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CONTRACTOR REPORT ARSCD-CR-78009

THERMAL ANALYSIS OF FOLDED AMMUNITION

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CHRISTINE M. EASTBURN

JANUARY 1979

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US ARMY ARMAMENT RESEARCH AND DEVELOPMENT COMMAND
SMALL CALIBER
WEAPON SYSTEMS LABORATORY
DOVER, NEW JERSEY

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INTRODUCTION

The objective of this investigation was to determine the temperature response of the "U-shaped" folded cartridge. Figure 1 presents a comparison of the conventional and folded cartridges. The interpose, or web region between the projectile and propellant reservoir, is of particular interest. Because of the unique geometry of the folded cartridge, this region of the weapon will experience heating on both sides, and the possibility of excessive temperatures or structural distortion exists.

To determine the temperature response of the interpose region, the heat input during a typical firing of 100 msec had to be established. In addition, analytical models describing the heat transfer of the folded cartridge geometry had to be developed. These models, one and two-dimensional, describe the transient temperature distribution using a finite difference technique. The models have been checked numerically to insure that nodal size provides an accurate representation. Since a time period longer than 100 msec is required for the weapon's structure to experience temperature change, the analytical model includes multiple firing bursts.

AMMUNITION / WEAPON SYSTEM COMPARISON

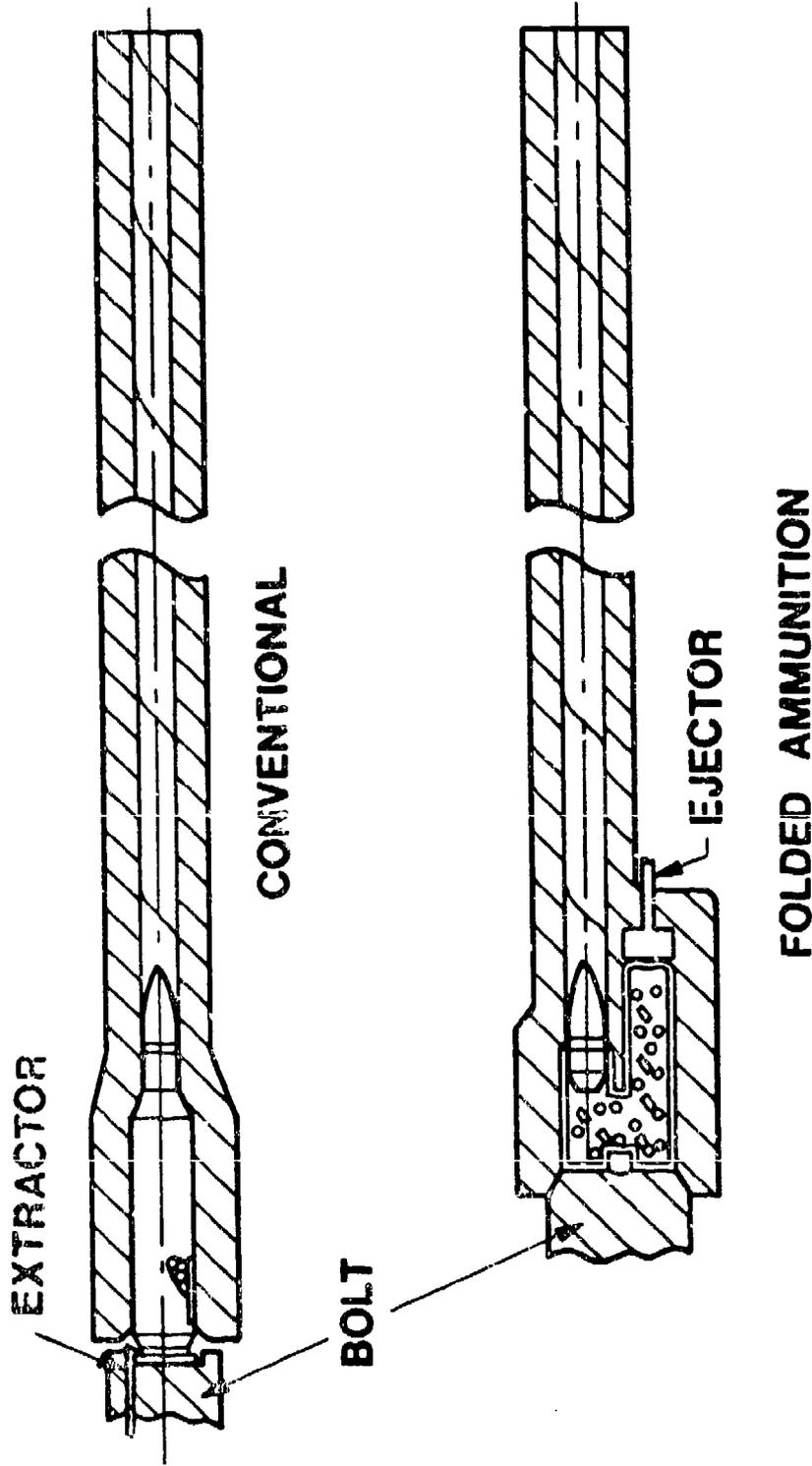


Figure 1. Geometry of folded ammunition.

HEATING ANALYSIS

The heat input cycle is 100 msec in duration. This period includes the firing of one cartridge, a dwell period, the extraction of the spent cartridge, and the insertion of a new cartridge. The highest heat fluxes occur during the firing time, or ballistic portion of the cycle. During this period the cartridge case is heated to its maximum temperature. Following this ballistic period, the energy deposited in the cartridge case dissipates into the weapon. To determine the heat transfer coefficient and gas temperature during the initial period, the local gas pressure, temperature, and velocity must be determined using the applicable ballistic model.

The ballistic model used in this study was supplied by U.S. ARRADCOM. This model, in the form of a computer code, is described in detail in reference 1. The code was developed for standard shape ammunition. The code, as supplied, has been modified to incorporate methods for determining thermochemical properties (ref. 2). Figure 2 presents the space mean pressure history for the folded cartridge calculated with the computer code.

As stated above, the model was not developed for ammunition of the present geometry; however, a review of reference 1 indicates that the model is insensitive to shape. The model does not provide for flame propagation through the propellant bed. Nonsimultaneous ignition is provided by the surface ignition function, which is equal to the total propellant surface area divided by the ignition time range. The latter quantity, ignition time range, is determined by experiment. Since the test case for the folded cartridge compares favorably with test firings, it appears that the ignition time range currently in the program is satisfactory. The three major pressures calculated by the program are the space mean pressure, the breech pressure, and the projectile base pressure. Both the breech and projectile base pressures are constant functions of the space mean pressure. All three pressures are made equal prior to movement of the projectile (approximately 0.34 msec for the present case).

The test case considers the breech to be 9.14 in. from the projectile. This corresponds to a straight cartridge 9.14 in. long from the base of the projectile to the end of the cartridge. A test case with the breech located 1.5 in. from the projectile was run. This corresponds to the folded geometry. The resulting pressures, temperatures, and muzzle velocities were not affected. It appears, therefore, that the only significant parameter is the propellant volume. The gas temperature and velocity are uniform over the entire region behind the projectile. This is significant because it indicates uniform heating.

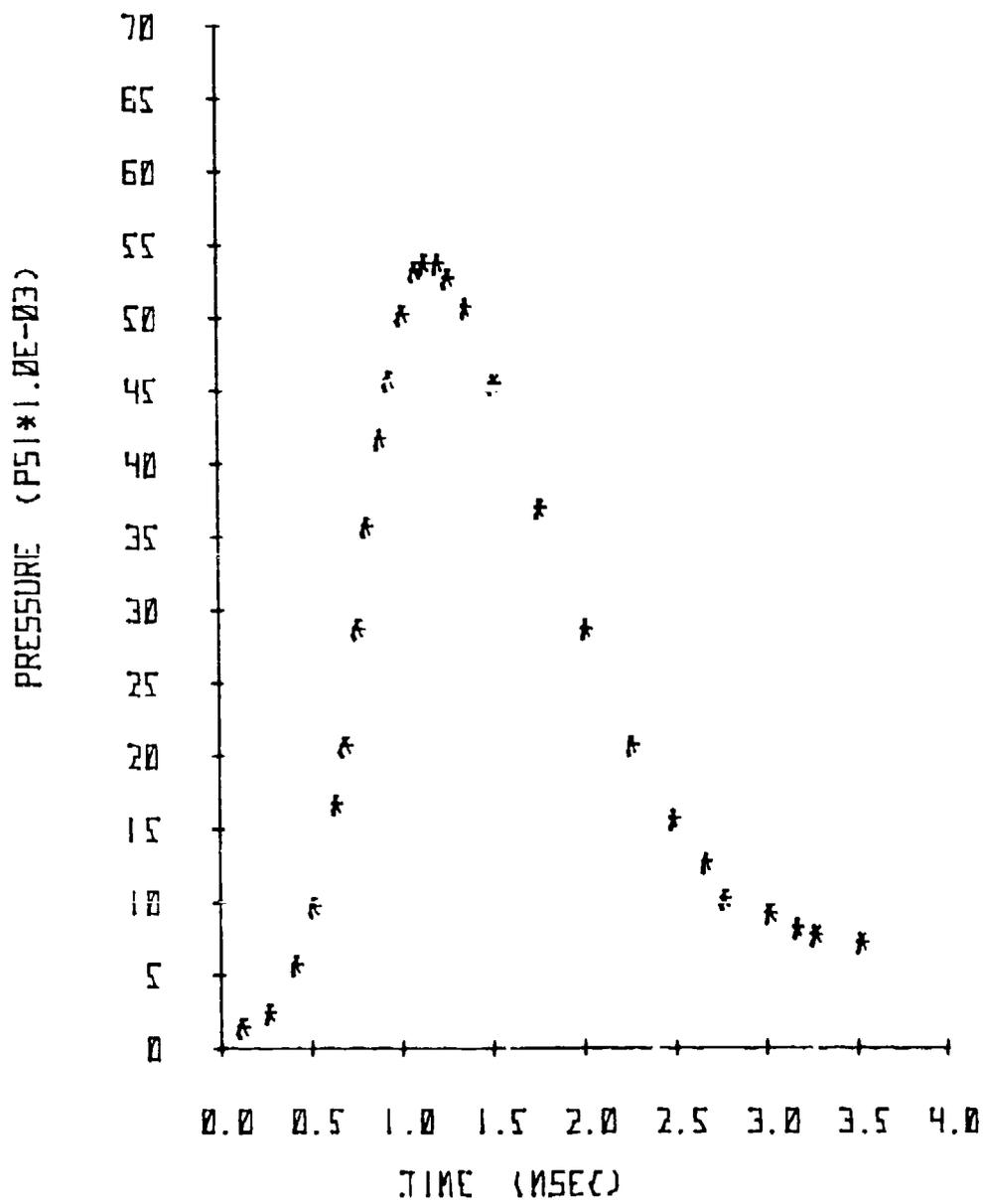


Figure 2. Space mean pressure versus time.

To accurately model nonuniform heating, a new ballistic model would have to be developed. This model would have to consider a one-dimensional transient flow allowing for finite flame propagation. This model could correctly position the initiation point and yield space dependent velocities, temperatures, and pressures. The decision regarding whether such a model should be developed will depend upon the need for improved ballistic data and the magnitude of the heating problem. However, the decision was made to use the existing model to establish the heating conditions in the folded geometry.

The flow conditions and temperature levels in the folded cartridge pose a severe heating environment for the cartridge case and surrounding weapon. Prediction of the heat transfer coefficient is also difficult because of the unique operating conditions.

Reference 3 provided heat transfer coefficients and temperature data for an environment similar to the present case. Unfortunately, the method used to make the calculations is not given. The maximum heat transfer coefficient is approximately 35×10^3 Btu/ft² hr⁰ F. In addition, the maximum adiabatic wall temperature is 5700⁰ F. This is considerably higher than the value of 4065⁰ F obtained with the ballistic code, and it is also higher than the adiabatic flame temperature of the propellant.

Reference 4 was reviewed as a possible source of a heat transfer coefficient prediction method. This report examined a variety of possible prediction methods, but all required more flow field information than presently available. In addition, the report was more involved with the gun barrel, rather than the breech area. Some of the data presented indicated heat transfer coefficients in the range from 30×10^3 Btu/ft² hr⁰ F to 70×10^3 Btu/ft² hr⁰ F. It should be noted that these extreme levels exist for only 1 or 2 msec. Reference 5 evaluated the results of a method of characteristics study of the interior ballistics problem. The Colburn analogy, which is given by:

$$St Pr^{2/3} = C_f/2$$

where,

St = Stanton Number = Nu/Re Pr
Nu = Nusselt Number
Pr = Prandtl Number
Re = Reynolds Number
Cf = Friction factor

was used. The resulting heat transfer coefficients were in the same range as the previous references.

Reference 6 considered the problem of flow field development in a solid rocket motor. The operating conditions and dimensions were similar to the folded cartridge problem. This reference recommends

the Dittus Boelter correlation for local heat transfer coefficient. The Dittus Boelter equation is

$$Nu = 0.023 Re^{.8} Pr^{.4}$$

The above equation is widely used in turbulent heat transfer calculations, and is relatively easy to apply to the present case. The decision was therefore made to use this relationship.

Assuming the gas to be ideal, and that the viscosity varies with the temperature to the 0.65 power, the heat transfer coefficient is given by:

$$h = 0.023 \left(\frac{PV}{R} \right)^{.8} C_p Pr^{-.6} \left(\frac{\mu_0}{D(530)^{.65}} \right)^2 T^{-.67}$$

where,

- P = pressure
- V = gas velocity
- C_p = specific heat
- μ₀ = viscosity at 530° R
- D = diameter of chamber.

For the present case, the various constants are given by:

$$\mu_0 = .87 \times 10^{-6} M^{.5} (530)^{.65} \quad (\text{ref. 6})$$

$$R = \frac{F}{T_0} = \frac{\text{propellant impetus}}{\text{flame temperature}} \quad (\text{ref. 2})$$

$$M = R_0/R$$

Using the thermochemical properties for the folded cartridge propellant and the ballistic data from the computer code, the heat transfer coefficients can be determined as a function of time. Figure 3 presents the variation of heat transfer coefficient versus time, based upon space mean pressure. Figure 4 presents the same data based upon projectile base pressure. For space mean pressure conditions, the peak heat transfer coefficient is approximately 52,500 Btu/ft²hr⁰F. This method for calculating heat transfer coefficient has been incorporated into the ballistic code. In addition, the stagnation temperature, which will be used in computing the convective heat inputs, has also been determined and is represented graphically in figure 5. It is worth noting that the maximum temperature in figure 5 is below that given in reference 3. This aspect of the results in reference 3 is difficult to explain. The peak temperatures reported in that reference are higher than the adiabatic flame temperature of the propellant used for the folded ammunition.

The heat transfer coefficients calculated with the Dittus Boelter equation are higher than those cited in reference 3, but are

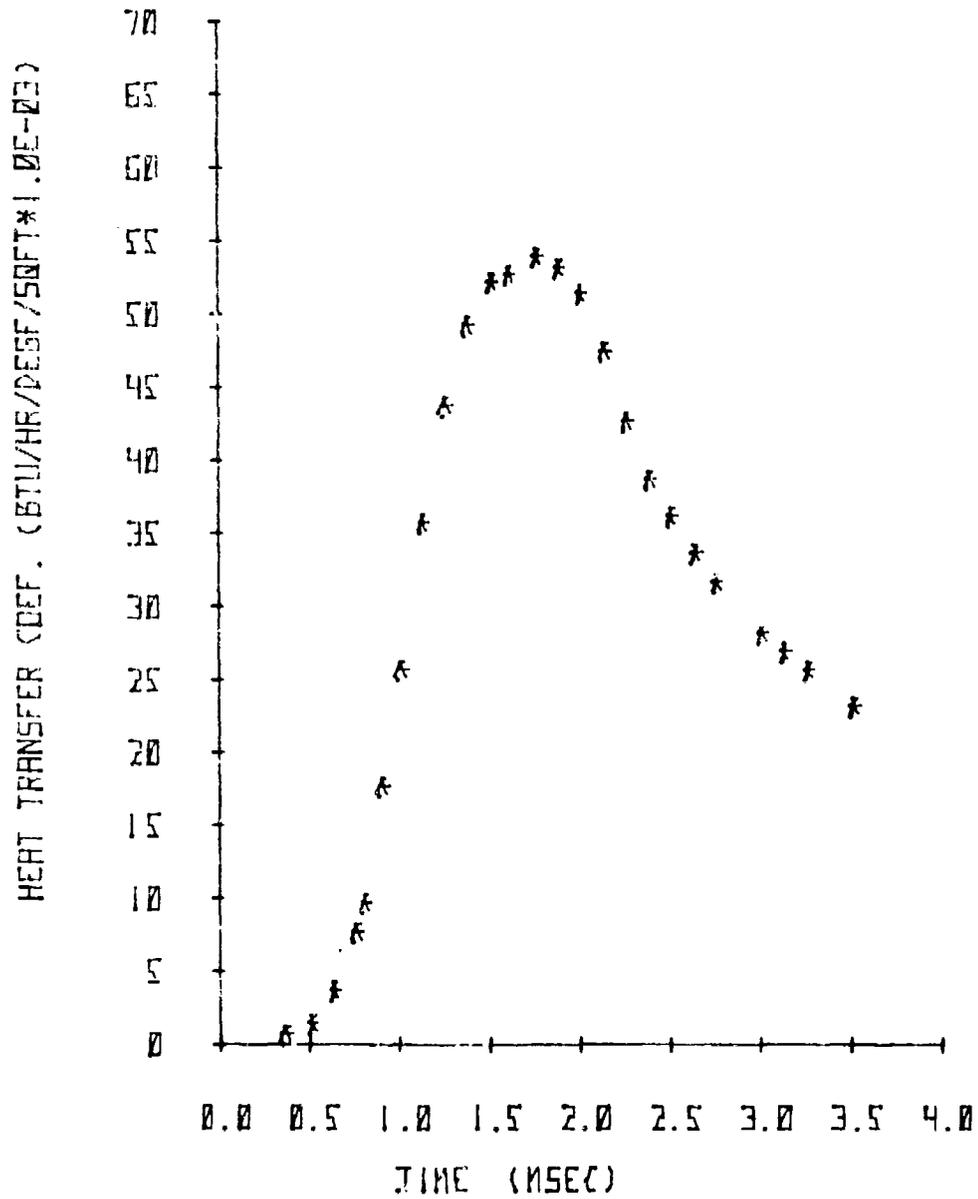


Figure 3. Heat transfer coefficient versus time - based upon space mech pressure.

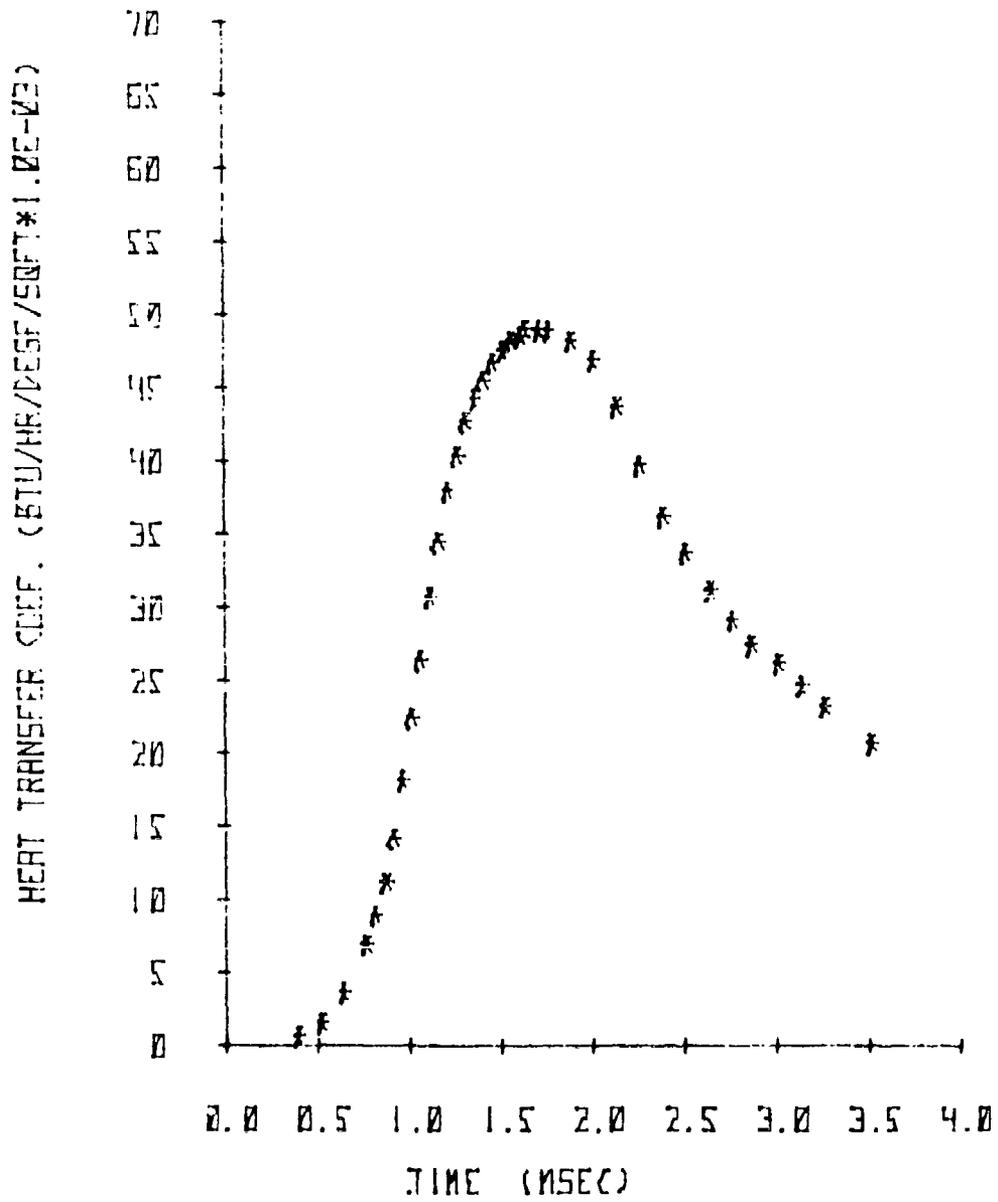


Figure 4. Heat transfer coefficient versus time - based upon projectile base pressure.

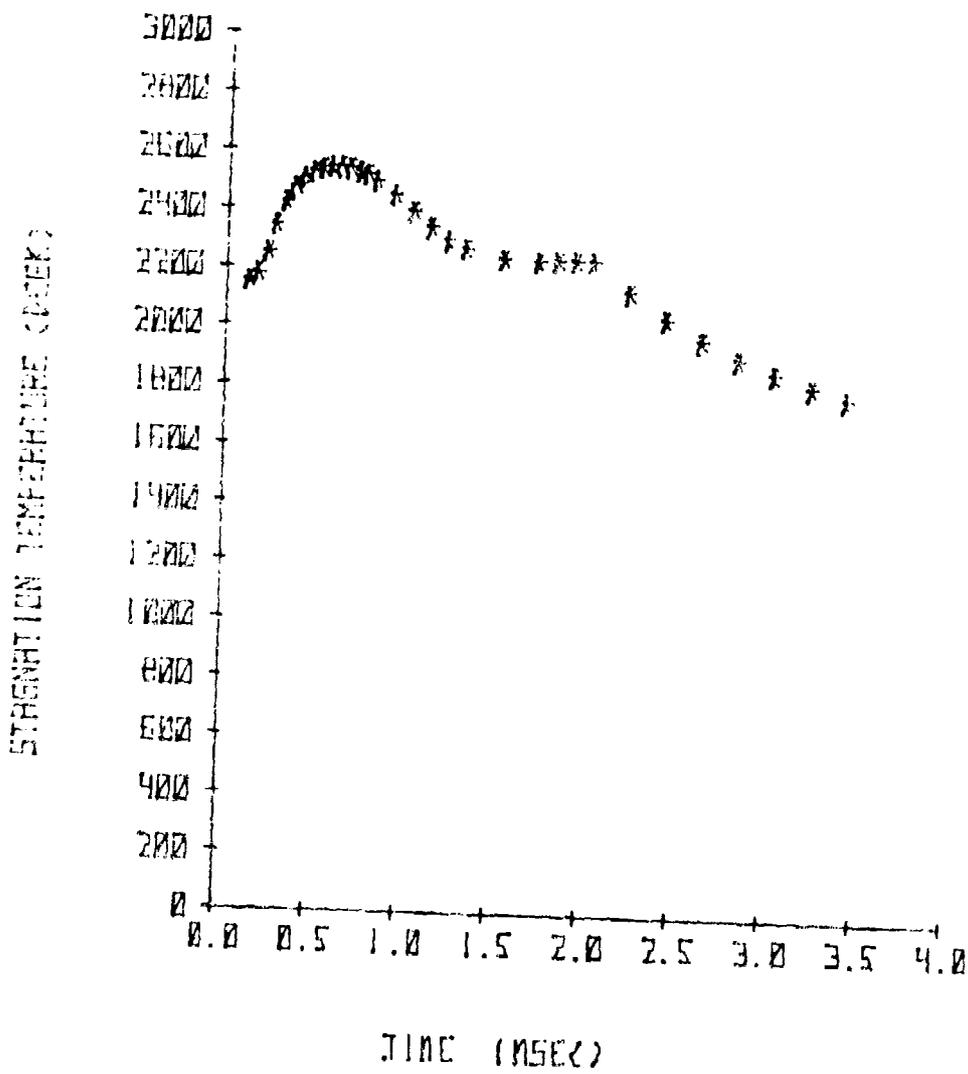


Figure 5. Stagnation temperature versus time.

well within the range reported by other references.

The heating conditions presented in figures 3 and 5 occur during the initial ballistic portion of the cycle, and last for approximately 3.5 msec. The peak surface temperature of the cartridge occurs during this portion of the cycle. However, the total time for the firing of a single cartridge is approximately 100 msec. It is during this longer period that the heat permeates into the weapon's structure and causes the highest structural temperatures.

The time history of a firing cycle is composed of the following four periods:

<u>Time</u>	<u>Event</u>
0 to 3.5 msec	Firing of projectile (ballistic portion)
3.5 to 35 msec	Dwell time
35 to 65 msec	Extraction of expended cartridge
65 to 100 msec	Loading of unspent cartridge

For the preceding firing cycle, the following heat inputs have been developed:

<u>Time</u>	<u>Heat Inputs</u>
0 to 3.5 msec	Convective heat input based upon Dittus Boelter equation (ref. 1). Figures 3 and 5 present the heat transfer coefficient and gas temperature profile.
3.5 to 35 msec	Exponential decay of the heat transfer coefficient to 100 Btu/ft ² ·hr ⁰ ·F and of the gas temperature to 200° F.
35 to 65 msec	Constant heat transfer coefficient and gas temperature at 100 Btu/ft ² ·hr ⁰ ·F and 200° F, respectively.
65 to 100 msec	Zero external heating. In addition, the cartridge nodes are returned to 100° F at 65 msec and then allowed to heat up due to the heat flow from the structure.

During the dwell time (3.5 to 35 msec), exponential decay of the gas pressure, gas temperature, and gas velocity is felt to be a reasonable approximation following the pressure decay. The value of 100 Btu/ft²·hr⁰·F for the heat transfer coefficient at 35 msec is conservative. The form of the heating boundary condition, during extraction and loading, is very difficult to describe. During future studies, experimental data would be examined to provide better insight into the thermal conditions during these phases of the operating cycle. For the present effort, the decision was made to use a conservative

approach and combine extraction and loading into a single instantaneous event at 65 msec. At 65 msec, the heat transfer coefficient was set at zero, and the nodes representing the cartridge case were set to 100° F. This simulates the chambering of an unspent cartridge. During the period from 65 to 100 msec, the cartridge case is cooler than the structure, and the heat flows into the cartridge.

Although the heat input profile used for this study is considered to be conservative, it should be examined further to determine the significance of the various assumptions made. This activity would be an interesting area for further effort.

ANALYTICAL MODELS

The detailed geometry of the folded cartridge case is given in figure 6. Since the heating during firing is uniform, based upon the present model, the problem reduces to the determination of two-dimensional temperature distributions. However, the results presented in reference 3 indicate that the depth of penetration of heat was small; therefore, the heat flow was essentially one-dimensional. Initial studies during this investigation verify this result. The majority of the results presented in the report will therefore be based upon one-dimensional models. However, in the web region near the breech, the geometry causes two-dimensional effects to occur. Two-dimensional results were obtained for this region.

One-Dimensional Model

The one-dimensional model consists of a 0.04 in. cartridge case thickness, and 0.50 in. of steel, which represents the weapon structure. The properties of the case material, cartridge brass and steel (SAE 4340), were obtained from reference 3. The temperatures were determined by dividing the model into a number of isothermal nodes, and solving the following equation using a digital computer:

$$T_i^{t+\Delta t} = \frac{\Delta t}{m_i c_i} \left\{ \sum_j k_{ij} T_j^t + h A_i T_o \right\} + \left\{ 1 - \frac{\Delta t}{m_i c_i} (\sum_j k_{ij} + h A_i) \right\} T_i^t$$

where,

- k_{ij} = the conductive coupling between nodes
- $T_i^{t+\Delta t}$ = temperature of the i^{th} body at time $t + \Delta t$
- A_i = surface area of nodes exposed to convective (zero for all but surface node) heat transfer
- h = convective heat transfer coefficient
- T_o = gas temperature.

Appendix A presents a listing of the computer program. In this analysis, both the heat transfer coefficient and the gas temperature are functions of time, as described in the Heating Analysis section of this report.

Two grid systems were initially used in the analysis. These are presented in table 1. The first system contains 26 bodies, with 11 bodies in the cartridge case, while the second contains 31 bodies, with 21 in the cartridge case. The use of more bodies in the cartridge case was initially felt to be necessary because of the rapid temperature changes occurring in the region. However, examination of the results from both grid systems did not yield any significant differences (ref. 7). Since the 26 body model could be run at a larger time step, and therefore require less computer time, it was used for most of the analysis.

Table 1. One-dimensional model

1st Grid (26 bodies)		2nd Grid (31 bodies)	
Node Number	Thickness (in.)	Node Number	Thickness (in.)
1	.002	1,2	.001
(inter- 2 to 10	.004	(inter- 3 to 20	.002
face) 11	.002	face) 21	.004
12	.004	22	.004
13	.008	23	.008
14	.012	24	.012
15	.020	25	.024
16 to 25	.040	26,27,28	.05
26	.054	29,30,31	.1

Two-Dimensional Model

Although one-dimensional modeling is adequate for most of the folded geometry thermal analysis, certain regions require two-dimensional studies. The web, or interpose region near the breech end, will experience two-dimensional heating. To analyze this region, an isothermal model, shown in figure 7, has been developed. This model consists of 62 bodies, with 50 bodies in the cartridge case. Initially a smaller model consisting of 37 bodies, with 25 bodies in cartridge case, was considered. However, nodalization studies showed that this grid was too large. Because of the symmetry in the web, only half of the web was modeled. The rectangular region is 0.23 in. by 0.56 in. and contains the web region up to the base of the projectile.

The computer program used to analyze the two-dimensional problem is essentially the same as the one described in the One-Dimensional Model section. The major difference is that surface and

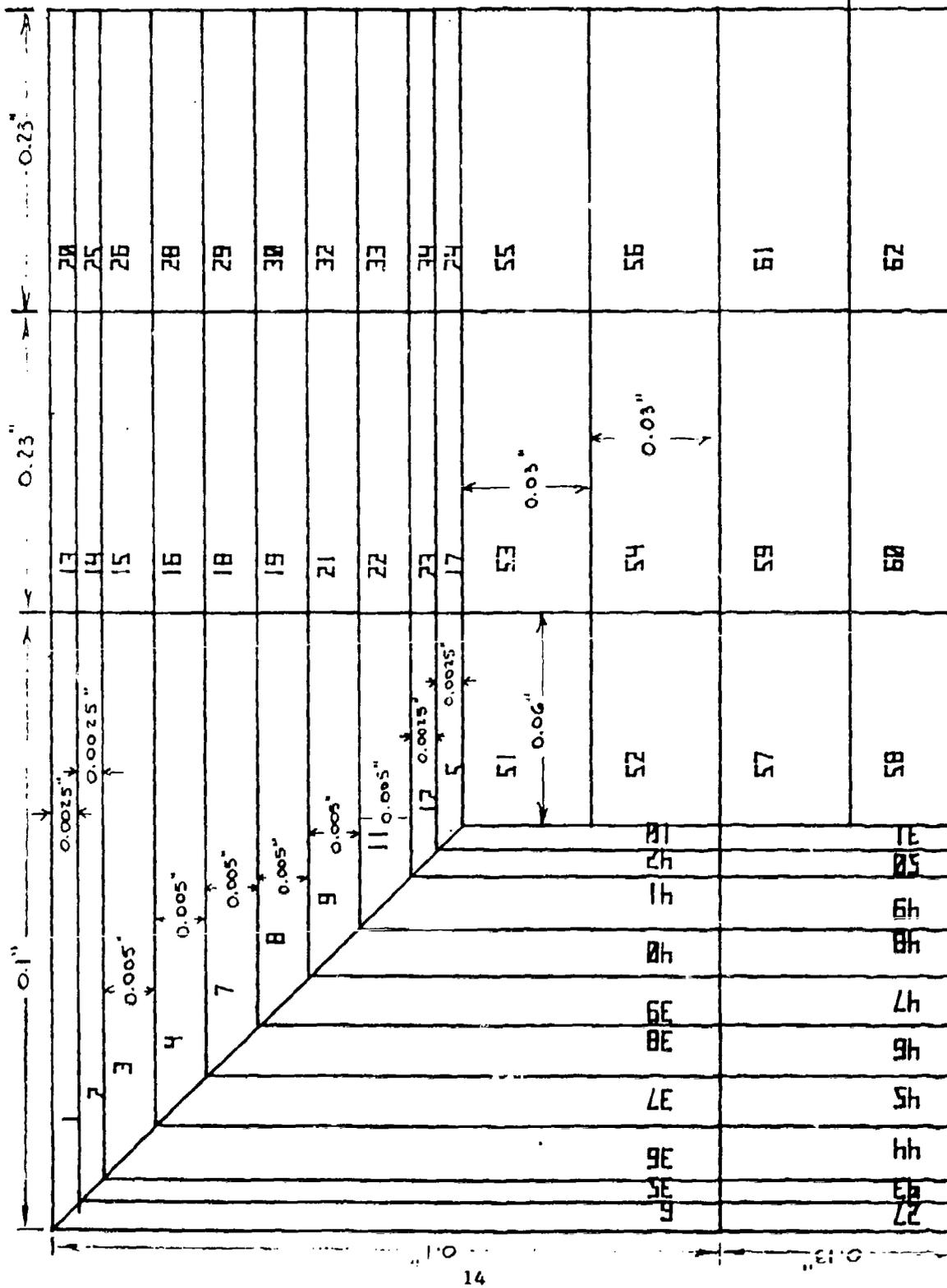


Figure 7. Two dimensional model.

interior nodes are identified, and surface nodes are allowed to experience convection. Appendix B contains the Fortran listing of the program.

RESULTS

One-Dimensional

Using the one-dimensional model described previously, temperature profiles have been determined for 29 cycles or rounds.

During any cycle, the surface of the cartridge exposed to the hot gases experiences a rapid temperature rise, with the peak occurring at approximately 2 msec. Figures 8 through 11 present the surface temperature history for rounds 1, 10, 20, and 29. These profiles are similar, with the peak temperature reaching a maximum value of approximately 1346° F. The sharp drop in temperature, occurring at 65 msec, results from the extraction of the spent round and the insertion of an unspent cartridge with the cartridge case at 100° F. The unspent cartridge is then heated by the warmer weapon chamber. Figure 12 presents a composite of the profile for rounds 1, 10, 20, and 29. From this figure it can be seen that by the tenth round, the surface temperature profile is stabilized. Figure 13 presents the peak surface temperature as a function of round number. It would appear that by round 4, the peak surface temperature has reached equilibrium at approximately 1345° F. However, examination of the detailed computer printouts shows that between rounds 19 and 29, the peak surface temperature rose 0.5° F. This rate, of course, is decreasing with each round, but extrapolation would yield a conservative estimate of the peak temperature after a one minute burst of 600 rounds. Using a rate of 1° F per 20 rounds, the extrapolated peak surface temperature is 1374° F.

During the ballistic portion of any firing cycle, there is very little penetration of the thermal wave into the weapon chamber. Reference 7 states that at the conclusion of the ballistic period (approximately 3.5 msec), the interface between the cartridge case and steel portion of the weapon experienced less than a 1° F rise. Over the 100 msec firing cycle, however, the thermal wave does penetrate into the steel structure. Figures 14 through 17 present the temperature distribution across the model at the end of the cycle for rounds 1, 10, 20, and 29. After 20 rounds, the temperature rise in the steel, at a depth of 0.25 in., is still less than 20° F. Figure 18 presents a composite of figures 14 through 17. From this figure, the propagation of the thermal wave can be seen. The peak temperature shown in figure 18 is not the highest steel temperature, but rather the condition existing at the end of the cycle for a particular round. Figure 19 presents the peak interface temperature as a function of round number. This represents the maximum steel

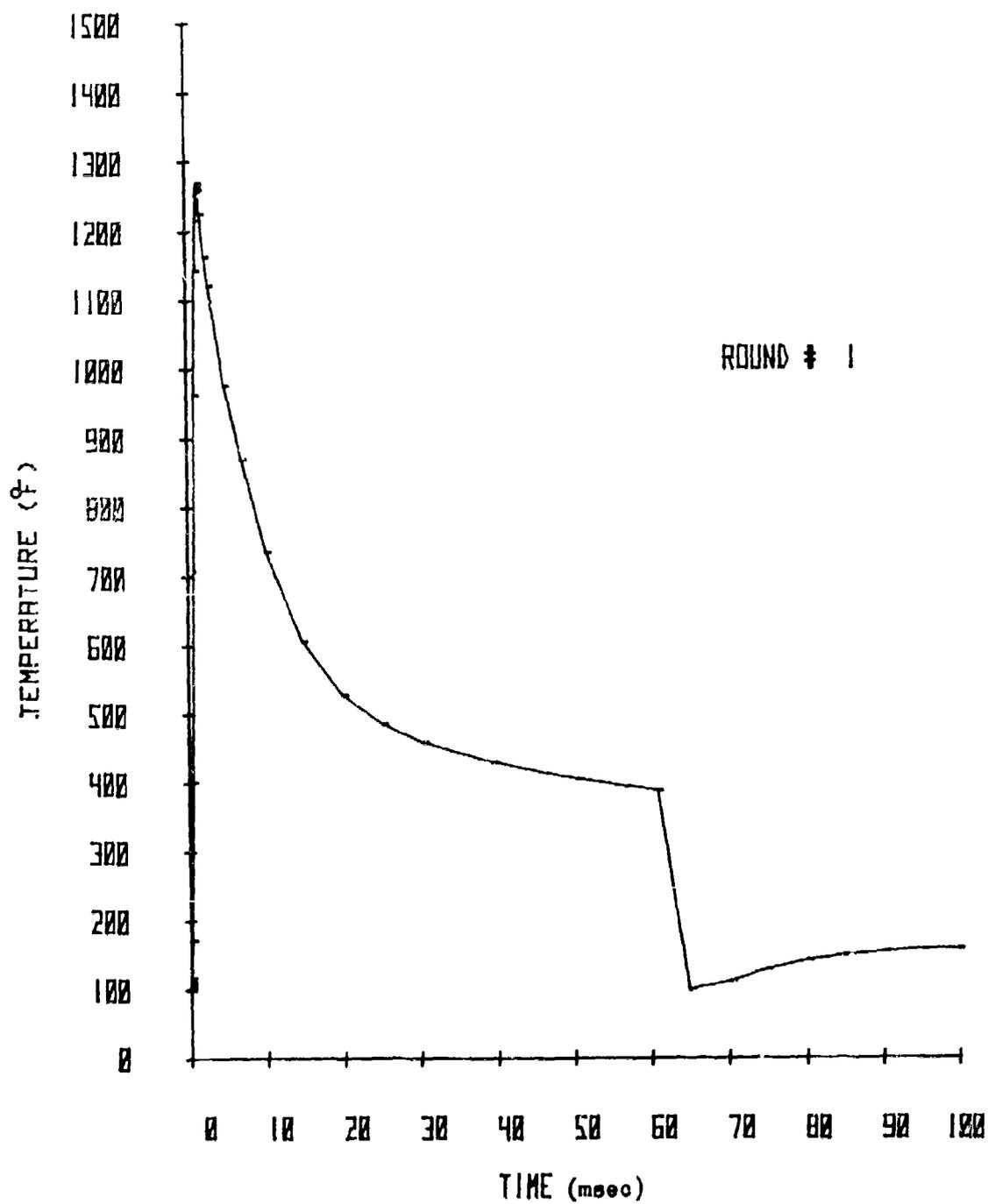


Figure 8. Surface temperature versus time, round 1.

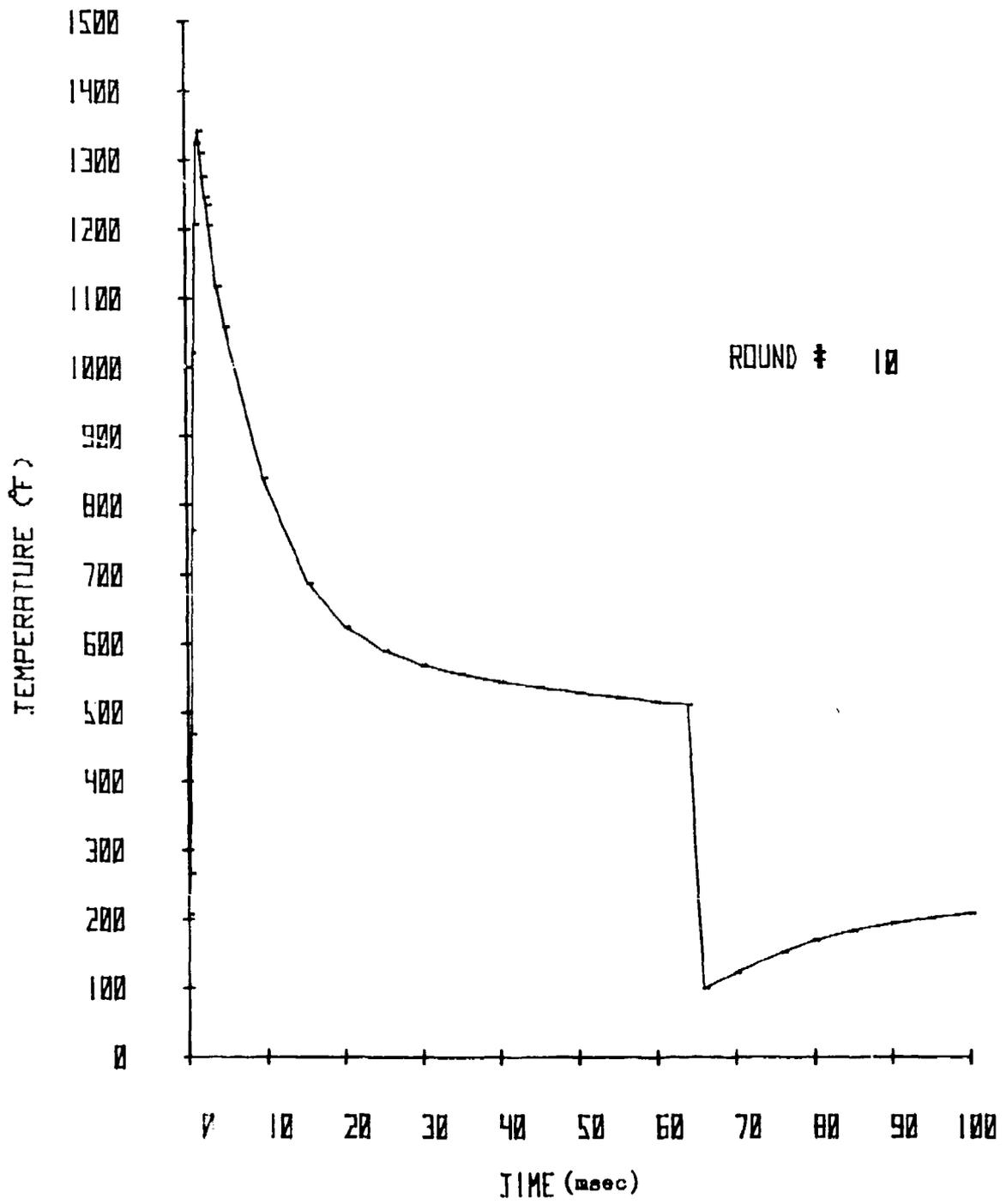


Figure 9. Surface temperature versus time, round 10.

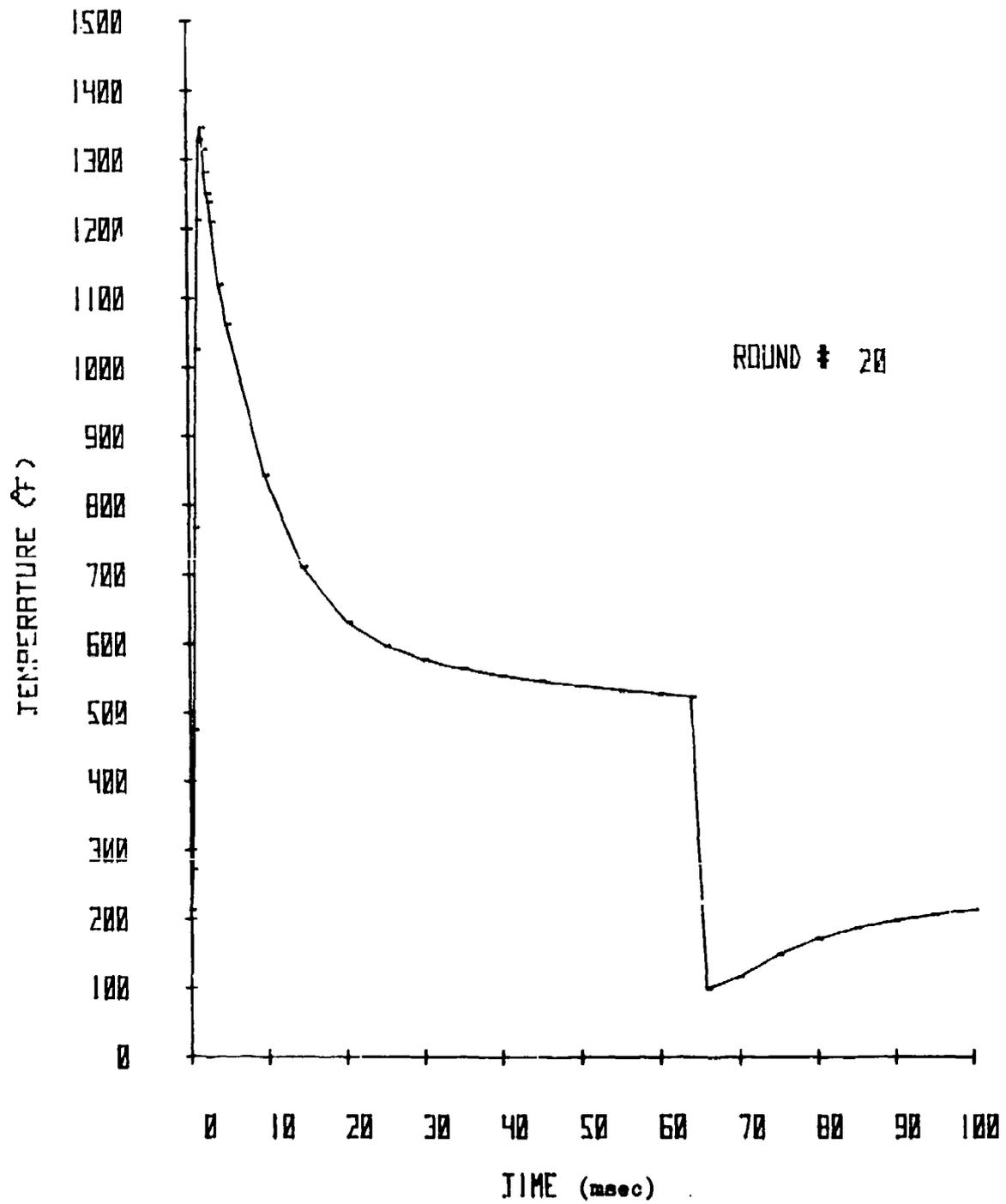


Figure 10. Surface temperature versus time, round 20.

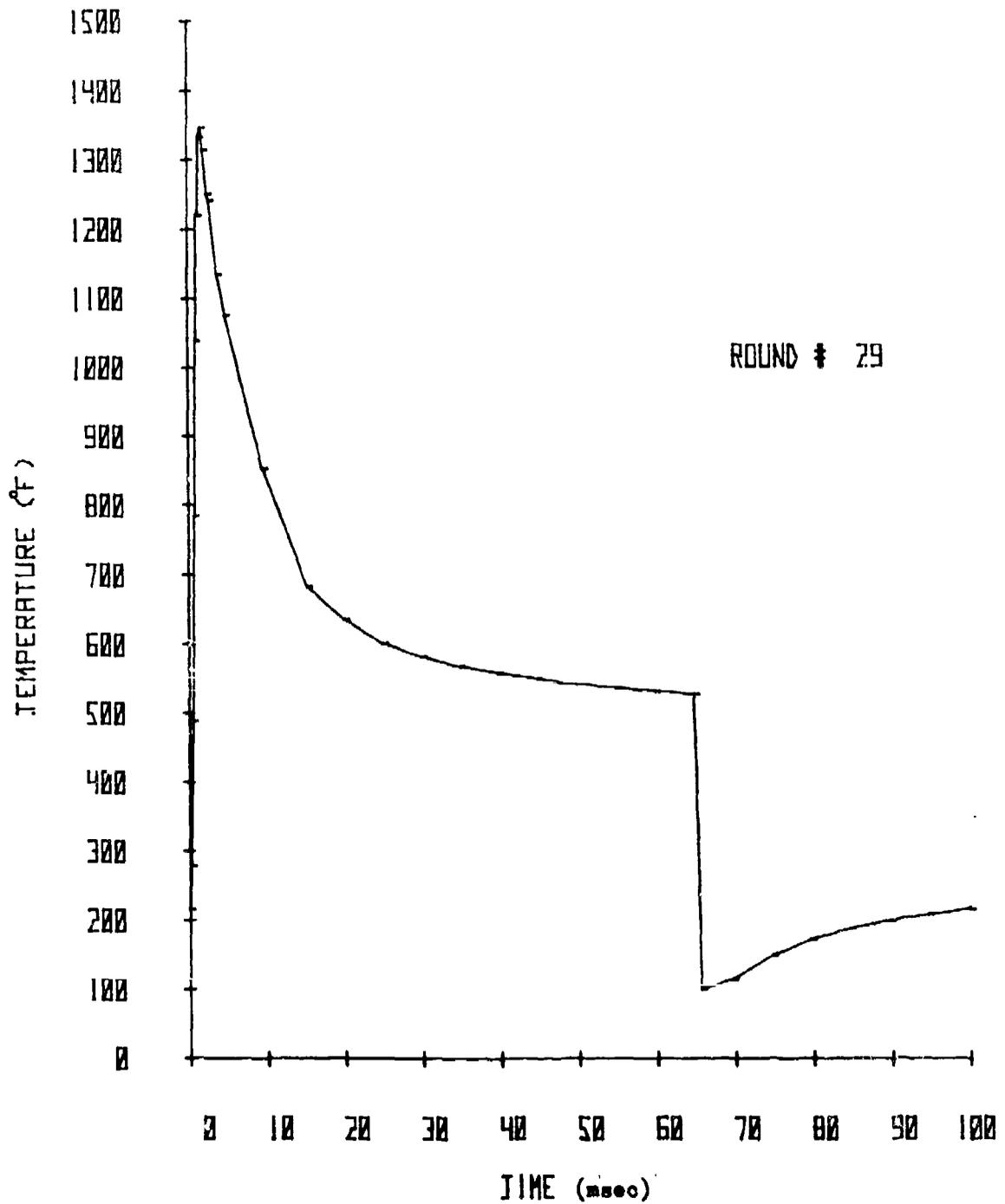


Figure 11. Surface temperature versus time, round 29.

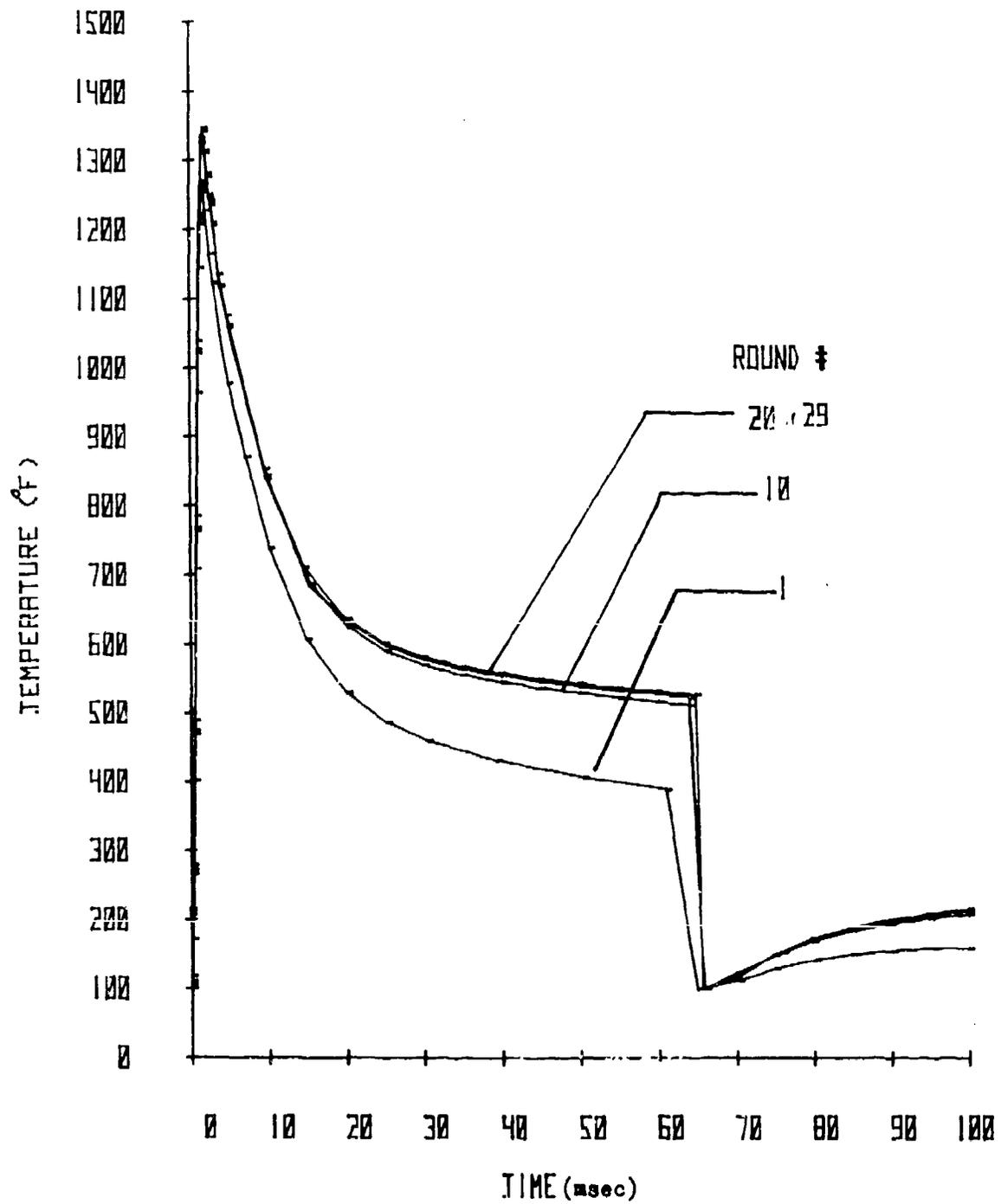


Figure 12. Surface temperature versus time, rounds 1, 10, 20 & 29.

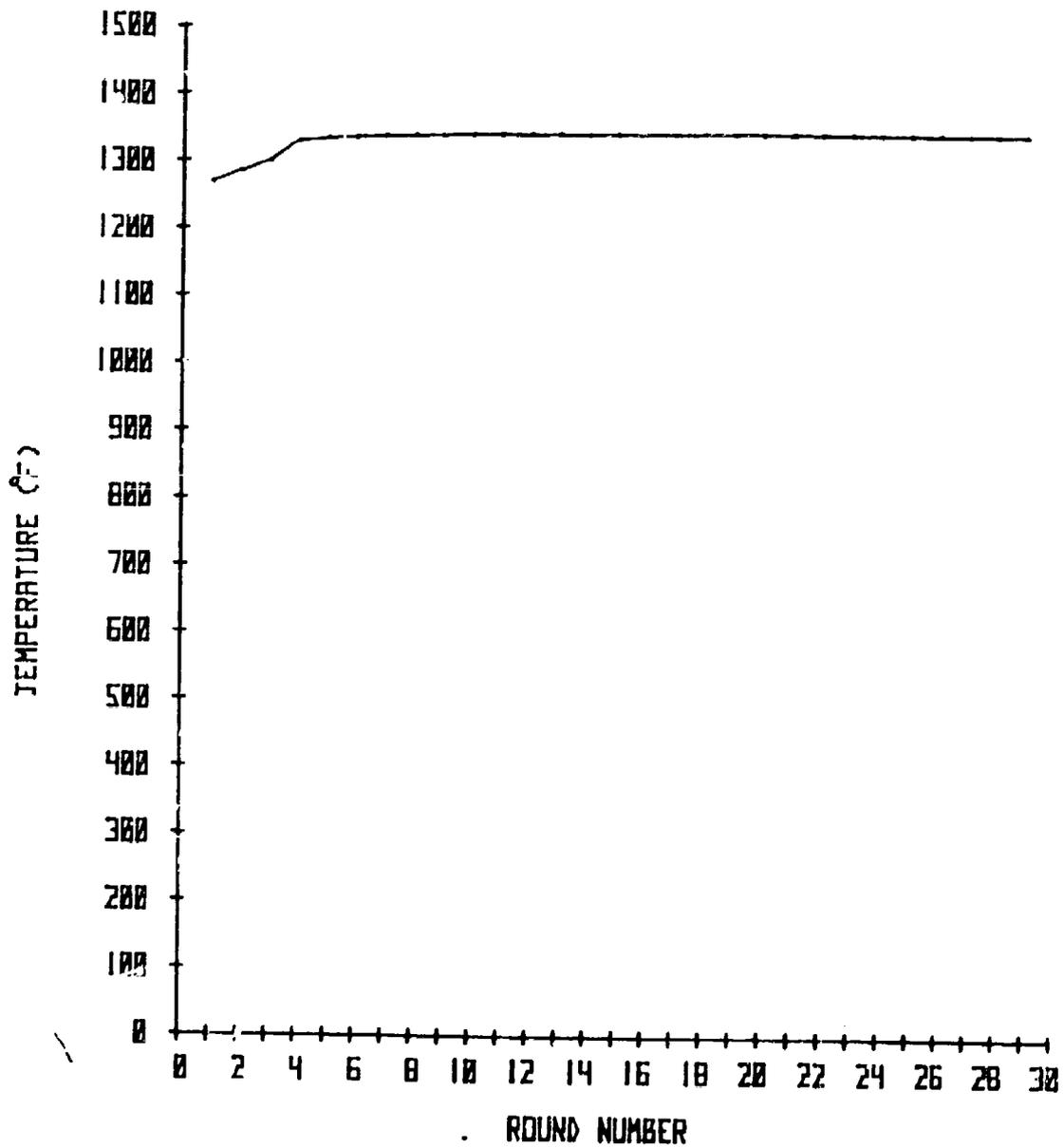


Figure 13. Peak surface temperature versus round number.

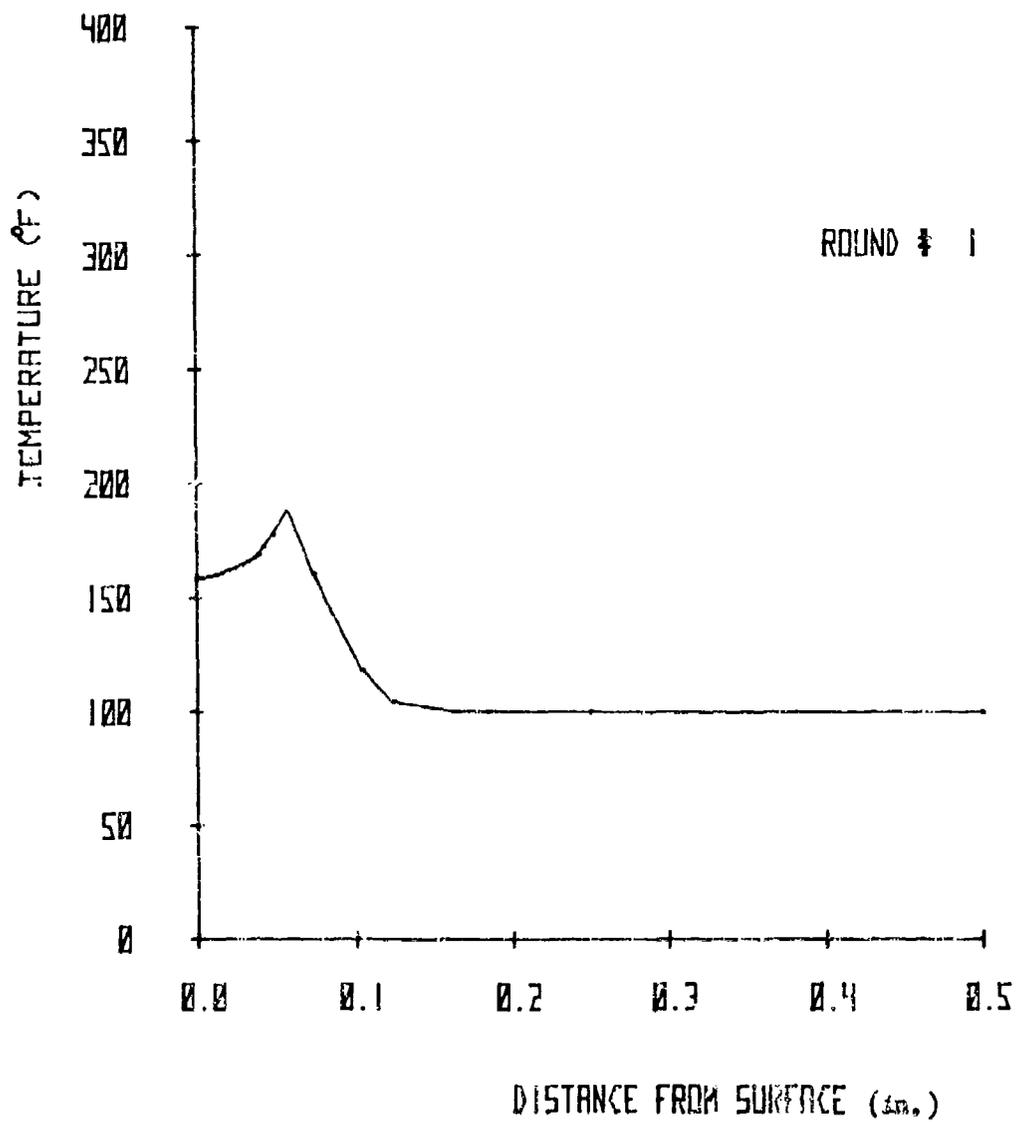


Figure 14. Temperature versus distance at 100 (msec).

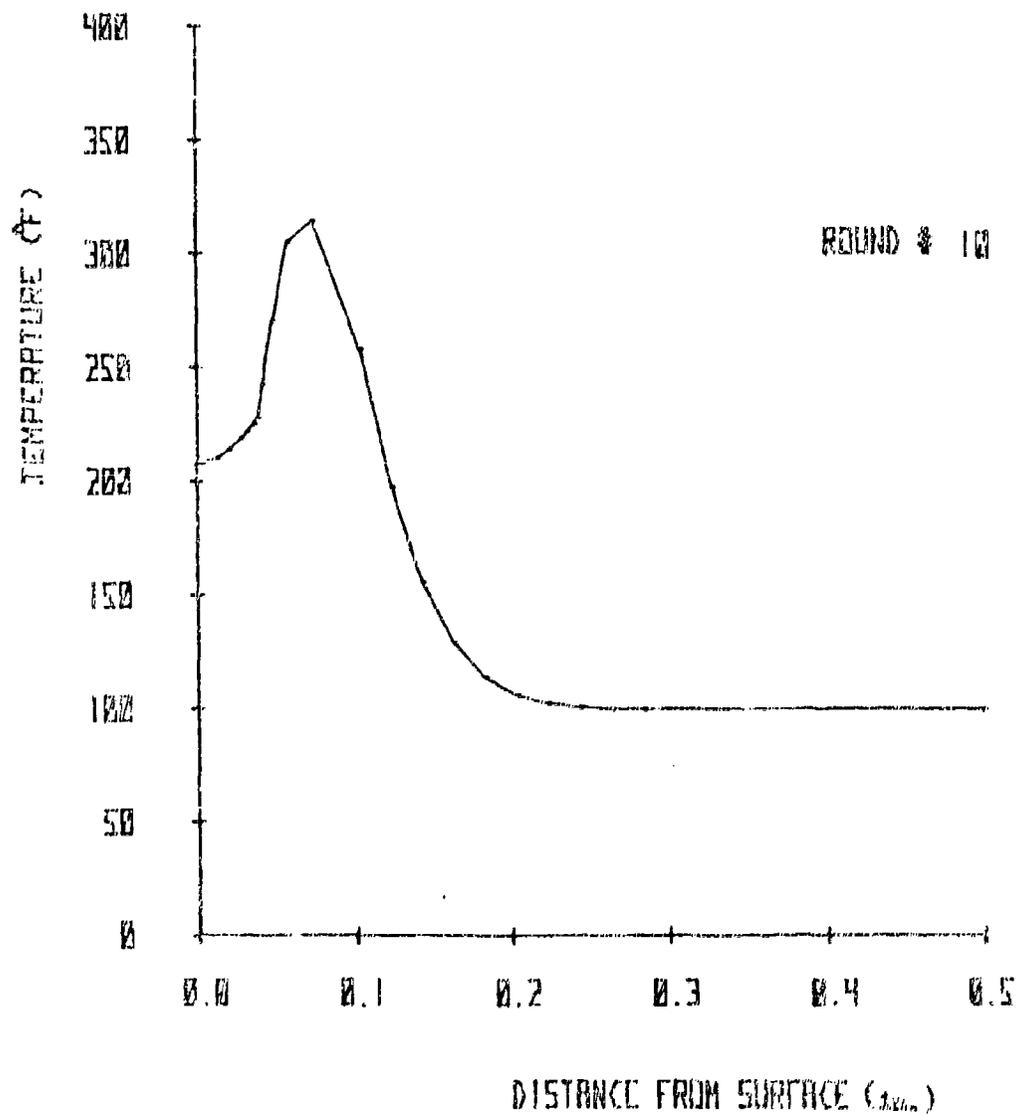


Figure 15. Temperature versus distance at 2000 (msec).

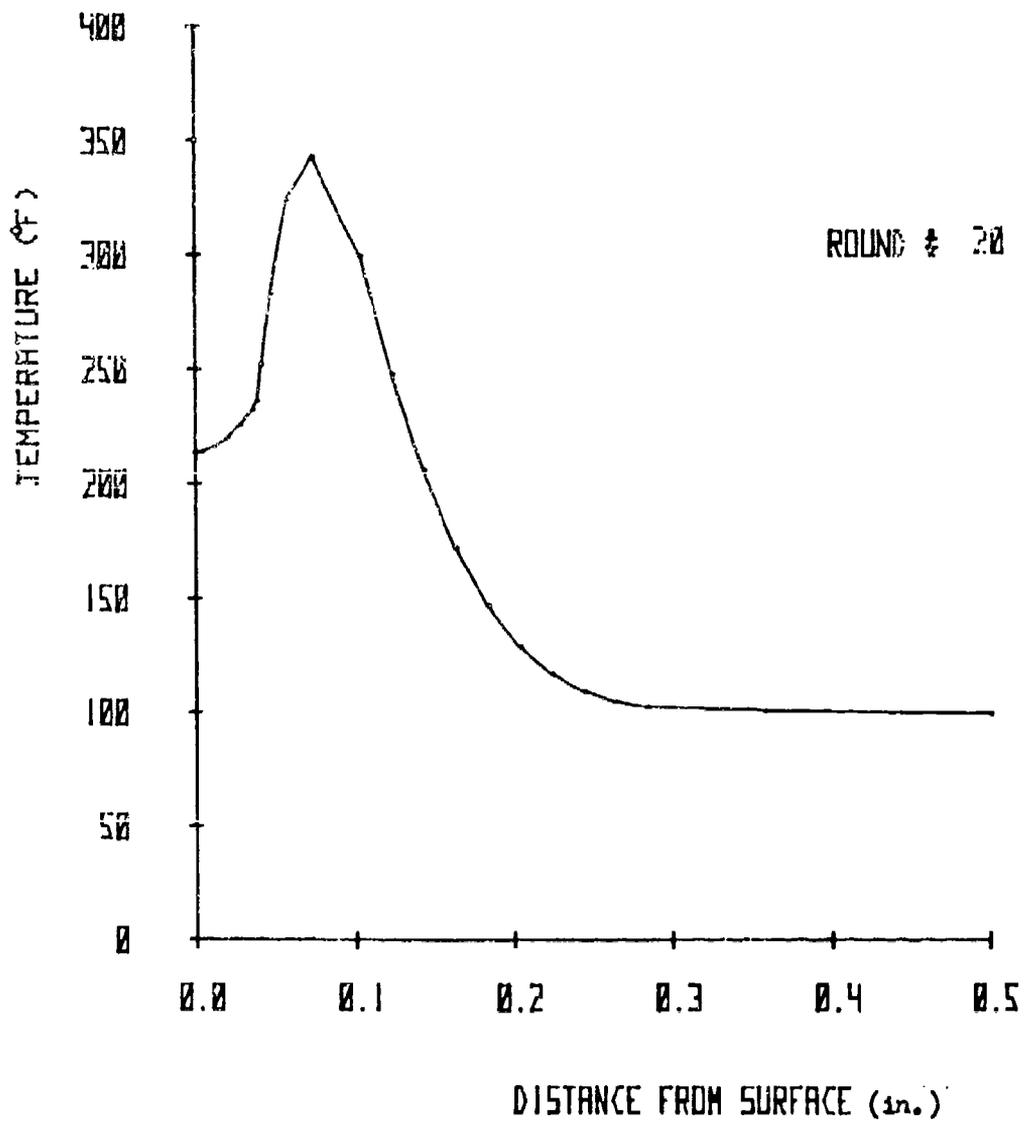


Figure 16. Temperature versus distance at 2000 (msec).

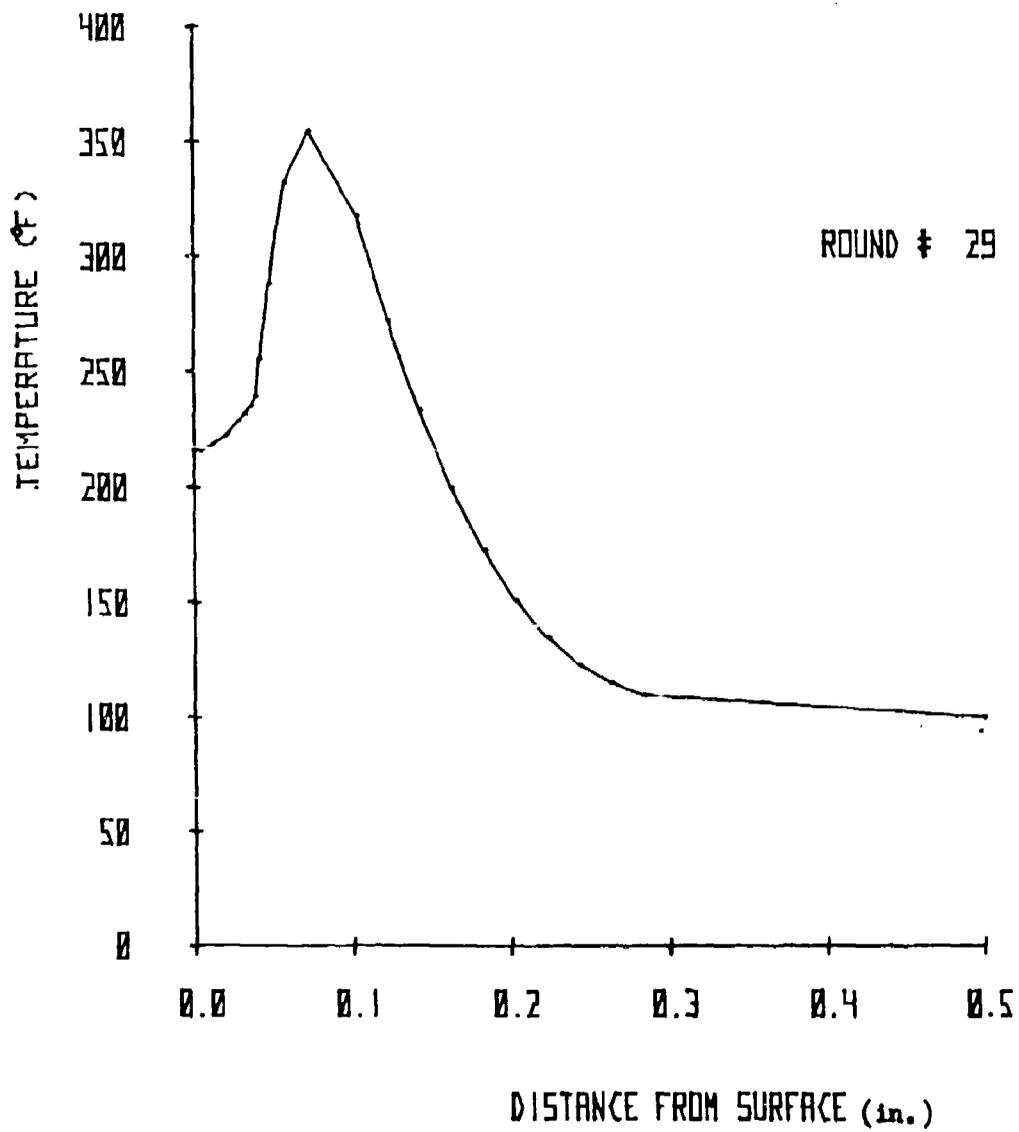


Figure 17. Temperature versus distance at 2900 (msec).

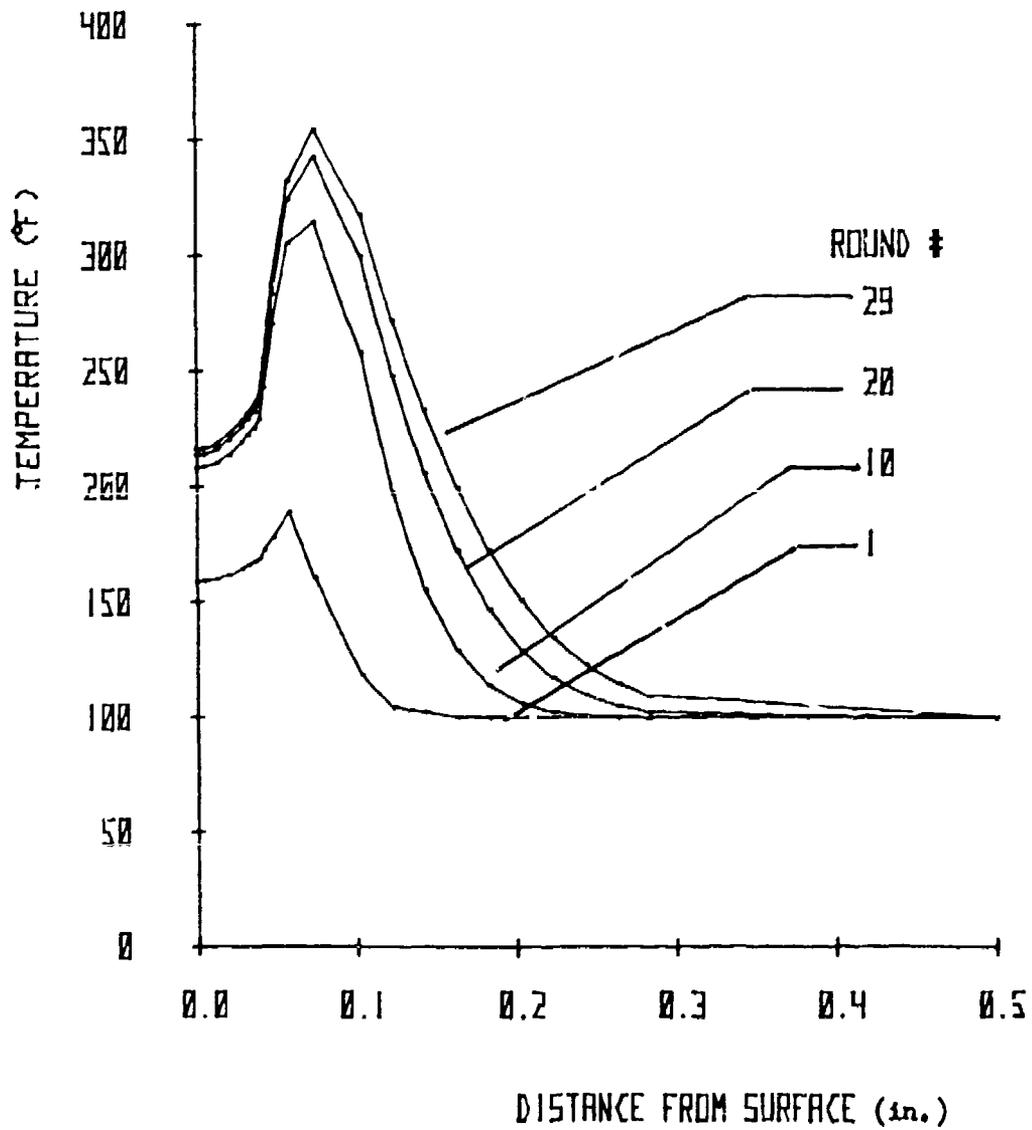


Figure 18. Temperature versus distance.

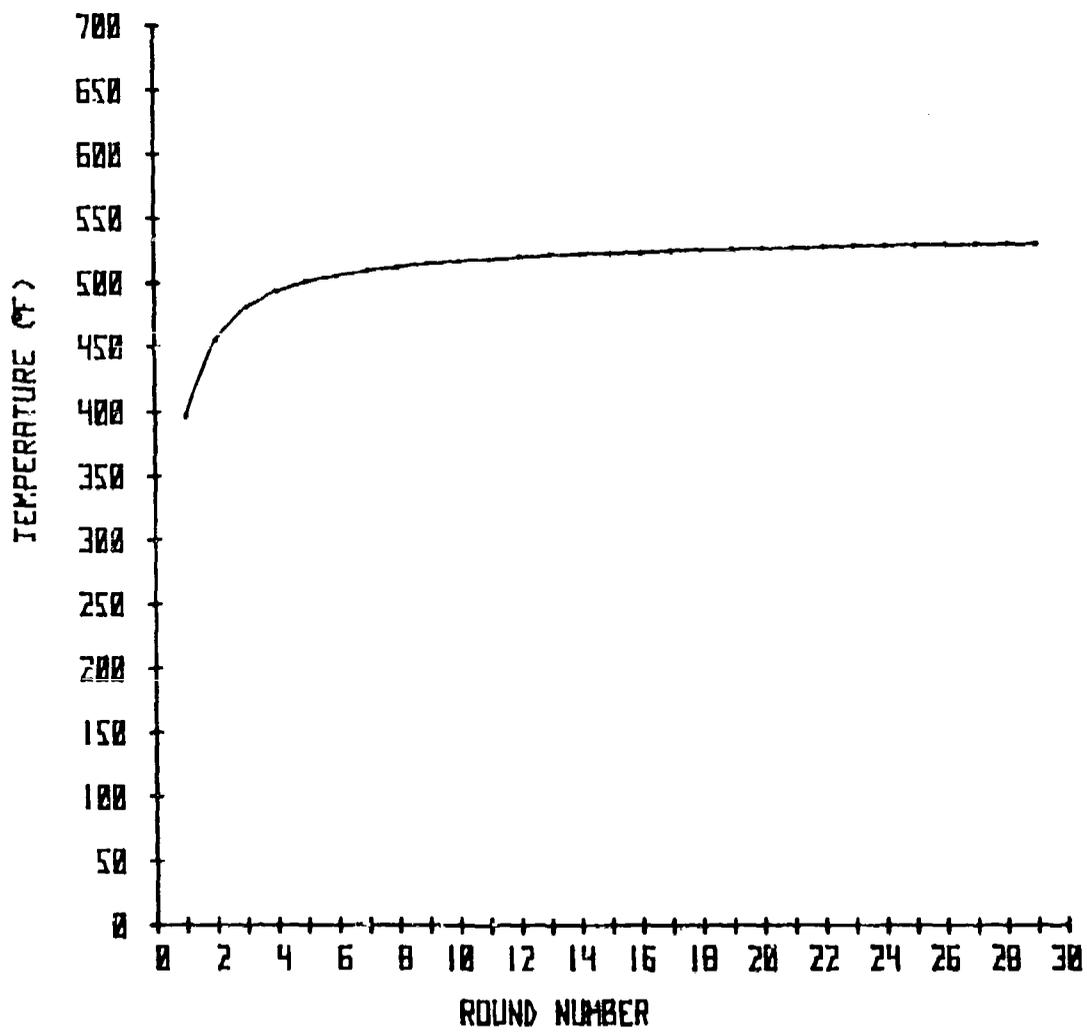


Figure 19. Peak interface temperature versus round number.

temperature. The peak temperature is still increasing at the rate of 0.3° F per round at round 29. Extrapolation of this rate out to 600 rounds is very dangerous. It would result in a peak steel temperature of approximately 700° F. This is an unrealistically high value. In addition, the question of depth of penetration must be considered. Examination of figure 18 shows that a 100° F temperature rise over the initial temperature has penetrated only 25% of the steel thickness. During future studies, additional cycles would be run.

The results of the one-dimensional analysis show that the penetration of the thermal wave is very slow, and therefore in those regions where this analysis is valid, conditions for the folded geometry are no worse than for conventional ammunition. In addition, since the severe heating during the ballistic phase is very short, this environment will heat the weapon uniformly and result in one-dimensional heat flow. A major factor affecting the structural temperature is the timing of the firing cycle. A shorter dwell time would greatly reduce the peak temperatures.

The major effect of the folded geometry is evidenced at the small interpose region near the base of the projectile. In this region the two-dimensional effect may be significant. The following section presents the two-dimensional results obtained for this region.

Two-Dimensional

Temperature distributions have been obtained for the model, and are described in the Two-Dimensional Model Section. The region examined is the base of the interpose region near the projectile base. Since the heating is uniform, only half of the web has to be considered. Figures 20 through 23 present the peak temperature occurring in this region for 4 complete cycles. The peak surface temperature in the corner is slightly higher than the peak temperature for the one-dimensional analysis. However, the temperature difference is approximately 1%, well within the accuracy of the analytical model. In figure 23, to the right of the model, the maximum temperatures corresponding to the same approximate location in the steel are given. For instance, the maximum interface temperature for round 4, node 24, in the one-dimensional analysis, is 493° F. The highest two-dimensional interface temperature at node 10 is 628° F. Node 10 is in the corner. However, node 24, located 0.41 in. from the corner, has a maximum of 521° F, as compared to 493° F in the one-dimensional interface temperature.

The results obtained from the two-dimensional calculations indicate that the temperature will be higher near the base of the interpose region. However, this level, at least after 4 rounds, is not prohibitive. In evaluating the temperature profiles, it will be necessary to assess the structural effect of small penetrations of high temperatures into the steel. This would be a major part of future efforts.

1	1298	13	1274	20	1274
2	980	14	936	25	936
3	816	15	754	26	754
4	695	16	621	28	621
5	630	18	546	29	545
6	618	19	508	30	507
7	574	21	485	32	484
8	554	22	466	33	466
9	539	23	452	34	451
10	523	17	438	24	437
11	560	51	319	55	238
12	570	52	198	56	139
13	579	53	161	59	103
14	592	54	140	60	101
15	608	57	160	62	100
16	633	31			
17	695	32			
18	816	33			
19	1298	34			
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68		83			
69		84			
70		85			
71		86			
72		87			
73		88			
74		89			
75		90			
76		91			
77		92			
78		93			
79		94			
80		95			
81		96			
82		97			
83		98			
84		99			
85		100			

Figure 20. Peak temperatures for round 1, (°F)

1	1328	13	1301	20	1302
2	1014	14	966	25	966
3	851	15	783	26	782
4	736	16	657	28	656
7	674	18	586	29	584
8	645	19	553	30	551
9	625	21	532	32	529
11	608	22	514	33	511
17	594	23	501	34	499
5	581	17	489	24	487
51	400	53	302	55	299
52	280	54	181	56	179
57	217	59	114	61	112
58	209	60	103	62	101
5	1319	5			
35	1004				
44	839				
45	723				
46	664				
47	641				
48	628				
49	617				
50	608				
51	601				
27	1290				
43	954				
44	770				
45	639				
46	577				
47	551				
48	538				
49	528				
50	519				
51	512				

Figure 21. Peak temperatures for round 2 (°F).

1	1338	13	1311	20	1311
2	1025	14	978	25	977
3	864	15	802	26	801
4	750	16	672	28	671
5	690	18	603	29	601
6	662	19	572	30	569
7	644	21	552	32	549
8	628	22	535	33	532
9	616	23	523	34	520
10	604	17	512	24	509
11	51	53	335	55	331
12	438	54	212	56	208
13	331	57	260	61	124
14	247	58	108	62	103
15	1296	27	1296	27	1296
16	962	43	962	43	962
17	849	35	849	35	849
18	733	37	733	37	733
19	676	38	676	38	676
20	654	39	654	39	654
21	642	40	642	40	642
22	632	41	632	41	632
23	625	42	625	42	625
24	618	43	618	43	618

Figure 22. Peak temperatures for round 3 (°F).

1	1343	13	1317	20	1316
2	1031	14	986	25	985
3	869	15	804	26	802
4	757	16	681	28	679
7	698	18	613	29	610
8	671	19	582	30	579
5	654	21	563	32	560
11	638	22	547	33	544
12	628	23	536	34	532
5	615	17	524	24	521 (493)
51	458	53	355	55	(445) (409) (347)
52	361	54	230	56	(270) (187) (134)
57	293	59	144	61	(113) (101)
58	277	60	114	62	(100) (100)
31	540				
50	546				
49	553				
48	562				
47	573				
46	596				
45	660				
44	784				
43	968				
5	1330				

Figure 23. Peak temperatures for round 4 (°F).

CONCLUSIONS

The thermal response of the folded geometry cartridge has been determined for one and two-dimensional models, and also for single and multiple firings. The results show that the surface of the cartridge case heats very rapidly due to the propellant combustion. The one-dimensional model has been run for 29 rounds, and the two-dimensional model for 4 rounds.

The peak cartridge case temperature occurs at the base of the interpose region, based upon two-dimensional calculations. At round 4, this peak is 1343°F . For the one-dimensional model, the corresponding peak is 1331°F . The two-dimensional effect on peak cartridge temperature is therefore very small. Extrapolating the one-dimensional results to 600 rounds would yield a peak of 1370°F . By applying the same difference of approximately 1.0%, for the worst two-dimensional results, the peak cartridge temperature would be less than 1400°F after 600 rounds. This result would indicate that with respect to peak cartridge temperature, the folded geometry cartridge is no worse than conventional rounds.

With respect to the weapon chamber, the interface temperature at the base of the interpose region will be hotter, due to the folded geometry. After 4 rounds, the peak interface temperature is 628°F for the two-dimensional analysis, compared to 493°F for the one-dimensional analysis. Extrapolation of these results, to conditions after 600 rounds, is not possible at this time. However, it seems reasonable to conclude that the base of the interpose region will experience more severe temperatures than would exist with conventional ammunition.

Future effort should concentrate on additional firings with the two-dimensional model, an investigation of the heating cycle during the period after the ballistic portion of the cycle, and thermal stress calculations of the effect of thermal distortion and elevated temperatures upon material strength.

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APPENDIX A

ONE-DIMENSIONAL PROGRAM LISTING

// JOB MCA>SEY,3610****10/*****CLASS=LIFZ/
// EXEC MATFIV,SIZE=0256K

/PROGRAM MCASSEY,PAGES=500,TIME=15

***** ONE DIMENSIONAL THERMAL ANALYSIS PROGRAM *****

```

1  REAL M
2  DIMENSION TFREE(50),FILM(20),TIME(50),M(100),C(100),AREA(100),TI(1
3  LOG),COY(100,100),TEMP(200),TNER(100)
4  READ(5,1000) DELTA,NBODY,NCOUP,N
5  WRITE(6,2002) DELTA,NBODY,NCOUP,N
6  READ(5,1001) ICT,INC,INTER,COEL
7  CDEL=COEL/3600.
8  WRITE(5,2008) ICT,INC,INTER,COEL
9  READ(5,1001) TFREE(I),FILM(I),TIME(I),I=1,N
10 PRINT=1.35E-07
11 PRO=0.0
12 DO 4 I=1,N
13   YFREE(I)=450.0
14   TIME(I)=TIME(I)/3600.0
15   WRITE(6,2004) (TFREE(I),FILM(I),TIME(I),I=1,N)
16   DO 5 I=1,NBODY
17     DREAL(5,1005) M(I),C(I),AREA(I),TI(I)
18     CO=0. I=1,NBODY
19     COY(I)=I,NBODY
20     COY(I)=I,NBODY
21     COY(I)=I,NBODY
22     COY(I)=I,NBODY
23     READ(5,1007) I,J,COND(I,J)
24     COND(I,J)=COND(I,J)
25     CONTINUE=1,NBODY
26     WRITE(6,2001) M(I),C(I),AREA(I),TI(I)
27     WRITE(6,2003) J,I,NBODY
28     DO 10 J=1,NBODY
29       CONTINUE=1
30       CONTINUE=1
31       CONTINUE=1
32       CONTINUE=1
33       CONTINUE=1
34       CONTINUE=1
35       CONTINUE=1
36       CONTINUE=1
37       CONTINUE=1
38       CONTINUE=1
39       CONTINUE=1
40       CONTINUE=1
41       CONTINUE=1
42       CONTINUE=1
43       CONTINUE=1
44       CONTINUE=1
45       CONTINUE=1
46       CONTINUE=1
47       CONTINUE=1
48       CONTINUE=1
49       CONTINUE=1
50       CONTINUE=1
51       CONTINUE=1
52       CONTINUE=1

```



```
111 $ / 1A, 'INTERFACE AT BODY NUMBER', I2, / 1A, 'NEW TIME STEP=', E20.5)
112 3600 FORMAT (1X, 'BODY NUMBER=', I5)
      END
      / GO
```

APPENDIX B

TWO-DIMENSIONAL PROGRAM LISTING

// JOB MCASSEY,3610****10/****,CLASS=L(FZ)
// EXEC MATFIV,SIZE=0256K

/PROGRAM MCASSEY,PAGES=500,TIME=b

C
C
C
C

***** TWC DIMENSIONAL TEMPERATURE PROGRAM *****

```

1 REAL M,ICN,TFREE(20),FILM(KU),TIME(20),M(100),C(100),AREA(100),TI(1
2 100),N(100),DELTA,INEM(100),INEM(100)
3 DIMEN,ICN,ITYPE(100)
4 READ(5,1000) DELTA,NBODY,NCUP,N
5 WRITE(6,2002) DELTA,NBODY,NCUP,N
6 READ(5,1001) TFREE(I),FILM(I),TIME(I),I=1,N)
7 PRINT=675E-07
8 C=0
9 I=1,N
10 TFREE(I)=TFREE(I)-460.0
11 TIME(I)=TIME(I)/3600.0
12 WRITE(6,2003)
13 WRITE(6,2004) (TFREE(I),FILM(I),TIME(I),I=1,N)
14 DO 5 I=1,NBODY
15 READ(5,1005) M(I),C(I),AREA(I),TI(I),ITYPE(I)
16 DO 6 J=1,NBODY
17 COND(I,J)=0.0
18 DO 7 K=1,NCUP
19 READ(5,1007) I,J,COND(I,J)
20 COND(I,J)=COND(I,J)
21 DO 10 I=1,NBODY
22 WRITE(6,3001) M(I),C(I),AREA(I),TI(I),ITYPE(I)
23 WRITE(6,2006)
24 DO 15 J=1,NBODY
25 IF COND(I,J).GT.0.0 WRITE(6,2007) I,J,COND(I,J)
26 CONTINUE
27 M(I)=M(I)*3.6E6
28 TT=TIME(I)
29 CC=16
30 I=1,NBODY
31 TIME(I)=I(I)
32 WRITE(6,3000) TT
33 WRITE(6,3001) J,TI(J),J=1,NBODY)
34 LCPK=TIME(I)+1-TIME(I)
35 SLCPE=(TFREE(I)+1)-TFREE(I)/WORK
36 HSLCP=(FILM(I)+1)-FILM(I)/WORK
37 TSTCPT=TIME(I)
38 TT=I*START
39 I=I+DELTA
40 IF(TT.GT.TSTOP) TT=TSTOP
41 DO 52 I=1,NBODY
42 TEMP(I)=INEM(I)
43 CUR=DELTA/M(I)/C(I)
44 IF(ITYPE(I).EQ.1) GC=TJ*OU
45 SUM=WORK*AREA(I)*(FILM(I))+HSLCP*(TT-TSTART))*TFREE(I)+SLOPE*(TT
46 1-TSTART)
47 GC=TC*61
48 SUM=0.0
49
50
51
52
60

```


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