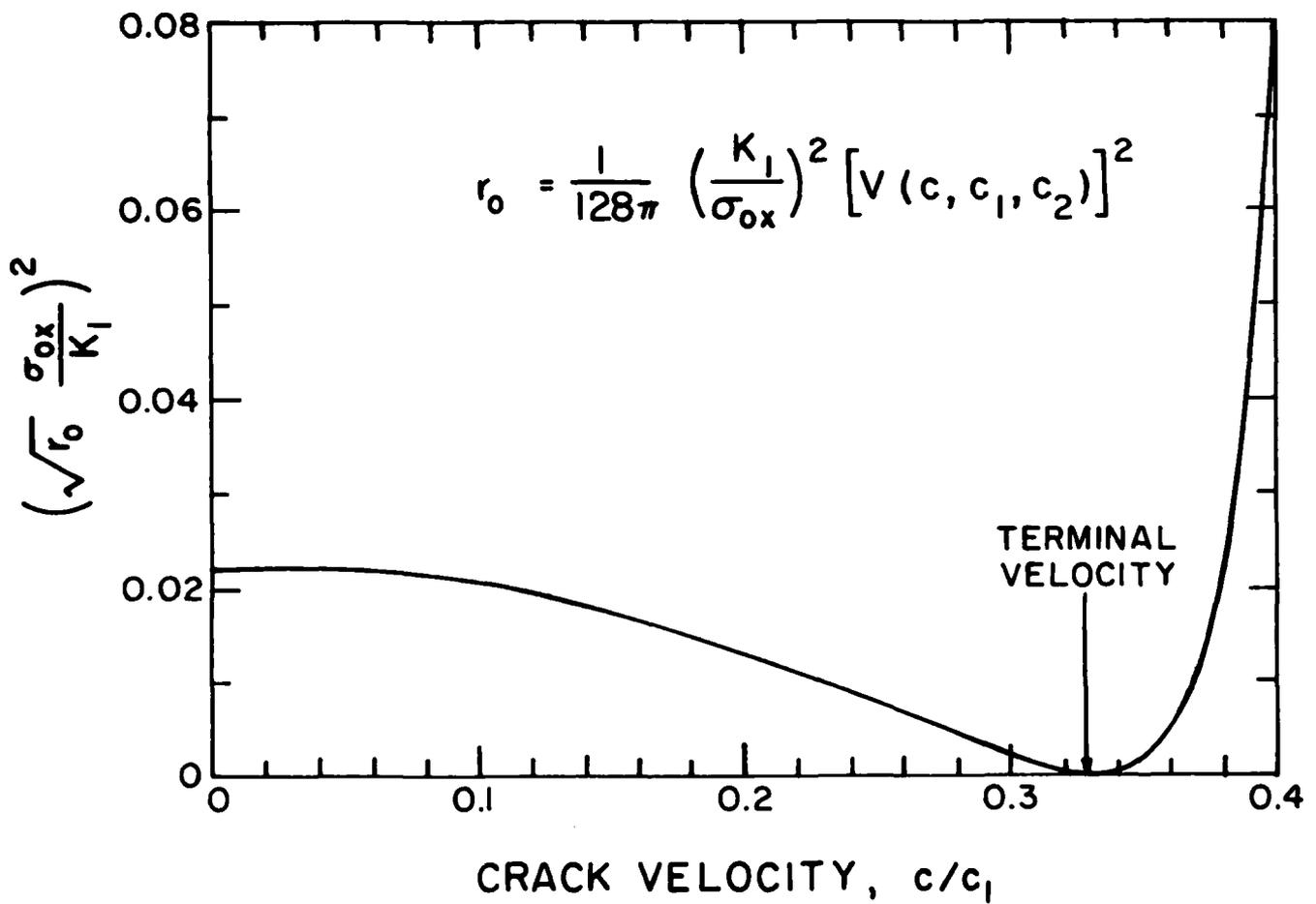


FIG. 1. NON DIMENSIONALIZED REMOTE STRESS vs CRACK VELOCITY.

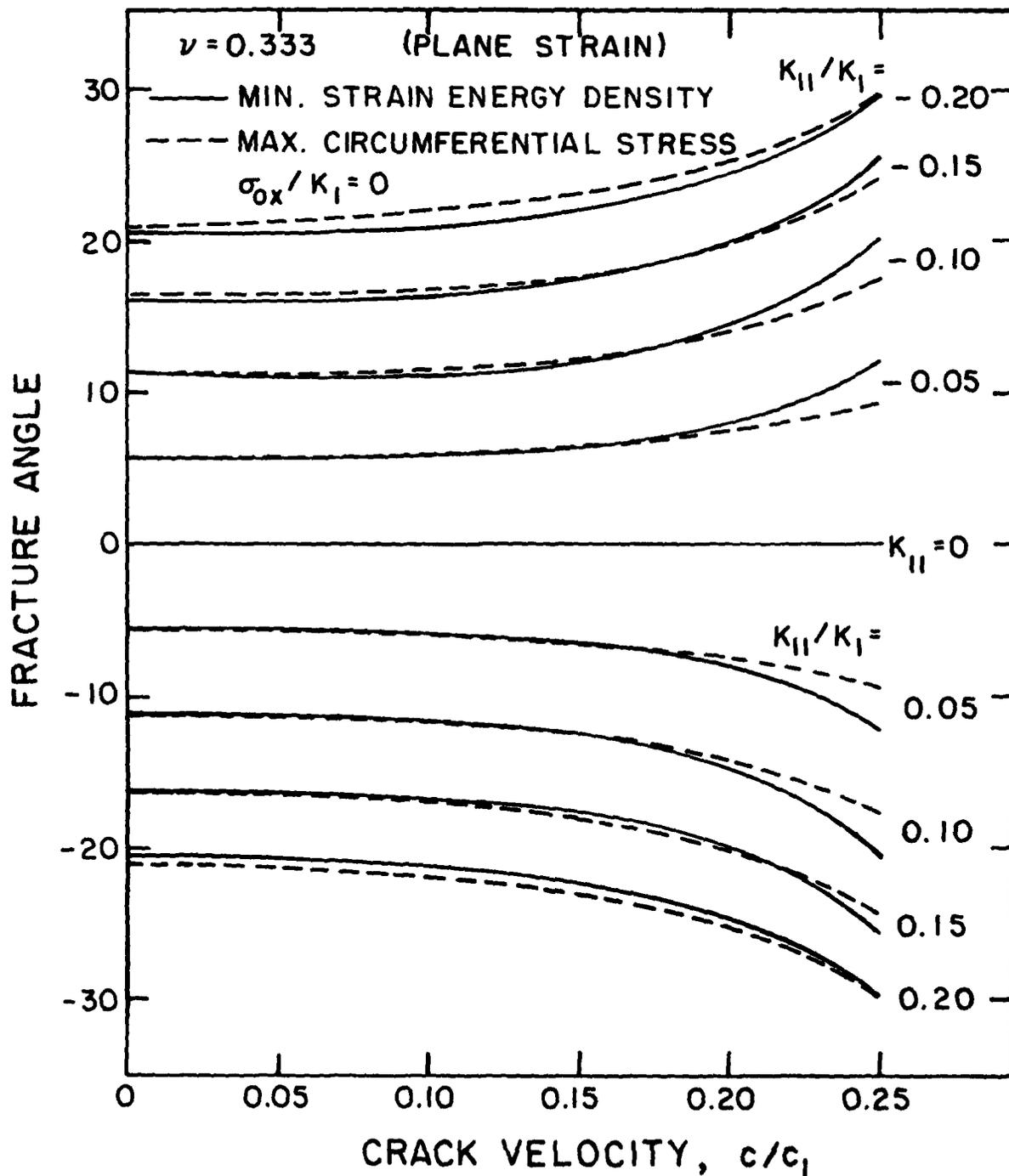


FIG. 2. DYNAMIC CRACK EXTENSION ANGLE FOR MIXED MODE LOADING BY MAX. CIRCUMFERENTIAL STRESS CRITERION AND MIN. STRAIN ENERGY DENSITY CRITERION WITHOUT NON-SINGULAR STRESS.

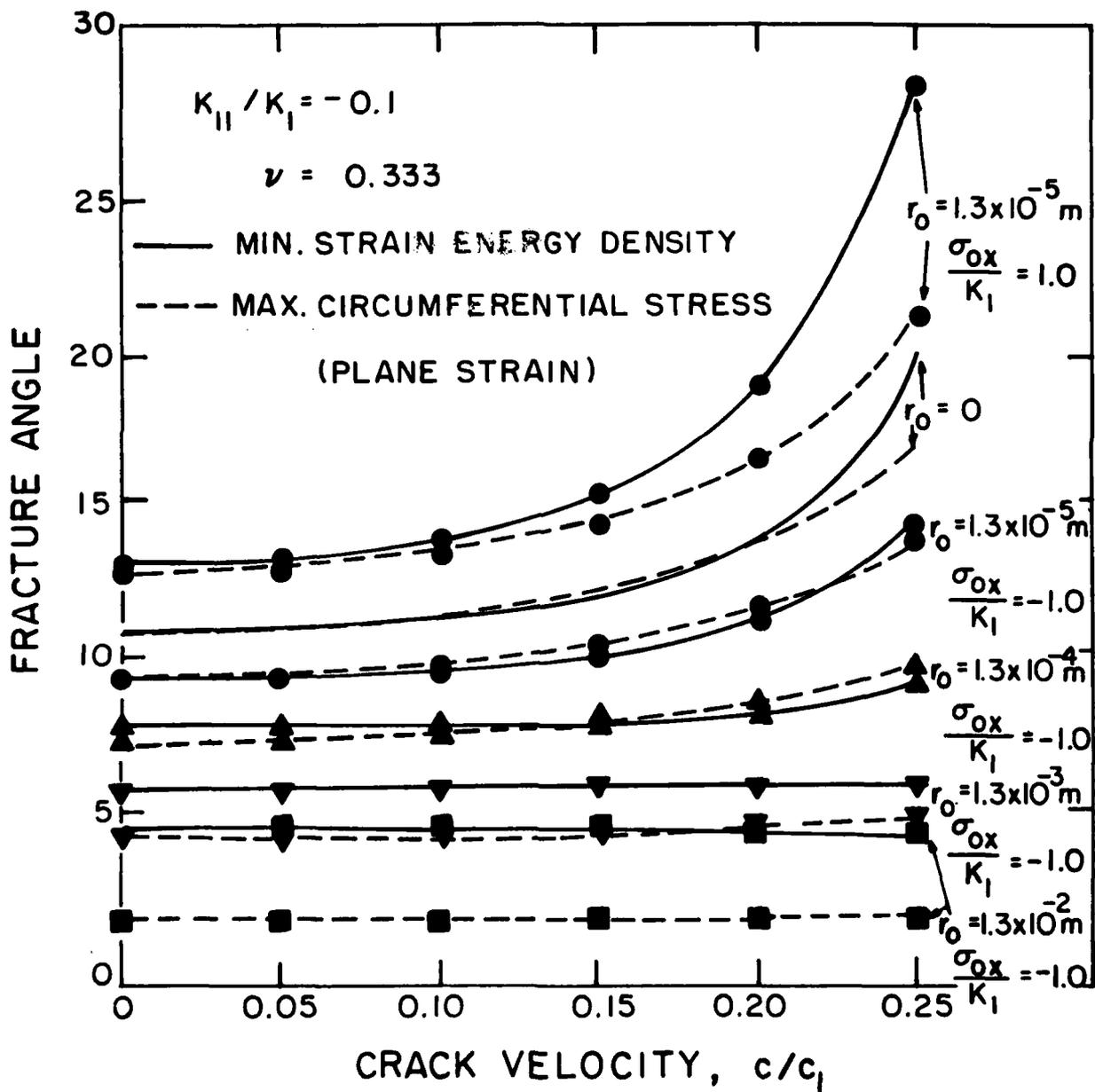
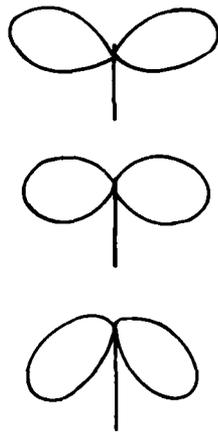


FIG. 3. EFFECT OF REFERENCE RADIUS  $r_0$  FOR PREDICTING DYNAMIC CRACK EXTENSION ANGLE AT  $K_{II} / K_I = -0.1$  BY MAX. CIRCUMFERENTIAL STRESS CRITERION AND MINIMUM STRAIN ENERGY DENSITY CRITERION WITH VARYING NON-SINGULAR STRESS.

CRACK INSTABILITY BY DYNAMIC PHOTOELASTICITY  
UNDER PURE MODE I CONDITIONS

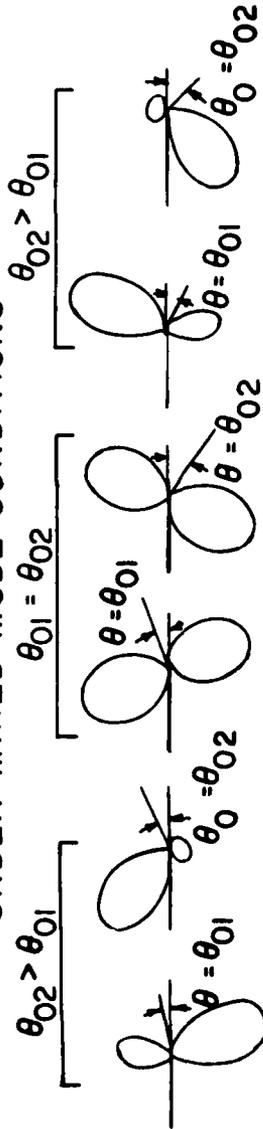


$$\sigma_{ox} > 0 \quad \sigma_{ox} = 0 \quad \sigma_{ox} < 0$$

INSTABILITY STABLE LESS UNSTABLE

FRACTURE PATH PREDICTIONS BY DYNAMIC PHOTOELASTICITY

UNDER MIXED MODE CONDITIONS



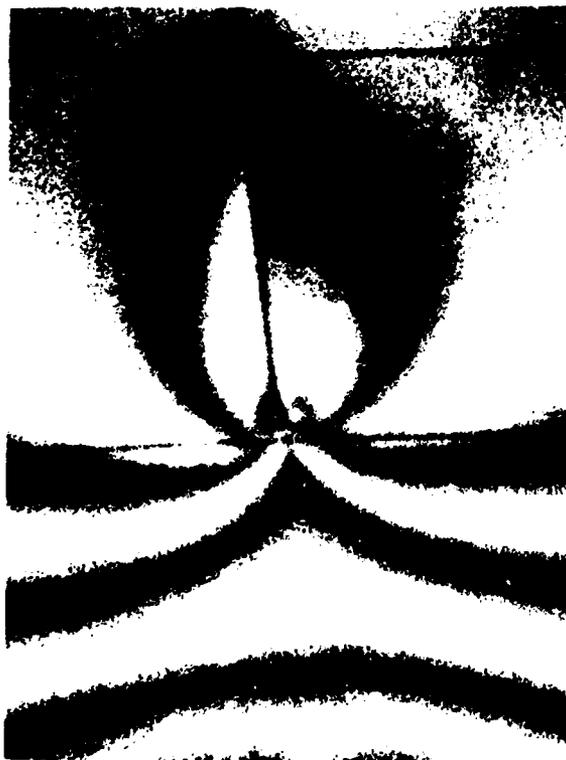
$$K_{II}/K_I < 0 \quad K_{II}/K_I < 0 \quad K_{II}/K_I < 0 \quad K_{II}/K_I > 0 \quad K_{II}/K_I > 0 \quad K_{II}/K_I > 0$$

$$\sigma_{ox} < 0 \quad \sigma_{ox} > 0 \quad \sigma_{ox} = 0 \quad \sigma_{ox} = 0 \quad \sigma_{ox} < 0 \quad \sigma_{ox} > 0$$

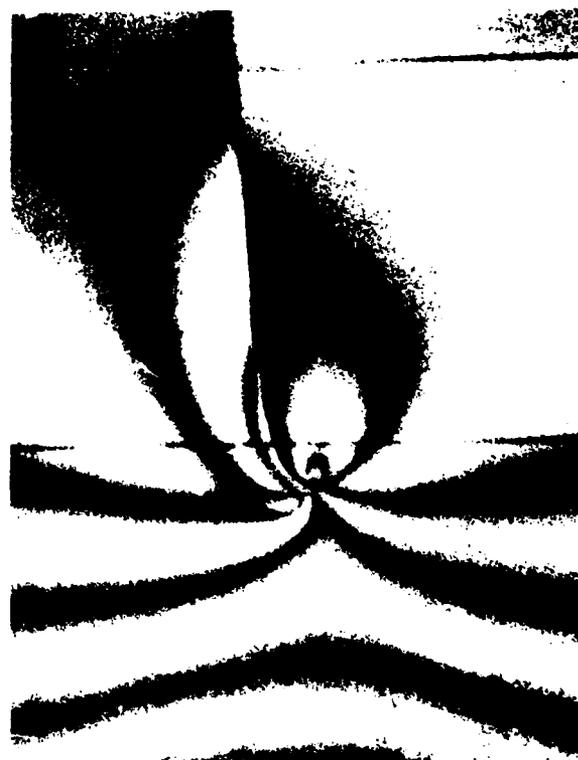
FIG. 4. EXPECTED FRACTURE PATHS BY DYNAMIC PHOTOELASTICITY.



(a) FIFTH FRAME 100  $\mu$  SECONDS



(b) EIGHTH FRAME 130  $\mu$  SECONDS



(c) TENTH FRAME 160  $\mu$  SECONDS

FIG. 5 . TYPICAL DYNAMIC ISOCHROMATICS OF A CURVED CRACK HOMOLITE-100 NOTCH BEND SPECIMEN NO. 6-C051074.

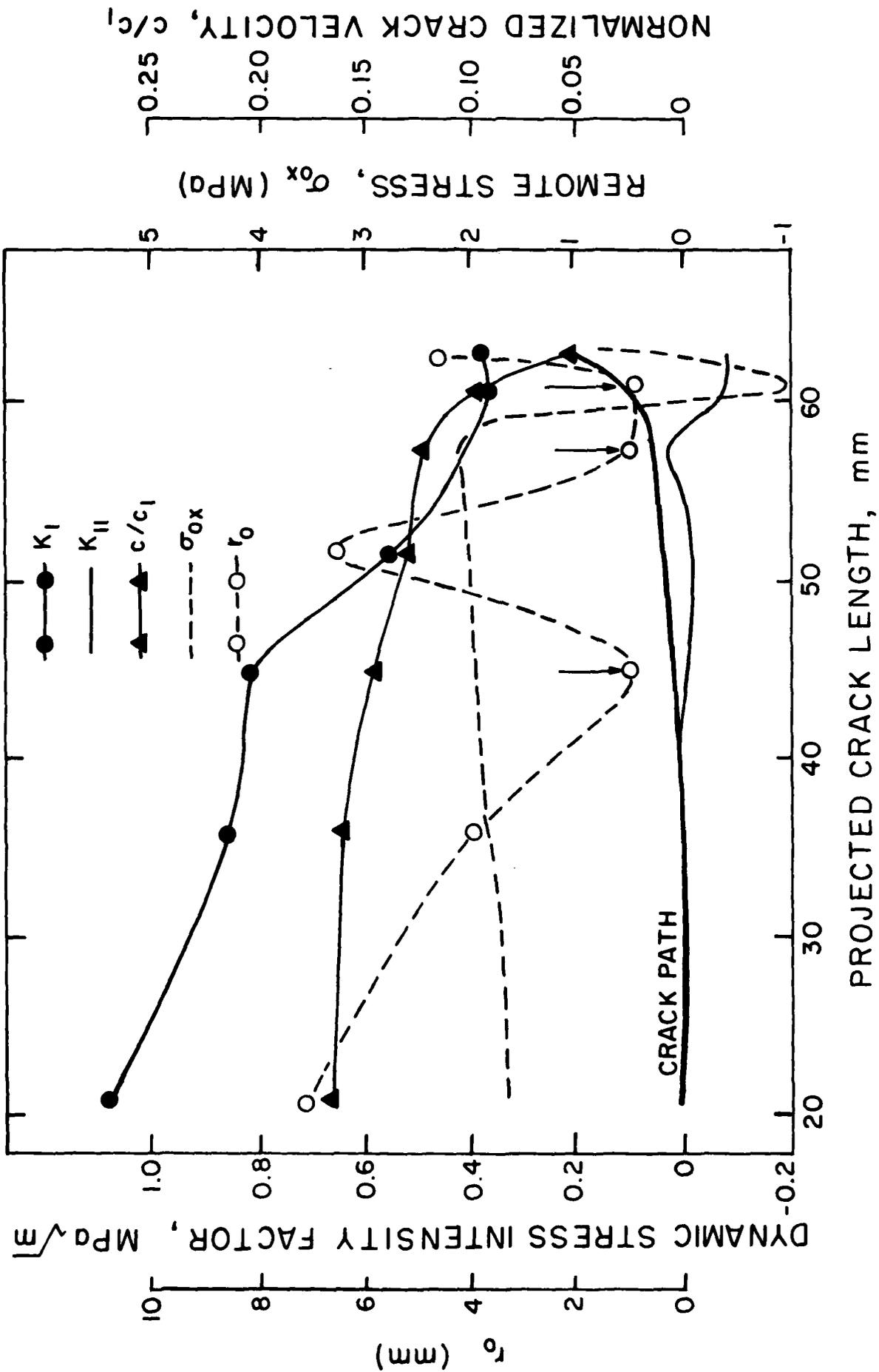


FIG. 6 . MODE I AND MODE II DYNAMIC STRESS INTENSITY FACTORS OF THE CURVED CRACK SHOWN IN FIG. 5 .

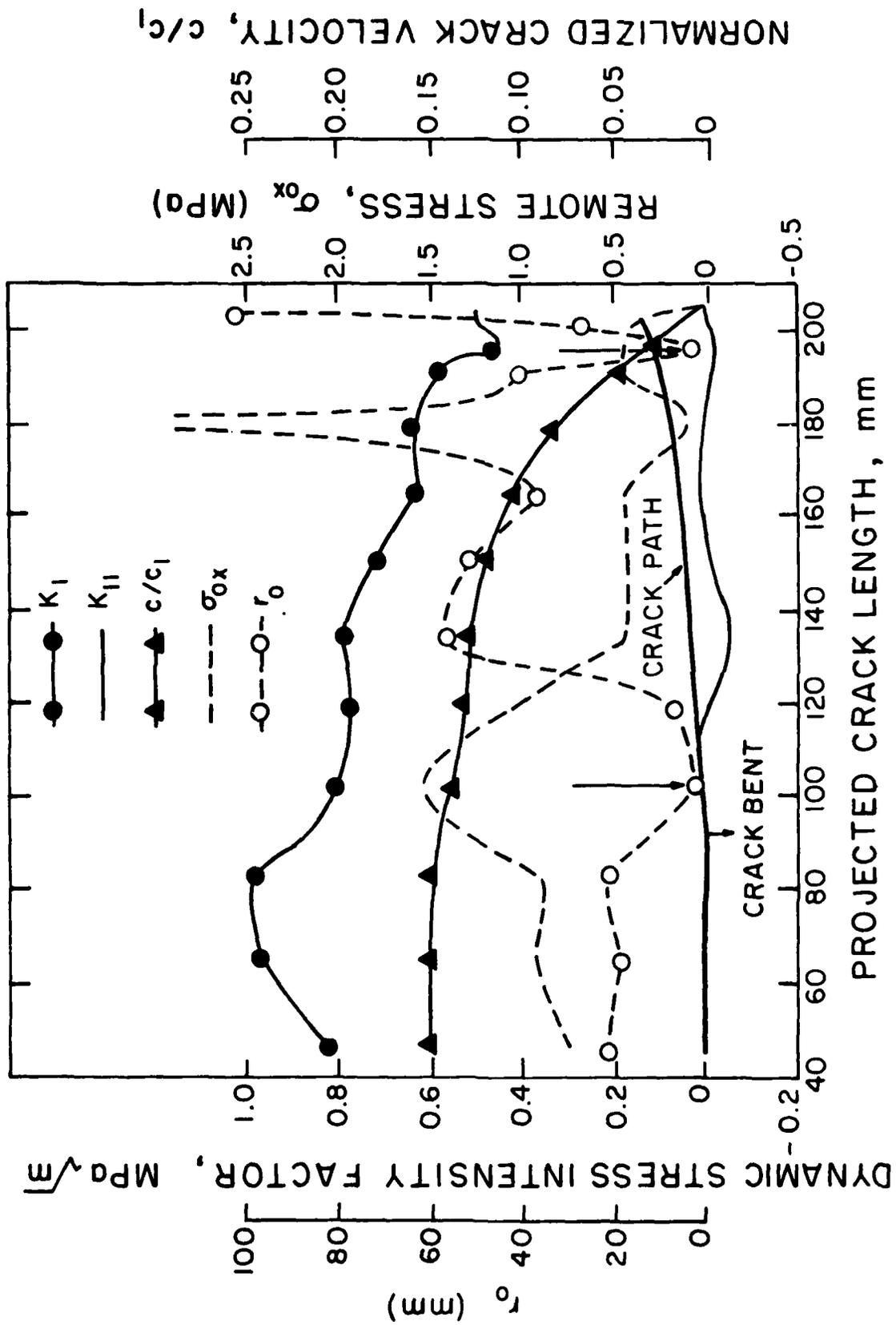


FIG. 7. MODE I AND MODE II DYNAMIC STRESS INTENSITY FACTORS OF SLIGHTLY CURVED CRACK IN A SINGLE-EDGED-NOTCH (SEN) TENSION PLATE. HOMALITE-100, SPECIMEN NO. B12.

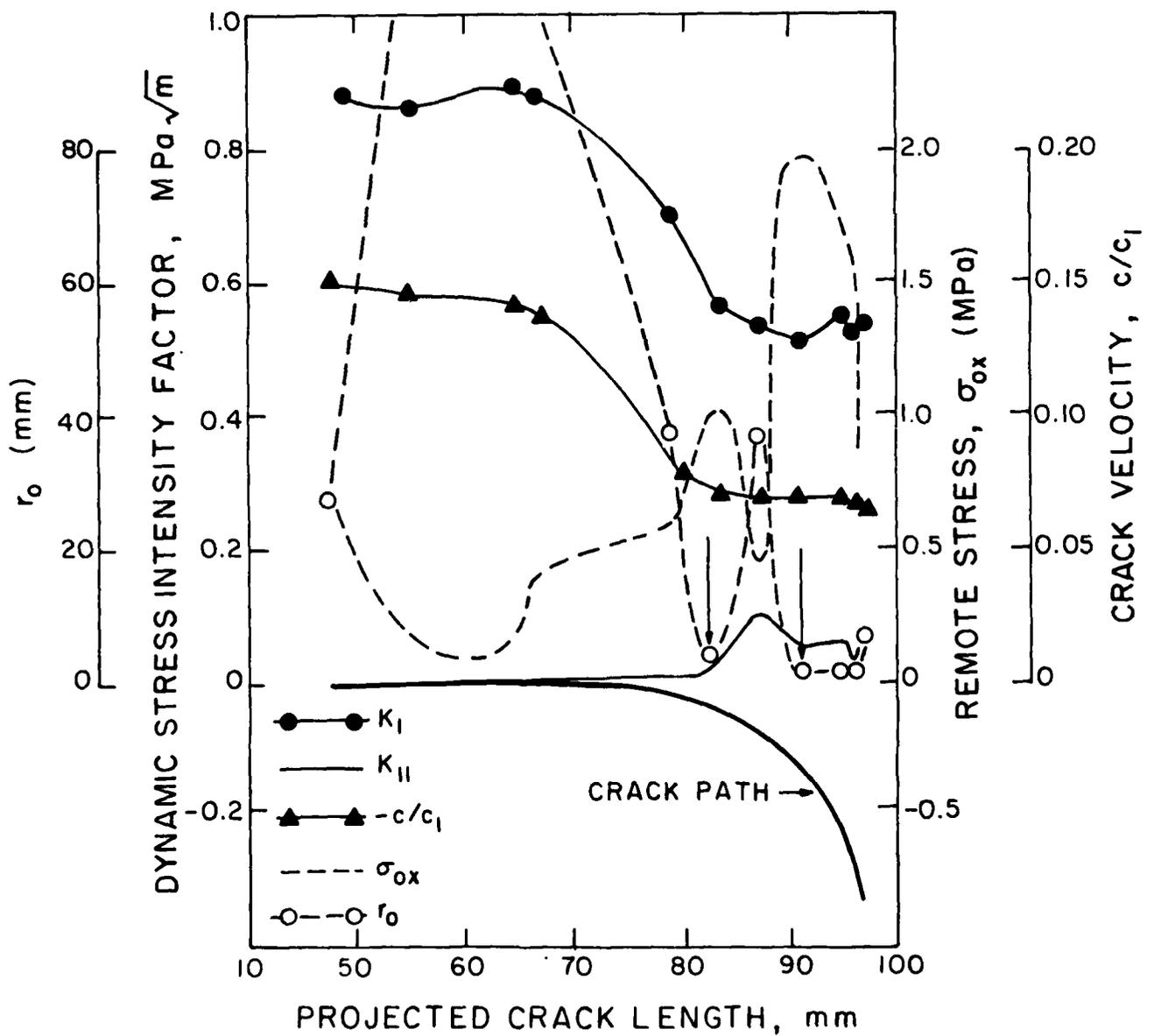
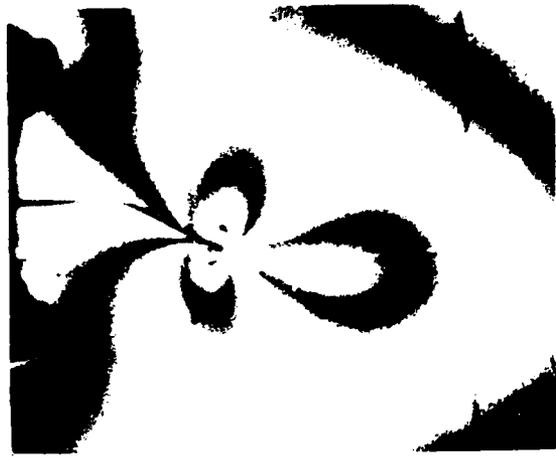


FIG. 8. MODE I AND MODE II DYNAMIC STRESS INTENSITY FACTORS OF A CURVED CRACK IN A WEDGE LOADED RECTANGULAR DOUBLE CANTILEVER SPECIMEN, HOMOLITE -100, SPECIMEN NO. L7B-051573.



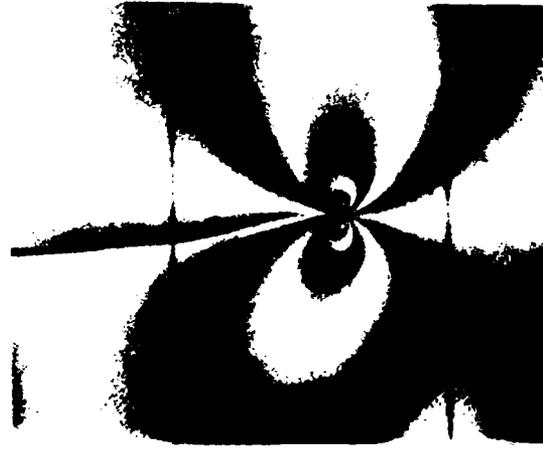
(a) SEVENTH FRAME 150  $\mu$ SECONDS



(b) TENTH FRAME 255  $\mu$ SECONDS



(c) TWELFTH FRAME 315  $\mu$ SECONDS



(d) FOURTEENTH FRAME 370  $\mu$ SECONDS



(e) FIFTEENTH FRAME 390  $\mu$ SECONDS

FIG. 9 . TYPICAL DYNAMIC ISOCHROMATICS OF A CURVED CRACK . HOMALITE-100  
EDGE - CRACKED TENSION PLATE IMPACTED BY A FLAT NOSE PROJECTILE ,  
SPECIMEN NO. 21 - W090771 .

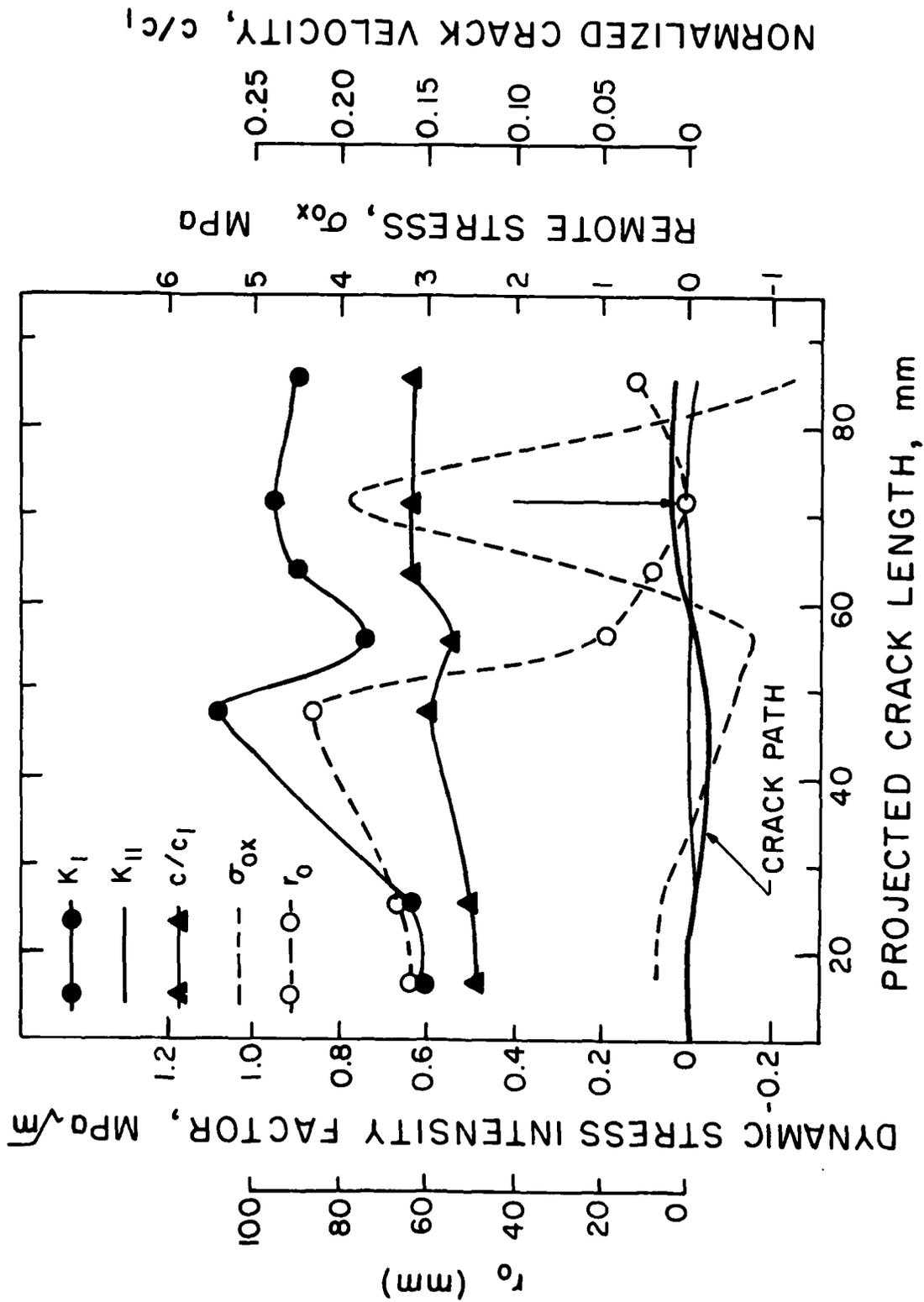


FIG.10 .MODE I AND II DYNAMIC STRESS INTENSITY FACTORS OF A CURVED CRACK IN AN ANGLE - EDGED - NOTCH TENSION PLATE IMPACTED BY A FLAT NOSE PROJECTILE. HOMALITE -100 SPECIMEN NO. 21-WO90771.

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stability of a mode I crack propagation is attained when  $r_0 > r_c$ , where  $r_c$  for Homalite-100 was determined from dynamic photoelastic experiments.

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