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TRANSONIC STATIC STABILITY CHARACTERISTICS OF BOMBLET MUNITION MODELS USED IN THE EVALUATION OF THE ZERO-CONING AERODYNAMIC DISPERSAL TECHNIQUE

T. O. Shadow
ARO, Inc.

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TRANSONIC STATIC STABILITY CHARACTERISTICS OF BOMBLET MUNITION MODELS USED IN THE EVALUATION OF THE ZERO-CONING AERODYNAMIC DISPERSAL TECHNIQUE

T. O. Shadow
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FOREWORD

The work reported herein was sponsored by the Air Force Armament Laboratory (DLRA/M. J. Bouffard), Armament Development and Test Center, Air Force Systems Command (AFSC), under Program Element 62602F, Project 2547.

The test results presented were obtained by ARO, Inc. (a subsidiary of Sverdrup & Parcel and Associates, Inc.), contract operator of the Arnold Engineering Development Center (AEDC), AFSC, Arnold Air Force Station, Tennessee, under Contract F40600-72-C-0003. The test was conducted September 16 and 17, 1971, under ARO Project No. PC0205. The manuscript was submitted for publication on October 12, 1971.

This technical report has been reviewed and is approved.

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ABSTRACT

A wind-tunnel investigation was conducted in the Aerodynamic Wind Tunnel (4T) to determine the static stability characteristics of bomblet munition models designed for the evaluation of the Zero-Coning Aerodynamic Dispersal Technique. Force and moment data were recorded at Mach numbers from 0.3 to 1.2 at a constant Reynolds number of $2.2 \times 10^6$ per foot. Angle of attack was varied from -8 to 27 deg. Roll angle was varied from 0 to 30 deg on one configuration. The test results indicate that the configurations tested are marginally acceptable to achieve Zero-Coning dispersion.
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### NOMENCLATURE

- **\( C_A \)**: Axial-force coefficient, measured axial force/\( q_m S \)
- **\( C_m \)**: Pitching-moment coefficient (see Fig. 3 for moment reference location), measured pitching moment/\( q_m S_d \)
- **\( C_N \)**: Normal-force coefficient, measured normal force/\( q_m S \)
- **\( d \)**: Model body diameter (reference diameter), 0.2500 ft
- **\( M_m \)**: Free-stream Mach number
- **\( q_m \)**: Free-stream dynamic pressure, psf
- **\( Re/ft \)**: Reynolds number per foot
- **\( S \)**: Model cross-sectional area (reference area), 0.0491 sq ft
- **\( \alpha \)**: Model angle of attack with respect to tunnel centerline, deg
- **\( \phi \)**: Model roll angle, deg

### MODEL NOMENCLATURE

- **\( B_S \)**: Cylindrical centerbody, \( d = 0.2500 \) ft, \( S = 0.0491 \) sq ft
- **\( N_{S1} \)**: Blunted nose, 0.750-in. shoulder radius
- **\( N_{S2} \)**: Spherical nose, 1.500-in. spherical radius
- **\( N_{S3} \)**: Blunted nose, 0.375-in. shoulder radius
- **\( N_{S4} \)**: Spherical-segment nose, 2.250-in. spherical radius
- **\( A_{S1} \)**: Cylindrical afterbody
- **\( A_{S2} \)**: Boattail afterbody
$F_{S2,5}$  Rectangular fins for cylindrical afterbody; 2 and 5 for 3.466- and 3.840-in. spans, respectively

$F_{S3,6}$  Rectangular fins for boattail afterbody; 3 and 6 for 3.466- and 3.840-in. spans, respectively
SECTION I
INTRODUCTION

A wind-tunnel investigation of a group of munition models was conducted in the Aerodynamic Wind Tunnel (4T), Propulsion Wind Tunnel Facility (PWT), to determine the static stability characteristics for use in the evaluation of the Zero-Coning aerodynamic dispersal technique. The tests were conducted at Mach numbers from 0.3 to 1.2 at a constant Reynolds number of $2.2 \times 10^6$ per foot.

With the Zero-Coning concept, effective bomblet dispersion attributable to lift is achieved by restriction of the rotation of the angle-of-attack plane in space. A bomblet which is unstable at low angles of attack, stable at high angles of attack, and has lift which is insensitive to roll angle will disperse along a curved flight path.

Previous investigations of the Zero-Coning concept reported in Refs. 1 through 4 have provided static and dynamic stability data on several configurations. The purpose of this investigation was to expand the knowledge of nose shape and fin span effects on static stability characteristics and to investigate configuration modifications to improve stability and reduce the bomblet susceptibility to coning under off-design conditions.

SECTION II
APPARATUS

2.1 TEST FACILITY

The Aerodynamic Wind Tunnel (4T) is a closed-loop, continuous flow, variable density tunnel in which the Mach number can be varied from 0.1 to 1.3. At all Mach numbers, the stagnation pressure can be varied from 300 to 3700 psf. The test section is 4 ft square and 12.5 ft long with perforated, variable porosity (0.5- to 10-percent open) walls. It is completely enclosed in a plenum chamber from which the air can be evacuated, allowing part of the tunnel airflow to be removed through the perforated walls of the test section. A more thorough description of the tunnel is given in Ref. 5.

The model support system consists of a pitch sector, boom, and sting which provide a pitch capability from -12 to 28 deg with respect to the tunnel centerline. The center of rotation is at tunnel station 108. In addition, a remote controlled roll mechanism allows roll angle variations of $\pm 180$ deg. A schematic of the test section, showing the location of the test model, is shown in Fig. 1 (Appendix I). A photograph of the test installation is presented in Fig. 2.

2.2 TEST ARTICLES

Details of the bomblet models are shown in Fig. 3. The models consisted of a 3-in.-diam cylindrical centerbody, four noses with varying bluntness, straight or boattail afterbody, and rectangular fins with two spans for each afterbody.
2.3 INSTRUMENTATION

An internal six-component strain-gage balance was used to measure forces and moments on the models. Differential pressure transducers, referenced to the tunnel plenum pressure, were used to measure the pressure at the base of the models. Electrical signals from the balance, pressure transducers, and standard tunnel instrumentation were processed by the PWT data acquisition system and digital computer.

SECTION III
TEST DESCRIPTION

3.1 TEST CONDITIONS

Data were obtained at nominal Mach numbers from 0.3 to 1.2 at a constant Reynolds number of $2.2 \times 10^6$ per foot. Stagnation temperature was maintained at approximately $110^\circ F$ throughout the test. Angle of attack was varied from -8 to 27 deg. One configuration was selected for roll angle variations of 0, 15, and 30 deg. Free transition was used throughout the test.

3.2 PRECISION OF MEASUREMENTS

The estimated precision of the data presented in this report, based on a 95-percent confidence level, is given below and in Fig. 4. The error sources considered for the coefficients were balance uncertainties, Mach number nonuniformities, instrument errors, and Mach number calibration accuracies.

$$\Delta a \pm 0.1$$
$$\Delta \phi \pm 0.1$$

The uncertainties in Mach number include variation of Mach number along the tunnel centerline, instrument errors. and errors in data acquisition techniques.

SECTION IV
RESULTS AND DISCUSSION

4.1 GENERAL

The recorded balance data were reduced to aerodynamic coefficients in the nonrolling body-axis coordinate system with the moment reference point 1.5 cal from the nose of the model.

Base pressure measurements were made on all models; however, no corrections have been made to the data presented. Since the models represent bomblets which are free-dropped with no rocket assist, it is assumed that the pressures acting on the base of the models in the wind tunnel are representative of those in free flight.
4.2 **ROLL ANGLE EFFECTS**

One essential characteristic that a bomblet must possess to disperse with zero coning is that lift be insensitive to model roll angle. The effects of roll angle on the longitudinal characteristics of a typical bomblet model are presented in Fig. 5. There were essentially no roll angle effects.

4.3 **NOSE SHAPE EFFECTS**

The effects of nose shape on \( C_N \), \( C_m \), and \( C_A \), for a cylindrical and a boattail afterbody, are shown in Figs. 6 and 7, respectively. The small-radius blunt nose produced the highest normal-force coefficients and the spherical nose the lowest at the high angles of attack throughout the Mach number range with the cylindrical afterbody (Fig. 6a). The criteria for zero-coning dispersal (unstable at low angles of attack, stable at high angles of attack) is obviously violated by the spherical-segment nose configuration for \( M_e < 1.1 \) (Fig. 6b). In addition, the spherical-nose configuration, which comes nearest to meeting the requirements, is at best marginal at the subsonic Mach numbers. Axial-force coefficients were highest with the spherical-segment nose configuration and lowest with the spherical-nose configuration (Fig. 6c).

Normal-force coefficients were essentially insensitive to nose shape with the boattail afterbody except at supersonic Mach numbers and high angles of attack where the coefficients were highest with the large-radius blunt nose and lowest with the spherical nose (Fig. 7a). Again, the low angle-of-attack stability characteristics are less desirable with the spherical-segment nose (Fig. 7b). Other nose shapes are marginally unstable near \( \alpha = 0 \) deg with no highly stable trim angles. Axial-force coefficients show the same trends as with the cylindrical afterbody (Fig. 7c).

4.4 **AFTERBODY EFFECTS**

Comparisons are made of the effects of cylindrical and boattail afterbodies on the longitudinal characteristics of models with spherical and spherical-segment noses in Figs. 8 and 9, respectively. Normal-force coefficients for the configuration with spherical nose were highest with the boattail afterbody at all Mach numbers except 0.8 and 0.9 where they were equal or greater with the cylindrical afterbody (Fig. 8a). The comparisons in pitching-moment coefficients (Fig. 8b) show that both afterbodies had approximately the same stability characteristics near \( \alpha = 0 \) for \( M_e < 0.8 \). At \( M_e = 0.9 \), the boattail base was slightly more unstable near \( \alpha = 0 \) deg; however, there were no strongly stable trim angles. In Fig. 8c, it can be seen that the cylindrical afterbody produced the higher axial force but to a lesser extent as Mach number was increased.

For the configuration with the spherical-segment nose, normal-force coefficients were highest with the boattail afterbody except in the region \( 0.7 \leq M_e \leq 0.9 \) where the coefficients were highest with the cylindrical afterbody (Fig. 9a). Both afterbody configurations showed stable characteristics at angles of attack near zero and subsonic Mach numbers with the spherical-segment nose. At supersonic Mach numbers, both configurations
became slightly unstable near $\alpha = 0$ deg (Fig. 9b). The same trends in axial-force coefficients, as seen with the spherical nose, were obtained with the spherical-segment nose (Fig. 9c).

4.5 FIN SPAN EFFECTS

The effects of the variation of fin span on model longitudinal characteristics are presented in Figs. 10 through 13 for various nose and afterbody configurations. In general, the larger fin span produced the largest normal-force coefficient for all configurations (Figs. 10a, 11a, 12a, and 13a). The fins with the shortest span caused the models to be more unstable near $\alpha = 0$ deg (Figs. 10b, 11b, 12b, and 13b); however, only two of the configurations showed unstable characteristics near $\alpha = 0$ deg and a high angle-of-attack trim point. At supersonic Mach numbers with the cylindrical afterbody, short-span fins, blunt- and spherical-segment noses, trim angles of approximately 12 deg were noted (Figs. 10b and 11b). Fin span had negligible effect on axial-force coefficients (Figs. 10c, 11c, 12c, and 13c).

SECTION V
CONCLUSIONS

The following conclusions were made from the results of the investigation:

1. Increasing fin span makes the bomblet models more stable near zero angle of attack, at which there is no lift for dispersion.

2. Making the model nose more blunt causes the models to be more stable near zero angle of attack.

3. Neither the cylindrical nor the boattail afterbody shape produces the desired instability near zero angle of attack.

4. The normal-force coefficients for the bomblet models are insensitive to roll angle.

REFERENCES


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ILLUSTRATIONS
Fig. 1 Schematic of Tunnel Installation
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ALL DIMENSIONS IN INCHES

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Fig. 5 Continued

\( a. \) Concluded
b. $C_m$ versus $\alpha$

Fig. 5 Continued
b. Concluded
Fig. 5 Continued
\[ C_A \text{ versus } \alpha \]

Fig. 5 Continued
Fig. 6 Effects of Nose Shape on Longitudinal Characteristics of Model with Cylindrical Afterbody, $B_{S}N_{S}A_{S_{1}}F_{S_{5}}, \phi = 0$
Fig. 6 Continued

- a. Concluded

- Fig. 6 Continued
b. $C_m$ versus $\alpha$

Fig. 6 Continued
b. Concluded

Fig. 6 Continued
c. $C_A$ versus $\alpha$

Fig. 6 Continued
Fig. 6  Concluded

- $M_∞ = 0.90$
- $M_∞ = 1.10$
- $M_∞ = 1.20$

Concluded
Fig. 7 Effects of Nose Shape on Longitudinal Characteristics of Model with Boattail
Afterbody, $B_s N_s x A_{s2} F_{ss}, \phi = 0$
Fig. 7 Continued
b. $C_m$ versus $\alpha$

Fig. 7 Continued
b. Concluded

Fig. 7 Continued
c. $C_A$ versus $\alpha$

Fig. 7 Continued
c. Concluded
Fig. 7 Concluded
Fig. 8 Effects of Afterbody Shape on Longitudinal Characteristics of Model with Spherical Nose, $B_S N_{S2} A_{Sx} F_{Sx}$, $\phi = 0$
Fig. 8 Continued

b. $C_m$ versus $\alpha$

Fig. 8 Continued
Fig. 8 Continued

b. Concluded
Fig. 8 Continued

c. $C_A$ versus $\alpha$

$M_\infty = 0.3$

$M_\infty = 0.5$

$M_\infty = 0.7$

$M_\infty = 0.8$
c. Concluded
Fig. 8 Concluded
Fig. 9 Effects of Afterbody Shape on Longitudinal Characteristics of Model with Spherical-Segment Nose, $B_s N_{94} A_{5x} \phi_0 = 0$
a. Concluded

Fig. 9 Continued
b. $C_m$ versus $\alpha$

Fig. 9 Continued
b. Concluded
Fig. 9 Continued
c. $C_A$ versus $\alpha$

Fig. 9 Continued
c. Concluded
Fig. 9 Concluded
Fig. 10 Effects of Fin Span on Longitudinal Characteristics of Model with Cylindrical Afterbody and Blunted Nose, $B_sN_{s3}A_{s1}F_{sX}$, $\phi = 0$
Fig. 10 Continued

\( C_N \) vs. \( \alpha \) for different Mach numbers:
- \( M_{\infty} = 0.9 \)
- \( M_{\infty} = 1.1 \)
- \( M_{\infty} = 1.2 \)

\( \square \) \( B_5N_3P_5F_{52} \)
\( \circ \) \( B_3N_3P_3F_{55} \)

a. Concluded
b. $C_m$ versus $\alpha$

Fig. 10 Continued
b. Concluded
Fig. 10 Continued
c. $C_A$ versus $\alpha$

Fig. 10 Continued
\[ M = 0.9 \]
\[ M = 1.1 \]
\[ M = 1.2 \]

\[ C_R \]

\[ M_\infty = 0.9 \]
\[ M_\infty = 1.1 \]
\[ M_\infty = 1.2 \]

\[ \alpha \]

**Fig. 10 Concluded**

c. Concluded
Fig. 11 Effects of Fin Span on Longitudinal Characteristics of Model with Cylindrical Afterbody and Spherical-Segment Nose, $B_5 N_s A_s F_X$, $\phi = 0$

\[ C_N \text{ versus } \alpha \]
Fig. 11 Continued
b. $C_m$ versus $\alpha$

Fig. 11 Continued
b. Concluded

Fig. 11 Continued
Fig. 11 Continued

c. $C_A$ versus $\alpha$
B_3 N_{54} A_{31} F_{52}
B_3 N_{54} A_{31} F_{55}

Fig. 11 Concluded
Fig. 12 Effects of Fin Span on Longitudinal Characteristics of Model with Boattail Afterbody and Blunted Nose, $B_sN_sA_sF_{sx}$, $\phi = 0$

a. $C_N$ versus $\alpha$
Fig. 12 Continued

- Concluded

\[ C_n \text{ vs. } \alpha \]

- Points represent data for different conditions:
  - \( B_3N_3\alpha_{90}F_{53} \)
  - \( B_3N_3\alpha_{90}F_{56} \)

Grid for plotting data points with axes labeled:
- \( M_\infty = 0.9 \)
- \( \alpha \)

Graph shows trends in aerodynamic forces with varying angles of attack.
b. $C_m$ versus $\alpha$

Fig. 12 Continued
b. Concluded
Fig. 12 Continued
Fig. 12  Continued

Cₐ versus α

Fig. 12  Continued
\[ C_A \]

- \( M_\infty = 0.9 \)
- \( M_\infty = 1.1 \)
- \( M_\infty = 1.2 \)

\[ \text{Fig. 12 Concluded} \]

c. Concluded
Fig. 13  Effects of Fin Span on Longitudinal Characteristics of Model with Boattail Afterbody and Spherical-Segment Nose, $B_SN_sA_sF_{sx}$, $\phi = 0$
a. Concluded
Fig. 13 Continued
b. $C_M$ versus $\alpha$

Fig. 13 Continued
b. Concluded

Fig. 13 Continued
Fig. 13 Continued

C_A versus \( \alpha \)

- \( M_\infty = 0.3 \)
- \( M_\infty = 0.5 \)
- \( M_\infty = 0.7 \)
- \( M_\infty = 0.8 \)

Symbols:
- \( \bigcirc \): \( B_5 N_{54} A_{52} F_{53} \)
- \( \square \): \( B_5 N_{54} A_{52} F_{56} \)
Fig. 13 Concluded

c. Concluded
A wind-tunnel investigation was conducted in the Aerodynamic Wind Tunnel (4T) to determine the static stability characteristics of bomblet munition models designed for the evaluation of the Zero-Coning Aerodynamic Dispersal Technique. Force and moment data were recorded at Mach numbers from 0.3 to 1.2 at a constant Reynolds number of $2.2 \times 10^6$ per foot. Angle of attack was varied from -8 to 27 deg. Roll angle was varied from 0 to 30 deg on one configuration. The test results indicate that the configurations tested are marginally acceptable to achieve Zero-Coning dispersion.
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