

THE OHIO STATE UNIVERSITY



RESEARCH FOUNDATION

1314 KINNEAR ROAD COLUMBUS, OHIO 43212

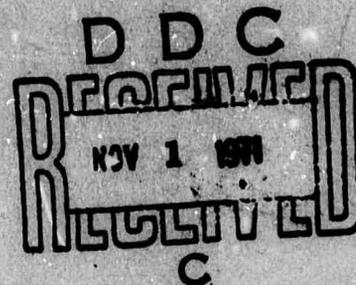
AD 731688

**STRESSES, DEFORMATIONS AND PROGRESSIVE
FAILURE OF NON-HOMOGENEOUS FISSURED ROCK**

Semi-annual Technical Report; August 30, 1971

U.S. Bureau of Mines, Contract No. HO210017

Sponsored by
Advanced Research Projects Agency
ARPA Order No. 1579, Amend. 2
Program Code 1F10



The views and conclusions contained in this document are those of the authors and should not be interpreted as necessarily representing the official policies, either expressed or implied, of the Advanced Research Projects Agency or the U.S. Government.

Reproduced by
**NATIONAL TECHNICAL
INFORMATION SERVICE**
Springfield, Va. 22151

**BEST
AVAILABLE COPY**

Mar 7, 66

Security Classification

DOCUMENT CONTROL DATA - R & D

(Security classification of title, body of abstract and indexing symbols must be entered when the overall report is classified)

1. ORIGINATING ACTIVITY (Corporate author)		2a. REPORT SECURITY CLASSIFICATION	
The Ohio State University Research Foundation		Unclassified	
2b. GROUP			
3. REPORT TITLE			
STRESSES, DEFORMATIONS AND PROGRESSIVE FAILURE OF NON-HOMOGENEOUS FISSURED ROCK			
4. DESCRIPTIVE NOTES (Type of report and inclusive dates)			
Semiannual Technical Report - February 1, 1971 - July 31, 1971			
5. AUTHOR(S) (First name, and the initial, last name)			
R. S. Sandhu, T. H. Wu and J. R. Hooper			
6. REPORT DATE		7a. TOTAL NO. OF PAGES	7b. NO. OF REFS
August 30, 1971		100	44 items
8a. CONTRACT OR GRANT NO.		9a. ORIGINATOR'S REPORT NUMBER(S)	
HO210017		None	
b. PROJECT NO.		9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)	
RF 3177 A1		None	
c.			
d.			
10. DISTRIBUTION STATEMENT			
Distribution of this document is unlimited.			
11. SUPPLEMENTARY NOTES		12. SPONSORING MILITARY ACTIVITY	
		Advanced Research Projects Agency Washington, D.C. 20301	
13. ABSTRACT			
<p>The objective of this research program is to develop computer programs, using the finite element method, to predict stresses and deformations in the vicinity of underground excavations. The computer programs should allow for arbitrary initial stresses in rock, arbitrary shape and size of the opening, any given sequence of construction, non-homogeneous material properties, interaction of rock with supporting structures, progressive damage and time dependent deformation and load development on supporting structure. Limited experimental work to verify key points in the theory is planned.</p> <p>Research during the first year is directed towards literature survey, selection of mathematical models for behavior of rock and development of computer programs for elastic-plastic, elastic-brittle rock and for progressive failure of rock around underground openings.</p> <p>At this reporting, selection of mathematical models has been completed. The report presents finite element computer programs for i) Plane Strain Analysis of Elastic-Plastic Mohr-Coulomb Materials and ii) Two-Dimensional Analysis of a No Tension System. Plans for further requirements of these programs in line with research objectives are discussed.</p>			

KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
computation						
deformation						
elasticity						
excavation						
failure						
finite element method						
foundations						
mining						
plasticity						
progressive failure						
research						
rock mechanics						
stresses						
tunnels						
underground excavation						

SEMIANNUAL TECHNICAL REPORT

FEBRUARY 1, 1971 - JULY 31, 1971

ARPA Order Number: 1579, Amend 2

Contract Number: HO210017

Program Code Number: 1F10

Principal Investigators:

R. S. Sandhu

T. H. Wu

J. R. Hooper

Telephone Number: (614) 422-7531

Name of Contractor:

**The Ohio State University
Research Foundation**

Project Scientist or Engineer:

R. S. Sandhu

Telephone Number: (614) 422-7531

Effective Date of Contract:

February 1, 1971

Short Title of Work:

**Stresses, Deformations and
Progressive Failure of Non-
Homogeneous Fissured Rock**

Contract Expiration Date:

January 31, 1972

Amount of Contract:

\$46,113

This research was supported by the Advanced Research Projects Agency of the Department of Defense and was monitored by Bureau of Mines under Contract Number HO210017.

Distribution of this document is unlimited.

Technical Report Summary

The objective of this research program is to develop computer programs, using the finite element method, to predict stresses and deformations in the vicinity of underground excavations. The computer programs will have the capability to allow for arbitrary initial stresses in rock, arbitrary shape and size of the opening, any given sequence of construction, nonhomogeneous material properties, interaction of rock with supporting structures, progressive damage, and time dependent deformation and load development on supporting structure. Limited experimental work to verify key points in the theory is planned.

Research during the first year is directed towards survey of literature on the subject, selection of mathematical models for mechanical behavior of rock, and development of computer programs for elastic-plastic Mohr-Coulomb materials, for brittle rock following Griffith's theory, and for progressive deformation and fracture of rock around underground openings under stress changes associated with excavation.

At this reporting, selection of mathematical models for elastic-plastic Mohr-Coulomb materials and for elastic-brittle materials failing according to Griffith theory has been completed. Chapter 1 of the report describes the theoretical considerations leading to the model selected. The stress-strain relations for incremental or rate type theory of plasticity are generally based on the normality rule and convexity and regularity of the yield surface in a 'stress-space'. Using these

concepts, various investigators have proposed conflicting constitutive equations. In this report the elastic-plastic behavior of materials has been re-examined as a mathematical generalization of observations on a one-dimensional test. The role of kinematic constraints upon yield conditions has been studied and adequacy of certain postulates examined. Current theories of elastic-plastic behavior are found to be inadequate as it is, in general, not possible to satisfy the 'normality' rule as well as continuity of stress path under plane strain conditions. Further research into this aspect of material behavior is needed to clear the air. Experimental phase of the research program is being planned with this requirement in view. In the mathematical model selected as the basis for development of computer programs, a modification of the yield surface is introduced to eliminate discontinuity in stress paths. For behavior of elastic-brittle rock, the model selected assumes elements of rock to be incapable of supporting tensile and shearing forces across a crack. A review of literature showed errors in similar formulations by other investigators. These have been corrected in the present development. The mathematical models of elastic-plastic and elastic-brittle rock have been incorporated into finite element computer programs for analyses of stresses and deformations of plane strain systems. Chapters III and IV of the report present two computer codes along with relevant description, instructions for usage and illustrative examples for:

- i. Plane Strain Analysis of Elastic-Plastic Mohr-Coulomb Materials
- ii. Two-Dimensional Analysis of a Non-Tension System

Further work on these computer programs is continuing. However, even in their present form, program capabilities include consideration of arbitrary initial stresses, arbitrary shape of openings with or without linings, and considerable variation in material properties. These computer programs should be of immediate application to a variety of problems.

Adequate mathematical models of rock behavior have been chosen. The finite element method has been used successfully to develop computer codes for analysis of complex problems of stresses, deformations and fracture in rock. The method appears to be suitable for further development to realize the objectives of the current research program.

Experimental work so far has been directed towards development of suitable laboratory material (exhibiting elastic-plastic behavior). No equipment has so far been purchased under the contract. However, procurement of a plane strain testing machine has been initiated. It is expected to be received in September 1971.

PREFACE

The terrestrial crust is in a complex state of stress. Underground excavations in this stressed medium profoundly influence the distribution of stress which in turn determines the stability of the opening and of the rock in the vicinity. Traditional methods based upon the theory of linear elastic solids are inadequate. It is necessary that the sequence of construction and realistic material properties be taken into account in calculation of stresses and deformations in rock.

The objective of the present research program is development of finite element techniques to predict stresses and deformations in the vicinity of underground excavations allowing for arbitrary initial stresses in the rock, arbitrary shape and size of the opening and any given sequence of construction. The procedures will allow for nonhomogeneous material properties, interaction of rock with supporting structures, progressive damage, time dependent deformation and load development on supporting structures. Limited experimental work to verify key points in the theory is also planned. The entire program is expected to extend over three years.

This is the first semi-annual progress report covering the period 2/1/71 to 7/31/71. The main activity in this period has been a review of the work done by other investigators in mathematical simulation of mechanical behavior of elastic-plastic solids and of jointed rock. Models of stress-strain behavior

have been proposed and work on development of relevant computer programs started. A sequential approach has been followed whereby a basic program is coded and then modified to include all the ramifications of material behavior and actual loading sequences. Two computer programs, viz.

i. Plane Strain Analysis of Elastic-Plastic Mohr-Coulomb Materials

ii. Two-Dimensional Analysis of a No-Tension System

are included. The present capabilities of each program are indicated in the program descriptions. Further development on all these is continuing and will be included in future reports.

The work is supported by the U.S. Government through the Advanced Research Projects Agency, ARPA, and its agent the Bureau of Mines, Department of the Interior. At the Ohio State University the work is under direct supervision of Professors T.H. Wu, R.S. Sandhu, and J.R. Hooper. Messrs. S.W. Huang, R.D. Singh, C.W. Chang and T. Chang, graduate students in the Department of Civil Engineering, have contributed to the research reported. Dr. William Karwowski of the Spokane Mining Research Center, Spokane, Washington is the Project Officer designated by the sponsor. In early stages Dr. Syd Peng of Twin Cities Mining Research Center, Twin Cities, Minnesota acted as the Project Officer.

The opinions, findings and conclusions expressed in the report are those of the authors and not necessarily those of the U.S. Bureau of Mines, Department of the Interior or the Advanced Research Project Agency.

R. S. Sandhu
Project Supervisor

TABLE OF CONTENTS

	<u>Page</u>
Technical Report Summary	i
Preface	iii
List of Figures	vi
Chapter I. Theoretical Considerations	1
1.1 Mechanical Behavior of Rock	
1.2 Stress-Strain Relation in Elastic-Plastic Solids	
1.3 Stress-Strain Behavior of Jointed Rock	
Chapter II. The Finite Element Method	26
2.1 Basic Concepts	
2.2 A Potential Energy Formulation	
2.3 Incremental Analysis	
Chapter III. Computer Program for Plane Strain Analysis of Elastic-Plastic Mohr-Coulomb Materials	32
3.1 Program Capability and Organization	
3.2 Input Data Preparation	
3.3 Program Listing	
3.4 Illustrative Example	
Chapter IV. Computer Program for Analysis of Jointed Rock	66
4.1 Program Capability and Organization	
4.2 Input Data Preparation	
4.3 Program Listing	
4.4 Illustrative Example	
Chapter V. Additional Comments	96
List of References	97

LIST OF FIGURES

<u>Figure</u>	<u>Title</u>	<u>Page</u>
1-1	Stress-Strain Curves for Cedar City Granite	2
1-2	Stress-Strain Curves for Marble	2
1-3	Idealization of Elastic-Plastic Stress-Strain Behavior for Rocks	4
1-4	Typical Axial and Lateral Stress-Strain Behavior of Brittle Rock	4
1-5	Proportional Limit in Shear for Westerly Granite	5
1-6	Effect of Confining Pressure on Stress-Strain Behavior of Limestone	5
1-7	Plane Strain Mohr-Coulomb Criteria	19
3-1	Calculation of Stress-Ratio in Elastic-Plastic Computer Program	61
3-2	Finite Element Idealization for Elastic-Plastic Wedge	62
3-3	Distribution of Radial Stress for Wedge at Various Loads	63
3-4	Distribution of Circumferential Stress for Wedge at Various Loads	64
3-5	Distribution of Radial Displacements	65
4-1	Finite Element Idealization for Lined Tunnel	93
4-2	Elastic Solution of the Lined Tunnel	94
4-3	Final 'No Tension' Solution of the Lined Tunnel	95

CHAPTER I
THEORETICAL CONSIDERATIONS

Chapter I. Theoretical Considerations

1.1. Mechanical Behavior of Rock

Figs. 1-1 and 1-2 show, respectively, typical stress-strain plots for a granite and a marble (Swanson 1970). Upon loading the stress-strain curve is almost linear and reversible over a short portion. Unloading from higher loads does not coincide with initial loading. This characteristic along with rate independence distinguishes elastic-plastic behavior. Reloading closely follows unloading until the previous maximum is reached; whereupon the original curve is followed. This leads to some simplifying assumptions.

- i. A yield point exists below which the material is linear elastic.
- ii. The yield point corresponds to the maximum stress level previously attained.
- iii. Unloading and reloading paths are linear, coincident and parallel to the initial elastic loading curve.

Fig. 1-3 shows this simplification. Clearly the yield point can be described by the permanent or irrecoverable strain or the area bounded by the loading curve, the unloading curve and the horizontal axis. Whereas in generalization to the three-dimensional case, the stresses and strains become second rank tensors and are, therefore, unordered, the area is still a scalar product and retains its ordering characteristics. To this extent, it is often preferred as a measure of the elastic limit.

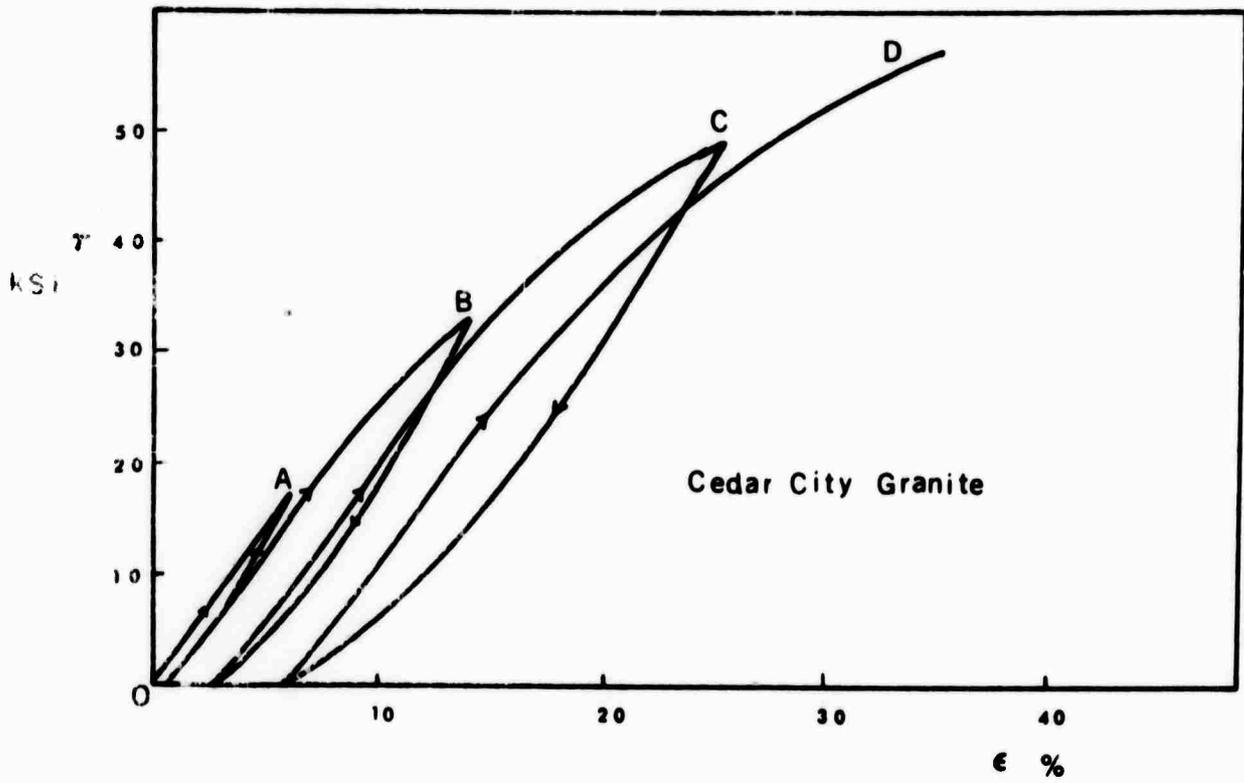


Figure 1-1. Stress-Strain Curves for Cedar City Granite

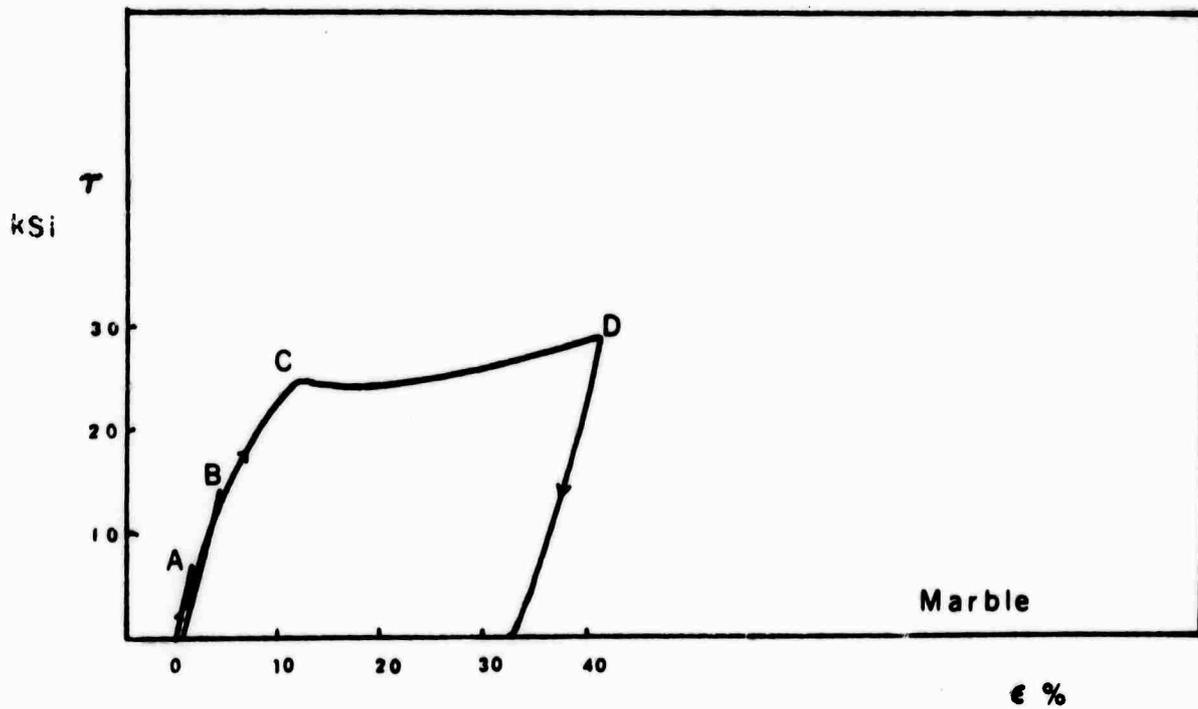


Figure 1-2. Stress-Strain Curves for Marble

Mechanical behavior of rock under polyaxial state of stress has been examined in the light of brittle failure theories (Brace, 1964; Bieniawski, 1967, 1969; Brady, 1969, 1970). Four regions of behavior are identified in Figure 1-4. The first region corresponds to closure of pre-existing open cracks and is peculiar to compressional loading. In region II material behavior is linear elastic. Fracture initiation occurs near the end of this region in accordance with Griffith or modified Griffith Theory. This stage also corresponds to onset of nonlinearity in the stress to volumetric strain curve (Brace, 1966). Stable fracture propagation characterizes region III. In region IV, unstable fracture propagation results in strength failure and rupture. Differences in loading and unloading behavior are observed (Walsh, 1965).

We have, thus, two general approaches to the characterization of stress-strain behavior of rock. One follows the theory of elastic-plastic solids without consideration of micro-mechanics of the system. The other uses Griffith theory or modified Griffith theory to relate deformation and failure to initiation and propagation of fracture. It has been observed (Swanson 1970) that Mohr-Coulomb failure law applies for moderate values of confining pressure and that at low confining pressures, failure is by rupture. Contrary to plastic behavior, strength of material drops to almost zero in the direction normal to the crack if rupture theory is followed. Figs. 1-5 and 1-6 depict typical relationships of failure strength and post failure behavior in relation to confining pressures. It is reasonable to assume that the material is linear elastic upto yield or rupture, as the case may

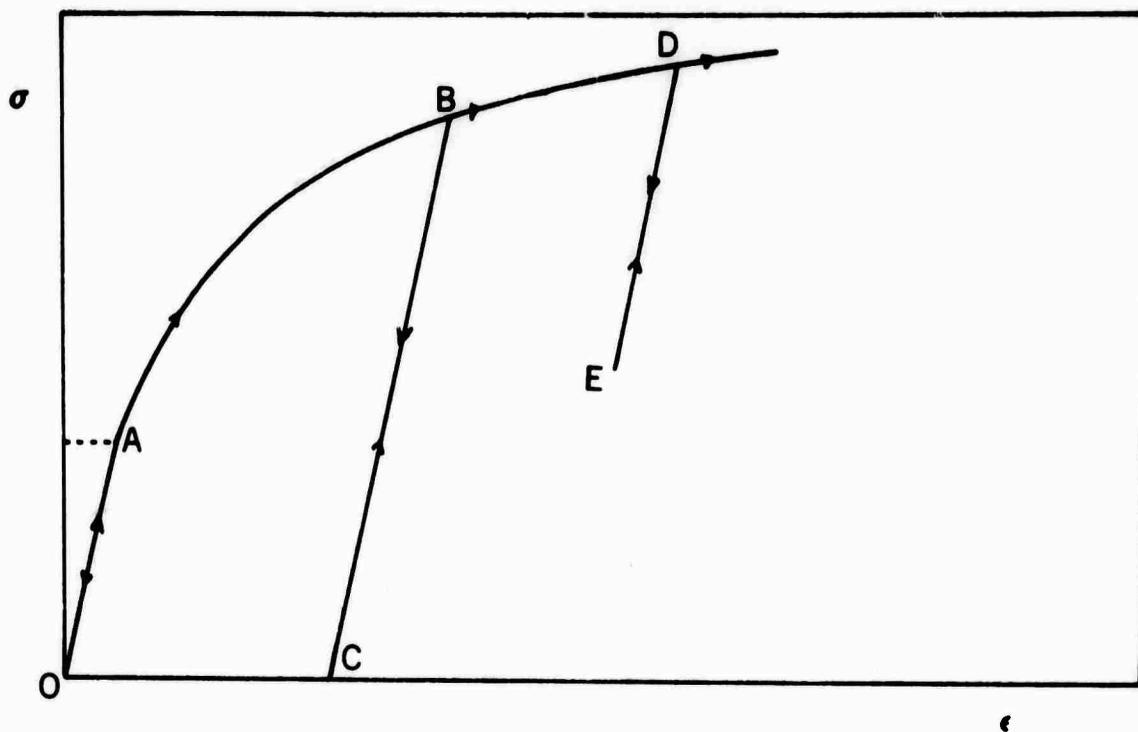


Figure 1-3. Idealization of Elastic-Plastic Stress-Strain Behavior for Rocks

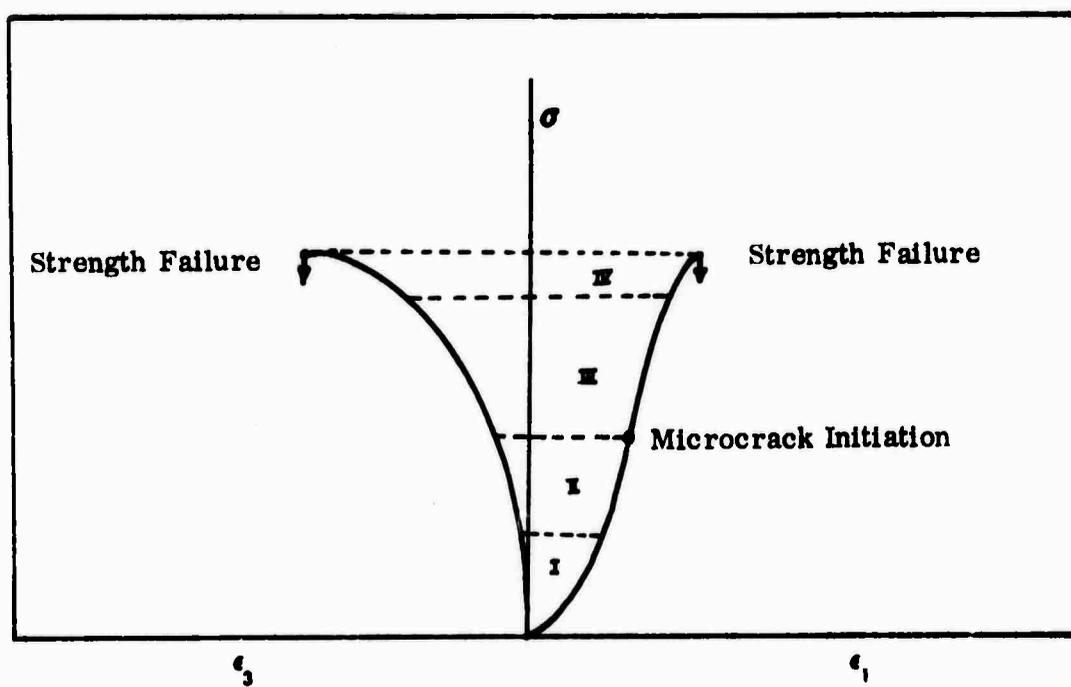


Figure 1-4. Typical Axial and Lateral Stress-Strain Behavior of Brittle Rock

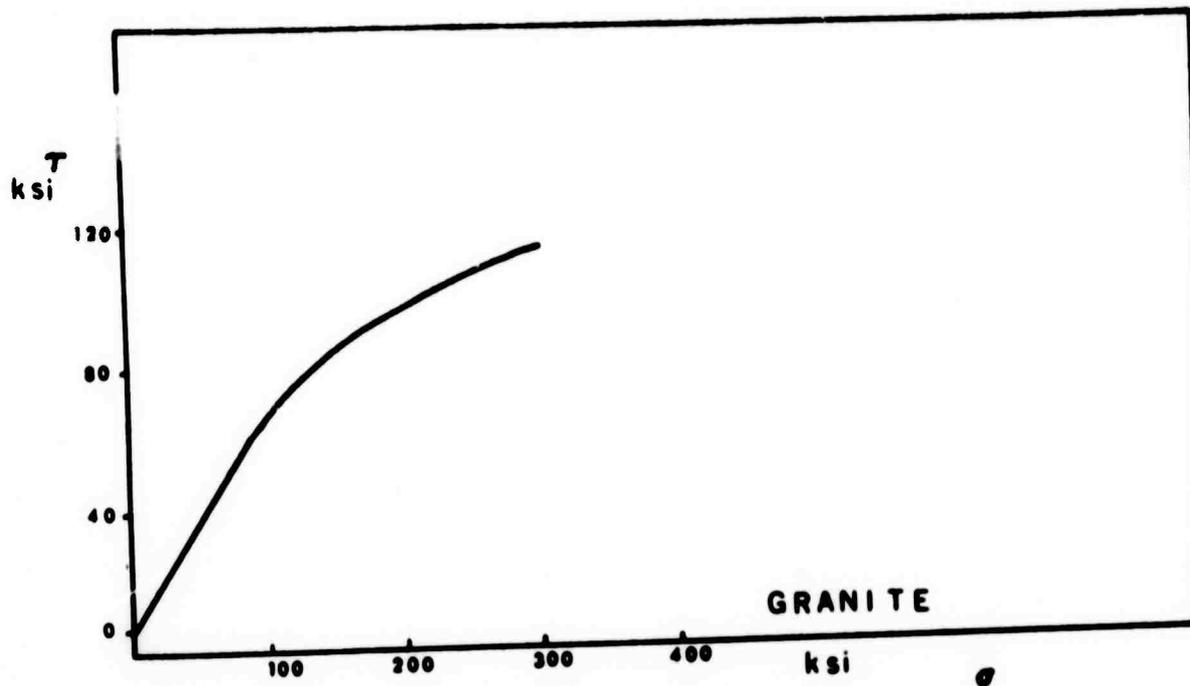


Figure 1-5. Proportional Limit in Shear for Westerly Granite

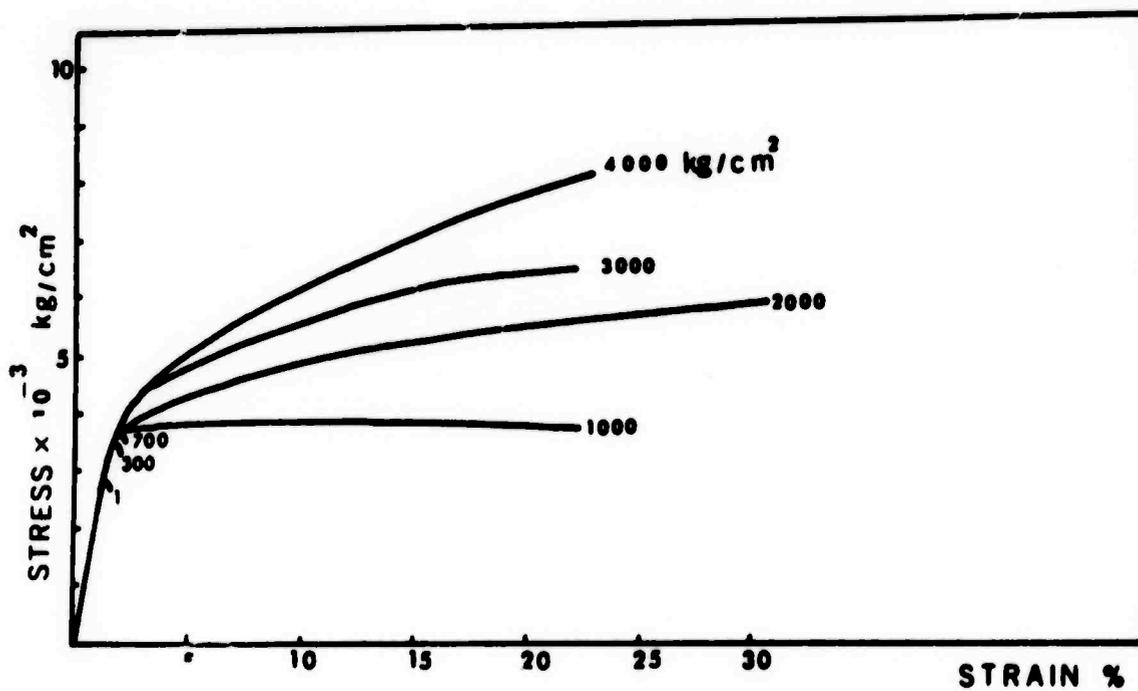


Figure 1-6. Effect of Confining Pressure on Stress-Strain Behavior of Limestone

be, and that the post-failure behavior only is governed by the theory used to define failure. Both the elastic-plastic Mohr-Coulomb failure theory and the Griffith theory have been used in the course of present research to develop computer programs for analysis of stress and deformation in rock.

1.2 Stress-Strain Relations for Elastic-Plastic Solids

Several approaches have been used for formulation of elastic-plastic behavior. Excellent presentations of the theory are available in literature (Drucker, 1951; Naghdi, 1960; Green and Naghdi, 1965; Koiter, 1953; Hill, 1950). Specializations to Mohr-Coulomb solids under plane strain (Drucker and Prager, 1952; Drucker, Gibson and Henkel, 1955; Reyes, 1965; Reyes & Deere, 1966) have been presented. Not intending to survey the entire range of different philosophies, we present briefly a discussion of elastic-plastic behavior under kinematic constraints leading to the plane strain formulation for Mohr-Coulomb materials used in the computer program in Chapter III. We present a generalization of observations on the conventional uniaxial test to the case of three-dimensions before discussing the role of constraints.

An isothermal system undergoing infinitesimal deformations is of interest to the present report. Extension to more comprehensive situations is direct.

1.2.1. A Generalization of One-Dimensional Test Results to Three-Dimensional Theory

Analogous to the case of one-dimensional test, we assume the existence

of a set D consisting of all admissible states of stress. D is convex,* has boundary B and includes the original or reference state. Clearly D is six-dimensional (assuming symmetric stress tensors, i.e. absence of body couples) and is contained in the six-dimensional linear vector space V spanned by the six components σ_{kl} of the stress tensor. The stress in uniaxial test is given by a real number and in this case D is ordered and convex. To introduce an ordering in the six-dimensional stress space so that 'increase' and 'decrease' of stress are meaningful, a mapping g is defined

$$g: V \longrightarrow R \quad (I-1)$$

with the following properties

- i. The image of D under g is a positive interval $I \subset R$.
- ii. Image of complement of D in V is the complement of I in R and g maps the interior D_I of D to the interior of I . (I-2)
- iii. g maps the 'original' or 'reference' state to zero $\in I$.

Then D is ordered by its image i.e. for $\nu_1, \nu_2 \in D$, if $g(\nu_1) = g_1, g(\nu_2) = g_2, \nu_2$ is greater than, equivalent**to, or less than ν_1 depending upon g_2 being greater than,

*Convexity of D is the property that $\nu_1, \nu_2 \in D \Rightarrow a\nu_1 + (1-a)\nu_2 \in D \forall a \in [0, 1]$. In literature, there are frequent references to a convex yield surface. This is inaccurate. It is easy to see that convexity of D does not imply convexity of its boundary B . Indeed, B is in general not convex.

**The term equivalent is used because g , in general, is not one to one. Thus, ν_1, ν_2 may be distinct while their images coincide.

equal to or less than g_1 . The mapping preserves convexity of D . The interval $I = [0, f]$ where f is the image of boundary B of D .

In one-dimensional tests, the set of admissible stress states may include negative points. The limiting stress states in tension and compression provide a positive supremum and a negative infimum to this set. The boundary B of this set is clearly discontinuous. In a multi-dimensional stress space, B may be continuous. To ensure correspondence between the one-dimensional and multi-dimensional cases, g is a two to one mapping in the case of one-dimensional loading. Thus the image of B is in all cases the supremum of non-negative interval I .

The boundary B of D and hence its image f under g are defined by prior deformation and load history. Considering components ϵ_{kl} of the strain tensor, on the analogy of the results of one-dimensional test, a plastic strain tensor with components ϵ''_{kl} is defined such that

- i. For a given ϵ''_{kl} , there is one to one correspondence between elements of D and a set of points in the six-dimensional space spanned by $\epsilon'_{kl} = \epsilon_{kl} - \epsilon''_{kl}$. ϵ'_{kl} are identified as components of the elastic strain tensor. (I-3)
- ii. If a generalization of Prandtl's simplifying assumption is admitted, the one to one correspondence between $\sigma_{kl} \in D$ and ϵ'_{kl} is independent of prior deformation history, and for stress states defined by interior D_I of D , ϵ''_{kl} vanish.

A positive measure of history of deformation can be defined in various ways. If plastic strain is used as representative of deformation history, a

mapping χ on H the six-dimensional linear vector space spanned by components ϵ''_{kl} of the plastic strain tensor is introduced

$$\chi : H \rightarrow P \quad (I-4)$$

where P is the positive class of real numbers. Other measures using bilinear or nonlinear maps involving both the stress and plastic strain components are in use. An example is*

$$\chi(\sigma_{kl}, \epsilon''_{kl}) = \int_{\tau=-\infty}^t \sigma_{kl}(\tau) d\epsilon''_{kl}(\tau) \quad (I-5)$$

In all cases the objective is to define a positive number k such that it equals the image f under g of boundary B of D . For elastic-perfectly plastic solids, k is constant but, in general, for stable material, k is a monotone increasing function of history of deformation and stress. In certain cases, the mapping g may also vary with plastic deformation. This happens when kinematic constraints are present. Theory of kinematic hardening is an instance in which the reference point in D , having image zero in I , depends upon stress and deformation history.

Considering, for the present discussion,** $\chi(\epsilon''_{ij}) = k$,

$$g(\sigma_{kl}) \leq f = k = \chi(\epsilon''_{kl}) \quad (I-6)$$

*Here and in subsequent work, standard indicial notation is used. Summation on repeated indices is implied unless otherwise indicated.

**A more general assumption uses $\chi = \chi(\kappa, \epsilon''_{ij})$ where $\kappa = \int_{\tau=-\infty}^t \sigma_{ij}(\tau) d\epsilon''_{ij}(\tau)$.
 In that case $\dot{\chi} = \frac{\partial \chi}{\partial \kappa} \dot{\kappa} + \frac{\partial \chi}{\partial \epsilon''_{ij}} \dot{\epsilon}''_{ij}$.

In the interior of D , $\epsilon''_{kl} = 0$, $g(\sigma_{kl}) < f = k$ and

$$\begin{aligned}\sigma_{kl} &= E_{kl}(\epsilon'_{mn}) \\ &= E_{kl}(\epsilon_{mn})\end{aligned}\quad (I-7)$$

For differential changes in stress and strain components, using a superposed dot to indicate differential quantities,

$$\begin{aligned}\dot{\sigma}_{kl} &= E_{klmn} \dot{\epsilon}'_{mn} \\ &= E_{klmn} \dot{\epsilon}_{mn}\end{aligned}\quad (I-8)$$

assuming that E_{kl} is sufficiently smooth and its derivative E_{klmn} , a tensor of fourth rank, exists.

On the boundary B of D ,

$$g(\sigma_{kl}) = f = k = \chi(\epsilon''_{kl})\quad (I-9)$$

For g , χ sufficiently smooth in their arguments

$$\dot{\chi} = \dot{\chi}(\epsilon''_{kl}) = h_{kl} \dot{\epsilon}''_{kl}\quad (I-10)$$

$$\dot{g} = \dot{g}(\sigma_{kl}) = q_{kl} \dot{\sigma}_{kl}\quad (I-11)$$

In the case of elastio-perfectly plastic solids, $h_{kl} = 0$ and arbitrary plastic straining can occur for $\dot{\chi} = 0$ i.e. $\chi = k$, a positive constant. Also $g = f = k$ requires $\dot{g} = 0$ leading to the relationship

$$q_{kl} \dot{\sigma}_{kl} = 0\quad (I-12)$$

Equation I-12 requires the stress changes to be in a plane tangent to the hyperplane defined by $g(\sigma_{kl}) = k$. However, for $h_{kl} \neq 0$, for nonvanishing $\dot{\epsilon}''_{kl}$, $\dot{\chi} > 0$ and $\dot{g} = \dot{f} > 0$. This is termed loading and $\dot{g} = \dot{f} = \dot{k} = \dot{\chi}$. For vanishing $\dot{\epsilon}''_{kl}$, $\dot{\chi} = 0$ and

once again equation I-12 applies. This is the case of neutral loading. In all cases $\dot{g} < 0$ implies decreasing load. This is the case of unloading and equation I-8 applies with $\epsilon''_{kl} = 0$.

1.2.2 Evaluation of Incremental Plastic Strain in Loading

Equations I-10 and I-11 suggest a relationship of the type

$$\dot{\epsilon}''_{kl} = S_{klmn} \dot{\sigma}_{mn} \quad (I-13)$$

where S_{klmn} may depend upon ϵ''_{mn} , σ_{mn} . Resolving $\dot{\sigma}_{mn}$ into components along the boundary B and normal to it, the plastic strain is due only to the normal component. Prager (1949) showed that

$$\begin{aligned} S_{klmn} &= S_{kl ij} \frac{\left(\frac{\partial g}{\partial \sigma_{ij}} \right)}{\left(\frac{\partial g}{\partial \sigma_{pq}} \right) \left(\frac{\partial g}{\partial \sigma_{pq}} \right)} \cdot \left(\frac{\partial g}{\partial \sigma_{mn}} \right) \\ &= B_{kl} \frac{\partial g}{\partial \sigma_{mn}} \end{aligned} \quad (I-14)$$

independent of $\dot{\sigma}_{mn}$. Hence direction of ϵ''_{ij} is independent of the direction of stress change given by $\dot{\sigma}_{mn}$. Other relations for plastic strain increments have been proposed. Using a thermo-dynamic postulate, Drucker (1951) obtained the normality rule

$$\epsilon''_{ij} = \lambda \frac{\partial g}{\partial \sigma_{ij}} \quad \text{at } g = f \quad (I-15)$$

where λ is a positive scalar which for rate independence must be homogeneous of order one in $\dot{\sigma}_{kl}$.

Drucker used the normality rule to evaluate λ . Defining

$$\dot{\epsilon}''_e = \left[\frac{1}{2} \dot{\epsilon}''_{ij} \dot{\epsilon}''_{ij} \right]^{\frac{1}{2}} \quad (I-16)$$

$$\lambda = \frac{\dot{\epsilon}''_e}{\left[\frac{1}{2} \frac{\partial g}{\partial \sigma_{ij}} \frac{\partial g}{\partial \sigma_{ij}} \right]^{\frac{1}{2}}} \quad \text{at } g = f \quad (I-17)$$

Now defining $\sigma_e = \left[\frac{1}{2} S_{ij} S_{ij} \right]^{\frac{1}{2}} \quad (I-18)$

where $S_{ij} = \sigma_{ij} - \delta_{ij} \frac{\sigma_{kk}}{3} \quad (I-19)$

and writing $\frac{\dot{\sigma}_e}{\dot{\epsilon}''_e}$, the slope of the σ_e, ϵ''_e curve as H, equation I-17 gives, with the normality rule

$$\begin{aligned} \dot{\epsilon}''_{ij} &= \frac{\frac{\partial g}{\partial \sigma_{ij}} \dot{\sigma}_e}{H \left[\frac{1}{2} \frac{\partial g}{\partial \sigma_{mn}} \frac{\partial g}{\partial \sigma_{mn}} \right]^{\frac{1}{2}}} \\ &= \frac{\frac{\partial g}{\partial \sigma_{ij}} S_{kl} \dot{S}_{kl}}{2 \sigma_e H \left[\frac{1}{2} \frac{\partial g}{\partial \sigma_{mn}} \frac{\partial g}{\partial \sigma_{mn}} \right]^{\frac{1}{2}}} \end{aligned} \quad (I-20)$$

This formulation was used in the so-called tangent modulus methods e.g. Swedlow and Yang (1965).

Hill used the normality rule assuming λ to be a fourth rank tensor linear in $\dot{\sigma}_{kl}$ and introduced a plastic potential. Using normality as well as the condition $\dot{g} = \dot{f} = \dot{k} = \dot{X}$ on the boundary B of D, Prager (1949) obtained

$$\lambda = \frac{\frac{\partial g}{\partial \sigma_{ij}} \dot{\sigma}_{ij}}{\frac{\partial X}{\partial \epsilon''_{mn}} \frac{\partial g}{\partial \sigma_{mn}}} \quad (I-21)$$

This formulation breaks down for elastic-perfectly solids where χ is independent of ϵ''_{mn} . Felippa (1966) obtained λ in terms of $\dot{\epsilon}_{ij}$, increment in the total strain tensor. In this procedure

$$\dot{\sigma}_{ij} = \left[E_{ijkl} \delta_{km} \delta_{ln} - E_{ijkl} L_{klmn} \right] \dot{\epsilon}_{mn} \quad (I-22)$$

where

$$L_{klmn} = \left[\frac{\partial \chi}{\partial \epsilon''_{ij}} \frac{\partial g}{\partial \sigma_{ij}} + \frac{\partial g}{\partial \sigma_{ij}} E_{ijpq} \frac{\partial g}{\partial \sigma_{pq}} \right]^{-1} \frac{\partial g}{\partial \sigma_{rs}} E_{rskl} \frac{\partial g}{\partial \sigma_{mn}} \quad (I-23)$$

This approach is valid for all cases including perfect plasticity and was used by Zienkiewicz, Valliappan and King (1969) in developing finite element procedures.

Using rate of work equations, it is possible to evaluate λ in terms of stress rates for materials of von Mises or Mohr-Coulomb type. Yamada (1968) used this approach for finite element analysis of von Mises solids. Using Drucker and Prager's generalization of Mohr-Coulomb law, Reyes (1965) developed the stress-strain equation for generalized Mohr-Coulomb elastic-perfectly plastic solids under plane strain conditions. The finite element procedures presented by Reyes and Deere (1966), Baker, Sandhu and Shieh (1969) and those included in Chapter III of this report were based on these equations. For plane strain

$$\begin{pmatrix} \dot{\sigma}_{11} \\ \dot{\sigma}_{22} \\ \dot{\sigma}_{12} \end{pmatrix} = 2G \begin{bmatrix} D_{11} & D_{12} & D_{13} \\ D_{21} & D_{22} & D_{23} \\ D_{31} & D_{32} & D_{33} \end{bmatrix} \begin{pmatrix} \dot{\epsilon}_{11} \\ \dot{\epsilon}_{22} \\ \dot{\gamma}_{12} \end{pmatrix} \quad (I-24)$$

where

$$D_{11} = 1 - h_2 - 2 h_1 \sigma_{11} - h_3 \sigma_{11}^2$$

$$D_{22} = 1 - h_2 - 2 h_1 \sigma_{22} - h_3 \sigma_{22}^2$$

$$\begin{aligned}
D_{33} &= \frac{1}{2} - h_3 \sigma_{12} \\
D_{12} = D_{21} &= -h_2 - h_1 \sigma_{22} - h_1 \sigma_{11} - h_3 \sigma_{22} \sigma_{11} \\
D_{13} = D_{31} &= -h_1 \sigma_{12} - h_3 \sigma_{12} \sigma_{11} \\
D_{23} = D_{32} &= -h_1 \sigma_{12} - h_3 \sigma_{12} \sigma_{22}
\end{aligned} \tag{I-25}$$

and

$$\begin{aligned}
h_1 &= \frac{\frac{3}{2} \alpha \frac{K}{G} - \frac{\sigma_{kk}}{6 J_2^{\frac{1}{2}}}}{J_2^{\frac{1}{2}} \left(1 + 9 \alpha^2 \frac{K}{G}\right)} \\
h_2 &= \frac{\alpha - \frac{\sigma_{kk}}{6 J_2^{\frac{1}{2}}} - 3 \alpha \frac{K}{G} - \frac{\sigma_{kk}}{3 J_2^{\frac{1}{2}}}}{\left(1 + 9 \alpha^2 \frac{K}{G}\right)} - \frac{3 \nu f K}{J_2^{\frac{1}{2}}} \frac{1}{E \left(1 + 9 \alpha^2 \frac{K}{G}\right)} \\
h_3 &= \frac{1}{2 J_2^{\frac{1}{2}} \left(1 + 9 \alpha^2 \frac{K}{G}\right)}
\end{aligned}$$

$$J_2 = \frac{1}{2} S_{ij} S_{ij}$$

E, K, G = elastic Young's modulus, bulk modulus and shear modulus, respectively.

$$\alpha = \frac{\tan \phi}{\sqrt{9 + 12 \tan^2 \phi}}$$

ϕ = the angle of internal friction

$$f = \alpha \frac{\sigma_{kk}}{3} + J_2^{\frac{1}{2}}$$

1.2.3. Kinematic Constraints

Plane strain conditions impose a kinematic constraint upon the deforming solid. In relation to elastic-plastic behavior, a consequence is that the yield sur-

surface from the elastic side and plastic side do not, in general, coincide (Baker et.al,1969). Consider the deformation of a body undergoing deformation. F , the set of all admissible deformation, is contained in the six-dimensional vector space S spanned by components ϵ_{kl} of the strain tensor. A kinematic constraint can be written as

$$C(\epsilon_{kl}) = 0 \quad (I-26)$$

and the admissible deformation is restricted to the intersection of F with the hyperplane in S defined by equation I-26. If several constraints are present, the admissible deformation is restricted to

$$\bigcap_{i=1}^n [C_i(\epsilon_{kl}) = 0] \quad (I-27)$$

As the multiple intersection reduces the dimension of the vector space by n , it is clear that n cannot exceed six.

Consider a single constraint. In differential form the equation is

$$C_{kl} \dot{\epsilon}_{kl} = 0 \quad (I-28)$$

where coefficients C_{kl} depend upon ϵ_{mn} .

As elastic-plastic behavior is studied with reference to loading paths in the stress space V , it is necessary that kinematic constraints be rewritten as constraints on stress. Here, for no plastic strain, we simply use the inverse of equation I-8 to write equation I-27 as

$$C_{kl} C_{klmn} \dot{\sigma}_{mn} = 0 \quad (I-29)$$

or

$$G_{mn} \dot{\sigma}_{mn} = 0 \quad (I-30)$$

where $C_{klmn} = [E_{klmn}]^{-1}$ (I-31)

and $G_{mn} = C_{kl} C_{klmn}$ (I-32)

For the case of not all of $\dot{\epsilon}''_{kl}$ vanishing, two alternative procedures are available.

Using definition of elastic strain tensor,

$$\dot{\epsilon}_{kl} = \dot{\epsilon}'_{kl} + \dot{\epsilon}''_{kl} \quad (I-33)$$

If T_{klmn} is the inverse of S_{klmn} in equation I-13

$$\dot{\epsilon}_{kl} = C_{klmn} \dot{\sigma}_{mn} + T_{klmn} \dot{\sigma}_{mn} \quad (I-34)$$

$$= [I_{klpq} + T_{klmn} E_{mnpq}] \dot{\epsilon}'_{pq} \quad (I-35)$$

$$= K_{klpq} \dot{\epsilon}'_{pq} \quad (I-36)$$

where I_{klpq} is a fourth rank identity tensor. Thus the constraint is expressed by

$$C_{kl} \dot{\epsilon}_{kl} = C_{kl} K_{klpq} \dot{\epsilon}'_{pq} \quad (I-37)$$

$$= C_{kl} K_{klpq} C_{pqmn} \dot{\sigma}_{mn} \quad (I-38)$$

$$= L_{klmn} \dot{\sigma}_{mn} = 0 \quad (I-39)$$

where $L_{klmn} = C_{kl} [I_{klpq} + T_{klrs} E_{rspq}] C_{pqmn}$ (I-40)

An alternative procedure is to use the normality rule and to satisfy the constraint

both upon loading and unloading i. e. $C_{kl} \dot{\epsilon}'_{kl} = 0 = C_{kl} \dot{\epsilon}''_{kl}$. Then the first equation

is identical with equation I-29 but the second equation gives

$$C_{kl} \lambda \frac{\partial g}{\partial \sigma_{kl}} = 0 \quad (I-41)$$

Equation I-41 may or may not coincide with equation I-39. Equations I-29 and I-39

have linear relationship between incremental stresses and describe hyperplanes

tangent to any loading path in the stress space V . As the two equations are in general different, there is a slope discontinuity in the stress path as plastic straining begins upon reaching the boundary B of D . We note in particular that proportional stress paths in V may not be possible in the presence of kinematic constraints. In the case of linear elasticity, let equation I-30 define a plane passing through the origin in V . A proportional loading path lying in this plane is possible upto the point of intersection with boundary B . Beyond that, upon loading, stress path has to be in the surface determined by Equation I-39 and this will in general be non-planar. If loading is continued to a certain point along this surface, unloading therefore will be along a path lying in a plane parallel to the original loading plane but different from it and not passing through the origin. Thus unloading to initial state is impossible. This corresponds to setting up of residual stresses corresponding to kinematic constraints.

Specifically considering plane strain conditions and elastic-perfectly plastic Mohr-Coulomb material, the mapping g from the set of all admissible stress states to the positive interval $[0, f]$ is

$$g(\sigma_{ij}) = \alpha \frac{\sigma_{kk}}{3} + J_2^{\frac{1}{2}} \quad (I-42)$$

$$J_2 = \frac{1}{2} \sigma_{ij} \sigma_{ij} - \frac{\sigma_{kk}^2}{6} \quad (I-43)$$

Linear isotropic elasticity implies, for $\dot{\epsilon}''_{ij} = 0$

$$\epsilon_{ij} = \epsilon'_{ij} = \frac{1}{2G} \sigma_{ij} - \frac{\nu \epsilon_{kk}}{E} \sigma_{kk} \quad (I-44)$$

Plane strain condition implies

$$\epsilon_{13} = \epsilon_{23} = \epsilon_{33} = 0 \quad (I-45)$$

From equations I-44 and I-45

$$\sigma_{13} = \sigma_{23} = 0 \quad (I-46)$$

$$\sigma_{33} = \nu (\sigma_{11} + \sigma_{22}) \quad (I-47)$$

For $\dot{\epsilon}_{ij}''$, not all vanishing, using normality rule

$$\dot{\epsilon}_{ij}'' = \lambda \left[\frac{\alpha}{3} \delta_{ij} + \frac{\sigma_{ij} - \frac{1}{3} \delta_{ij} \sigma_{kk}}{2 J_2^{\frac{1}{2}}} \right]$$

$$\frac{\alpha}{3} + \frac{\sigma_{33} - \frac{\sigma_{kk}}{3}}{2 J_2^{\frac{1}{2}}} = 0$$

$$\frac{2}{3} \sigma_{33} = - \frac{2\alpha J_2^{\frac{1}{2}}}{3} + \frac{\sigma_{11} + \sigma_{22}}{3}$$

or
$$\sigma_{33} = -\alpha J_2^{\frac{1}{2}} + \frac{1}{2} (\sigma_{11} + \sigma_{22}) \quad (I-48)$$

Equations I-47 and I-48 define different surfaces in a three-dimensional space spanned by σ_{11} , σ_{22} , σ_{33} . Let their intersections with B be respectively, P and Q. The stress path is constrained to lie in the plane defined by equation I-47 for stress states in the interior of D and for neutral loading. For plastic deformation to occur the stress path must lie in the surface defined by equation I-48. For a continuous stress path to be possible, P and Q must coincide. In general this is not the case. Figure I-7 illustrates the difference between the surfaces P and Q for Mohr-Coulomb plane strain case.

Considering that the stresses σ_{33} do not contribute to energy/work of the system, it appears reasonable to assume that progress from P to Q is possible with gradually increasing the value of σ_{33} . This would amount to following the

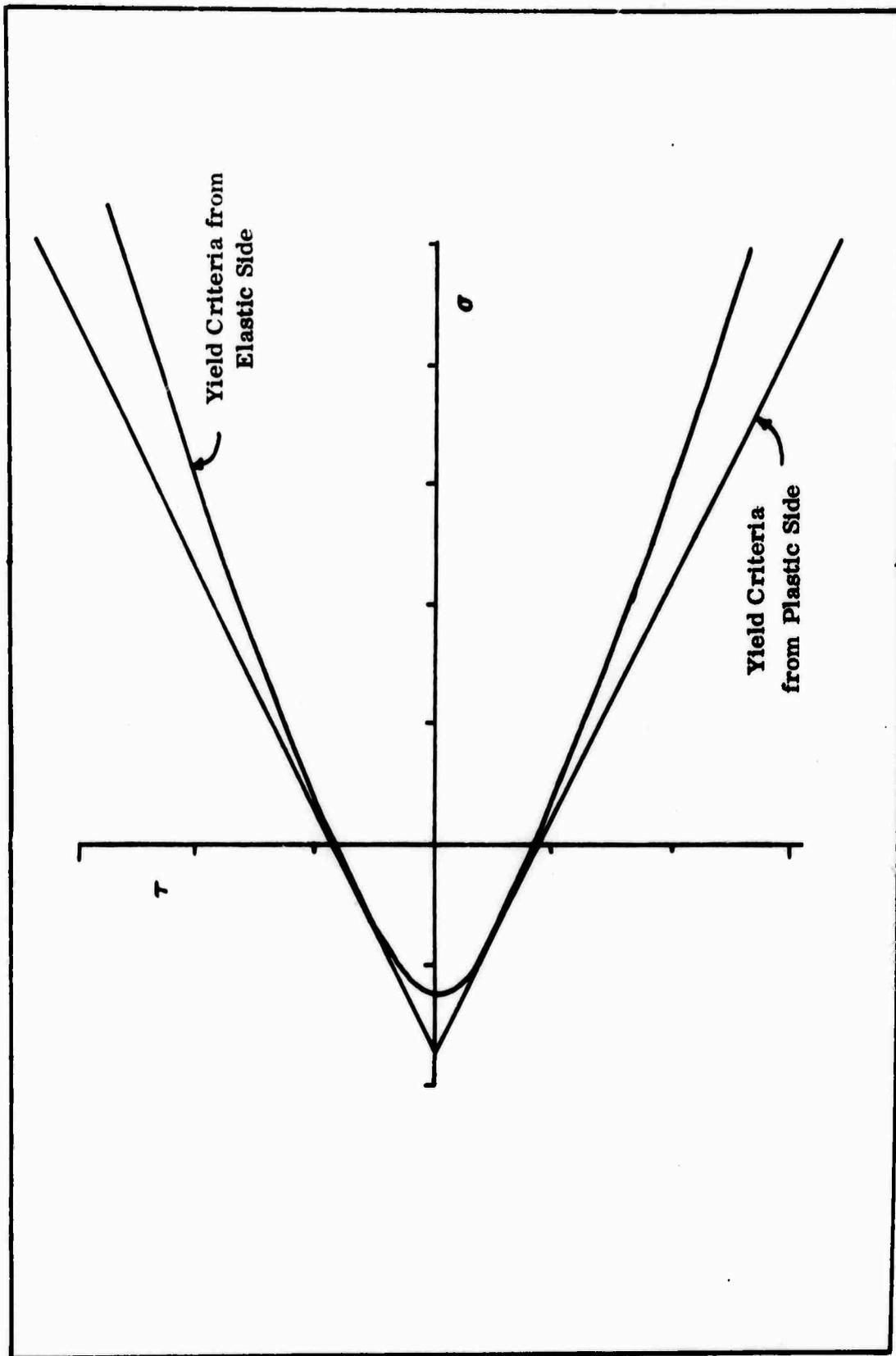


Figure 1-7. Plane Strain Mohr-Coulomb Criteria

boundary B. Growth of σ_{33} in progress from P to Q and the behavior upon unloading are not clearly understood. Later investigations may throw light on this aspect of material behavior. For the purpose of the computer program in Chapter III, it is assumed that elastic loading in plane strain can be continued upto a point from which the transition to plastic plane strain loading is possible merely by adding a residual value of σ_{33} . Referring to Drucker and Prager (1952), this is given by

$$\begin{aligned}
 k &= f = g(\sigma_{11}, \sigma_{22}, \tau_{12}) \\
 &= \frac{3}{2} \alpha (\sigma_{11} + \sigma_{22}) + \sqrt{1 - 3\alpha^2} \left[\frac{\sigma_{11} - \sigma_{22}}{2} + \tau_{12} \right]^{\frac{1}{2}}
 \end{aligned}
 \tag{I-49}$$

1.3. Stress-Strain Behavior of Jointed/Cracked Rock

Mathematical simulation of behavior of jointed rock must allow for closing of pre-existing open joints under compressive loads followed by linear elastic behavior upto initiation of fracture. After fracture occurs, the material cannot take any tension locally in the direction normal to the plane of crack. Non-monotonic loads may involve closing-cracking-closing cycles.

The finite element method has been applied to jointed rock (Anderson and Dodd, 1966; Goodman et al, 1968; Duncan and Goodman, 1968; Malma, 1971). Anderson and Dodd used pin ended one-dimensional elements across a fault to allow compressive stresses to be transferred in the direction normal to the fault. The fault plane was assumed to have no resistance against shear or tensile loads. This capability is now routinely incorporated in most finite element programs. A

two-dimensional 'soft' material element has long been used to represent weak joint planes in rock. Duncan and Goodman (1968) object to this on the basis of large number of elements needed to ensure a reasonable 'aspect ratio' in shape of elements. This becomes a problem for elements representing very thin joints. A one-dimensional element with shear and normal stiffness characteristics was proposed by Goodman et al (1968) to eliminate this objection. Recently (1971), the same investigators have introduced nonlinear properties in this type of element. This approach is quite effective for the case of pre-existing joints in rock. For well defined orthogonal joint systems, an orthotropic continuum approach was suggested by Duncan and Goodman (1968). Christian is credited (Einstein, Bruhn and Hirschfeld, 1970) with development of an element capable of simulating constant shear and residual shear characteristics.

In all these investigations, a distinct set of elements is used to represent the joint. This is alright for pre-existing joints but is impracticable for discontinuities arising as a result of fracture under applied load. To use the same procedure both for pre-existing joints as well as post failure cracks, it is necessary to allow cracks and joints within elements. Then, the mesh layout is more flexible and arbitrary failure laws can be used. Malina (1971) used this approach to study failure along joint planes and then went on to compute the amount of slip and accompanying stress redistribution on the basis of deformation or slip theory of plasticity.

Apparently, a bimodular analysis procedure (Sandhu and Wilson, 1970) can be used to represent pre-existing joints as well as fractures. Bimodularity would be dependent upon the joint opening. However, noting that fractured or open jointed rock has no resistance to tension in the direction normal to that of fracture, a simple approach following the procedure introduced by Zienkiewicz et al (1968) is more convenient. The 'no tension' method of Zienkiewicz consists of first obtaining a solution assuming the system to be linear elastic. Then the elements in tension are relieved of the tensile stresses by application of self-equilibrating forces in elements and at nodal points. This gives an iterative scheme for redistribution of loads to surrounding rock and a lower bound to the exact solution. This approach is essentially an orthotropic continuum approach with the orthotropy being applied to individual elements depending upon the orientation of the fracture plane. The fracture plane defines also a plane of material orthotropy. The relationship between principal stresses and strains can be written as

$$\begin{pmatrix} \sigma_1 \\ \sigma_2 \end{pmatrix} = \begin{bmatrix} C_{11} & C_{12} \\ C_{12} & C_{22} \end{bmatrix} \begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix} \quad (\text{I-50})$$

or symbolically

$$\sigma_p = C_p \epsilon_p \quad (\text{I-51})$$

The laws of transformation of stress and strain give,

$$\begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix} = \begin{bmatrix} \cos^2 \theta & \sin^2 \theta & \sin \theta \cos \theta \\ \sin^2 \theta & \cos^2 \theta & -\sin \theta \cos \theta \end{bmatrix} \begin{pmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{pmatrix} \quad (\text{I-52})$$

and

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \end{Bmatrix} = \begin{bmatrix} \cos^2 \theta & \sin^2 \theta \\ \sin^2 \theta & \cos^2 \theta \\ \sin \theta \cos \theta & -\sin \theta \cos \theta \end{bmatrix} \begin{Bmatrix} \sigma_1 \\ \sigma_2 \end{Bmatrix} \quad (I-53)$$

where $\sigma_x, \sigma_y, \tau_{xy}$ are components of stress and $\epsilon_x, \epsilon_y, \gamma_{xy}$ are components of strain in x, y coordinate system and θ is the angle between the principal direction 1 and x-axis. Symbolically, the above equations are

$$\sigma = J^T \sigma_p \quad (I-54)$$

$$\epsilon_p = J \epsilon \quad (I-55)$$

Substitution in I-51 gives

$$\sigma = J^T C_p J \epsilon \quad (I-56)$$

$$= C \epsilon \quad (I-57)$$

where $C = J^T C_p J \quad (I-58)$

Equation I-58 gives the transformation for stress-strain relation for principal direction to any arbitrary choice of coordinates. The matrix C is singular only for $\theta = 0$ or 90° . It is thus possible to use the relationship in principal stresses and principal strains as the starting point.

In finite element analysis procedures, the stiffness matrix for the system is the sum of element stiffnesses.

$$K = \sum_{m=1}^M k^m \quad (I-59)$$

where K is the system stiffness, k^m is the stiffness of the m th element, and Σ is viewed as a direct stiffness summation operator. Further, element stiffness is related to constitutive relationship through the equation

$$k^m = \int_{V_m} [b^T c b] dV \quad (I-60)$$

where b is the matrix relating strains to nodal point displacements, and V is the volume of the element. Using Equation I-58, the integrand in I-60 can be written

$$b^T c b = b^T J^T c_p J b \quad (I-61)$$

$$= b^T J^T \sigma_p \quad (I-62)$$

$$= B^T \sigma_p \quad (I-63)$$

where $B^T = b^T J^T$ relates principal stresses to nodal point forces.

Occurrence of fracture in an element reduces its ability to take tensile stresses normal to the fracture plane. Also there can be no shear transmitted across a crack which therefore is a principal direction. Thus, it is reasonable to reanalyze the system assigning an orthotropic constitutive relationship and a prescribed principal direction to the element containing a fracture. The procedure is to be repeated until no further fracturing occurs under a given load. To allow for nonlinearity introduced by progressive cracking, incremental procedures are required.

In Zienkiewicz et al (1968), a change in element stiffness was considered equivalent to a pseudo-load. Thus an iterative solution scheme was set up in

which each iteration only involved a back-substitution operation. The pseudo-loads were computed as equivalent to tensile principal stresses. This is satisfactory when both principal stresses are tensile. However, when only one of the principal stresses is tensile, use of pseudo-load corresponding to one principal stress introduces a non-symmetric constitutive law. Actually if the physical concept of 'unloading without any displacements' be followed, a change in the second principal stress corresponding to Poission's effect due to the first stress must be included. This modification is included in the computer program in Chapter IV.

CHAPTER II
THE FINITE ELEMENT METHOD

25a

Chapter II. The Finite Element Method

2.1. Basic Concepts

A boundary value problem can be stated in the form

$$A u = f \text{ on } F \quad (\text{II-1})$$

where u is the unknown function to be determined, A is an operator, and f is the 'forcing' function. F is the domain of interest and may be an open, connected, bounded spatial region embedded in R^3 or in a cartesian product, $R^3 \times [0, \infty)$ where $[0, \infty)$ is the non-negative time interval. In addition to the field equation II-1, there will be some conditions to be satisfied on boundary S of F . For A linear positive, it can be shown that equation II-1 has a unique solution. Necessarily, any approximate solution will in general not coincide with the unique solution of II-1 and consequently no approximate solution is expected to satisfy the field equation as well as the boundary conditions completely.

Solutions to engineering problems as well as the forcing functions are in general bounded and therefore belong to L_2 , the space of square integrable function. However u may be contained in a subset D of L_2 such that A is defined on D . We assume that D is dense in L_2 . If the set of functions $\{\phi_k, k=1,2,\dots,\infty\}$ is an orthonormal basis in D , then any function u can be expressed as an infinite sum:

$$u = \sum_{k=1}^N a_k \phi_k \quad (\text{II-2})$$

A scheme to generate approximate solutions is to use a finite set of terms in the

infinite sum above. Thus, we use

$$\bar{u} = \sum_{k=1}^N a_k \phi_k \quad (\text{II-3})$$

as an approximation. The approximation process then consists of appropriate choice of N , ϕ_k and the coefficients a_k . Several alternative procedures are available. The finite element method is a special process of selection of a finite subset of the basis $\{\phi_k\}$. The coefficients a_k are evaluated by an extension of Ritz method or other standard procedures.

The finite element method is well documented in literature (Zienkiewicz, 1967; Bell and Holand, WPAFB Conference, 1965, 1968; Felippa, 1966; Clough, 1960, 1965). Its theoretical basis (Oden, 1969; de Arantes e Oliveira, 1968, Zlamal, 1968; Melkes, 1970) and relationship to variational principles (Melosh, 1963, Pian and Tong, 1969) have been examined. Essentially, a finite element idealization partitions the spatial region F into a finite number of nontrivial discrete elements or subregions. The geometry of the elements is defined by a set of points in space called the nodal points of the system.

Over an element e let an approximation to u be

$$u^e = \sum_{k=1}^{N^e} a_k^e \bar{\phi}_k^e \quad (\text{II-3})$$

or in matrix form

$$u^e = [\bar{\phi}^e]^T \{a^e\} \quad (\text{II-4})$$

where $\{\bar{\phi}^e\}^T$ is a row vector consisting of $\bar{\phi}_k^e$ as its elements and $\{a^e\}$ is a

column vector of coefficients a_k^e . Evaluating the function at nodal points

$$\{u_1^e\} = \{\bar{\phi}_1^e\}^T \{a^e\} \quad (\text{II-5})$$

where $\{u_1^e\}$ is the vector of nodal point values of the function and $\{\bar{\phi}_1^e\}^T$ is the matrix of base functions evaluated at each nodal point. Rows and columns of $[\bar{\phi}_1^e]^T$ are linearly independent. If square, the matrix is invertible. If the number of nodal points is not equal to the number of independent base functions, a least square procedure can be used for inversion. Hence, we can write

$$\begin{aligned} \{a^e\} &= \left[[\bar{\phi}^e]^T \right]^{-1} \{u_1^e\} \\ &= [A]^{-1} \{u_1^e\} \end{aligned} \quad (\text{II-6})$$

where $A = [\bar{\phi}_1^e]^T$

Substituting II-6 in II-5

$$u^e = \{\bar{\phi}^e\}^T [A]^{-1} \{u_1^e\} \quad (\text{II-7})$$

$$= \{\phi^e\}^T \{u_1^e\} \quad (\text{II-8})$$

where $\{\phi^e\}$ can now be regarded as a set of interpolating functions relating nodal point values of a function to the value of an arbitrary point within the element.

2.2. A Potential Energy Formulation

We assume the rock continuum or 'discontinuum' to be stepwise linear for sufficiently small steps in loading. For such a case the governing equations are

$$\sigma_{kl,k} + \pi_{,k} \delta_{lk} + \rho F_l = 0 \quad (\text{II-9})$$

$$\sigma_{ij} = E_{ijkl} \epsilon_{kl} + \bar{\sigma}_{ij} \quad (\text{II-10})$$

$$\epsilon_{kl} = u_{(k,l)} \quad (\text{II-11})$$

where σ_{ij} , ϵ_{ij} , E_{ijkl} , F_i , u_i are components, respectively, of the symmetric stress tensor, the symmetric strain tensor, the isothermal elasticity tensor, the body forces vector per unit mass, and the displacement vector. ρ is the mass density and δ_{ik} is the kronecker delta. $\bar{\sigma}_{ij}$ are components of initial stress corresponding to zero displacement and π is the pore pressure. Potential energy formulation uses the functional

$$\begin{aligned} \Omega = & \int_F \left[\epsilon_{ij} E_{ijkl} \epsilon_{kl} - u_i \sigma_{ij,j} - 2 \epsilon_{ij} \sigma_{ij} + u_i \sigma_{ij} \right. \\ & \left. - 2 u_i \rho F_i - 2 u_i \pi_{,i} + 2 \epsilon_{ij} \bar{\sigma}_{ij} \right] dF \\ & + \int_{s_1} u_i (\sigma_{ij} n_j - \hat{t}_i) ds - \int_{s_2} (u_i - \hat{u}_i) \sigma_{ij} n_j ds \quad (\text{II-12}) \end{aligned}$$

where we have included the boundary condition

$$\sigma_{ij} n_j = \hat{t}_i \quad \text{on } s_1 \quad (\text{II-13})$$

$$u_i = \hat{u}_i \quad \text{on } s_2 \quad (\text{II-14})$$

s_1 , s_2 are complementary subsets of S the boundary of F and n_j are components of unit vector normal to surface.

Symmetry of the field equations leads to the functional

$$\begin{aligned} \Omega = & \int_F \left[\epsilon_{ij} E_{ijkl} \epsilon_{kl} - 2 \epsilon_{ij} \sigma_{ij} + 2 u_{i,j} \sigma_{ij} - 2 u_i \rho F_i \right. \\ & \left. - 2 u_i \pi_{,i} + 2 \epsilon_{ij} \bar{\sigma}_{ij} \right] dF \\ & - 2 \int_{s_1} u_i \hat{t}_i ds - 2 \int_{s_2} (u_i - \hat{u}_i) \sigma_{ij} n_j ds \quad (\text{II-15}) \end{aligned}$$

Further assuming that we restrict our choice of ϵ_{ij} , u_i such that II-11 and II-14

are identically satisfied, the functional in II-15 reduces to

$$\Omega = \int_F \left[u_{i,j} E_{ijkl} u_{k,l} - 2 u_i \rho F_i - 2 u_i \pi_{,i} + 2 u_{i,j} \bar{\sigma}_{ij} \right] dF - 2 \int_{s_1} u_i \hat{t}_i ds \quad (\text{II-16})$$

Replacing \int_F by $\sum_{m=1}^m \int_{F^m}$ where F^m represents the subregion or element m , and using suitable interpolation scheme to express the integrand in terms of nodal point vectors of displacement, vanishing of variation of the functional yields the matrix equation

$$[K] \{u\} = \{R\} \quad (\text{II-17})$$

where

$$[K] = \sum_{m=1}^m [k^m] \quad (\text{II-18})$$

$$\{R\} = \sum_{m=1}^m \left[\{L^m\} + \{P^m\} - \{Q^m\} + \{T^m\} \right] \quad (\text{II-19})$$

Components of element stiffness matrix and load vectors are:

$$k_{ij}^m = \int_{F^m} \phi_{im,n}^m E_{mnpq} \phi_{jp,q}^m \quad (\text{II-20})$$

$$L_i^m = \int_{F^m} \rho \phi_{ij}^m F_j \quad (\text{II-21})$$

$$P_i^m = \int_{F^m} \phi_{ij}^m \pi_{,j} \quad (\text{II-22})$$

$$Q_i^m = \int_{F^m} \phi_{im,n}^m \bar{\sigma}_{mn} \quad (\text{II-23})$$

$$T_i^m = \int_{F^m} \phi_{ij}^m \hat{t}_j \quad (\text{II-24})$$

and ϕ_{ij} are components of a matrix formed by the row vectors $\{\phi^e\}^T$ corresponding to each degree of displacement freedom. The vectors $\{L^m\}$, $\{P^m\}$, $\{Q^m\}$, $\{T^m\}$ represent the contribution to the load vector made respectively by the body forces, the pore pressure gradients, the initial stresses and the boundary loads in the element m .

2.3 Incremental Analysis

In case of incremental construction and incremental application of loads, the loads, stresses and displacements for any incremental step can be written as $\{R_n\}$, $\{\sigma_n\}$ and $\{u_n\}$. Then for the next stage, $\{\sigma_n\}$ and $\{u_n\}$ can be regarded as the initial stresses and the initial displacements for the structural system. Thus the matrix equations are

$$\left[K_{n+1} \right] \left\{ u_{n+1} - u_n \right\} = \left[K_{n+1} \right] \left\{ \Delta u_n \right\} = \left\{ \Delta R_n \right\} \quad (\text{II-25})$$

where

$$\Delta R_n = \sum_{m=1}^m \left[\left\{ \Delta L^m \right\} + \left\{ \Delta P^m \right\} - \left\{ \Delta Q^m \right\} + \left\{ \Delta T^m \right\} \right] \quad (\text{II-26})$$

and $\{\Delta L^m\}$, $\{\Delta P^m\}$, $\{\Delta Q^m\}$, $\{\Delta T^m\}$ are increments in the respective quantities.

For elastic-plastic analysis, the stiffness depends upon stress and has to be re-evaluated at small increments of load. To ensure manageable computation, the increments are kept at the largest practicable without loss of accuracy.

CHAPTER III

COMPUTER PROGRAM FOR PLANE STRAIN ANALYSIS OF
ELASTIC-PLASTIC MOHR-COULOMB MATERIALS

31a

Chapter III. Computer Program for Plane Strain Analysis of Elastic-Plastic Mohr-Coulomb Materials

3.1. Organization

Computer program described here is based on the theory presented earlier in this report. The program is written in Fortran IV language.

The program is intended to calculate stresses and strains for a plane strain problem in rock mechanics. Mohr-Coulomb yield criterion has been used. It makes allowance for the boundary conditions, residual stresses, stresses due to temperature change, and varying pressure boundaries. The structure may consist of different materials. It uses Wilson's (1965) quadrilateral elements and generates stiffness in line with the integration procedures discussed by Felippa (1966).

The principal program called MAIN controls all the data input and control information. It does the basic system initialization and prints the control data and material and geometrical properties of the structure. Stiffness formulation, equation solving and stress calculations are done by the subroutines called by MAIN.

3.12. Stiffness and Load Matrices

Stiffness matrix for each analysis is computed in blocks by the subroutine STIFF. For the element stiffness it calls QUAD for triangular and quadrilateral elements which have been allowed by this program. The element stiffness is added to the total stiffness using the direct stiffness technique. Concentrated loads are included in the load matrix. Equations are modified for the displacement boundary conditions by calling subroutine MODIFY. For the stress-strain matrix QUAD

calls STRSTR. With the constitutive law being available, stiffness of the system and load matrix are computed in the STIFF subroutine.

3.13. Calculations of Displacements

After the stiffness and load matrices for a stage have been computed, the resulting equations are solved by calling subroutine BANSOL. This uses Gaussian elimination technique for banded equations by Wilson (1963). In this the triangularization of stiffness matrix is done. Back substitution through the triangulized matrix gives the solution.

3.14. Calculations of Stresses

In the first cycle a purely elastic solution is obtained for the problem by solving

$$[K^e] \{r\} = \{R\} \quad (\text{III-1})$$

where $[K^e]$ is the elastic stiffness
 $\{r\}$ is the displacement vector
 $\{R\}$ is the load vector.

This can be done easily by assuming all the elements to be elastic to begin with. As the problem is a nonlinear one, therefore this solution will not be correct. In our analysis the system is assumed to be stepwise linear between the yielding of one element to the other. This is assumed not to cause any significant error. In Fig. 3-1 point A represents the initial stresses and C the final stresses in an element. The curve $f = k$ represents the yield surface. For those elements which become plastic under this loading \overrightarrow{AC} meets the surface $f = k$ at pt. B. It is seen easily that for the element it is not possible to be loaded to point C but it can be loaded to point B only, assuming proportional loading.

Let $S_r = \frac{|\vec{AB}|}{|\vec{AC}|}$. S_r is called the stress ratio. To calculate S_r we

know that

$$\begin{aligned}\vec{AC} &= (\sigma_{ij})_f - (\sigma_{ij})_i \\ \vec{OB} &= \vec{OA} + \vec{AB} \\ &= \vec{OA} + S_r \cdot \vec{AC} \\ &= (\sigma_{ij})_i + S_r \cdot (\sigma_{ij})_f - (\sigma_{ij})_i\end{aligned}$$

$$\sigma_{ij} = (\sigma_{ij})_i (1 - S_r) + S_r (\sigma_{ij})_f$$

As the point B lies on the yield surface $f = k$

$$f \left[(\sigma_{ij})_i (1 - S_r) + S_r (\sigma_{ij})_f \right] = k \quad (\text{III-2})$$

From equation (III-2) the value of S_r can be calculated. This stress ratio represents the fraction by which the increment in stress is to be scaled to bring the final load on the yield surface.

Value of S_r is calculated for all the elements. The element in which the final stress state is farthest from the yield surface will have the minimum stress ratio. If we scale down the displacements and stresses in this ratio, we shall have the stresses and strains precisely at the point when the system has its first element just going into the plastic region from purely elastic system.

In the next step the element having stress ratio equal to the minimum value is assumed to be plastic. To economize on computer time all such elements which have their value of stress ratio in the vicinity of minimum were also allowed to go plastic. As the stress-strain matrix is known the stiffness is calculated again and

equation
$$[K_1] \{r_1\} = \{R_1\} \quad (\text{III-3})$$

is solved. This procedure is repeated until the whole load has been applied to the system and cumulative stresses and displacements calculated. The stresses in (i-1)th step become the initial stresses for ith step.

3.2. Input Data Preparation

1. **Control Card (A6).** This card will carry the characters **START** in columns 1-5. This will start the processing of the data deck which consists of the following set of cards.
2. **Job Title (72H).** This card will give the descriptive identification for the job.
3. **Control Information (4I5, 2F10.2, I5)**

<u>Information</u>	<u>Columns</u>
Total number of nodal points	1 - 5
Total number of elements	6 - 10
Number of different materials	11 - 15
Number of pressure boundary cards	16 - 20
Body Force in x-direction	21 - 30
Body Force in y-direction	31 - 40
Number of Approximations	41 - 45

4. **Material Property Cards.** One set of 2 cards is provided for each material.

In each set:

- a. **first card (1I5, F10.0) will give the following information**

Material identification number	1 - 5
Mass density of the material	6 - 15

- b. **The second card will carry the following information (4F10.0)**

<u>Information</u>	<u>Columns</u>
Elastic Modulus	1 - 10
Poisson's Ratio	11 - 20
Cohesion	21 - 30
Friction Angle in Degrees	31 - 40

5. Nodal Point Cards (I5, F5.0, 5F 10.0).

One card for each nodal point with the following information:

Nodal Point number	1 - 5
Type of Nodal point	6 - 10
X-ordinate	11 - 20
Y-ordinate	21 - 30
XR	31 - 40
XZ	41 - 50

If the number in columns 6 - 10 is

Zero	XR is the specified X-load and XZ is the specified Y-load
1	XR is the specified X-displacement and XZ is the specified Y-load
2	XR is the specified X-load and XZ is the specified Y-displacement
3	XR is the specified X-displacement and XZ is the specified Y-displacement

All loads are considered to be total forces acting on an element of unit thickness.

Nodal point cards must be in numerical sequence. If cards are omitted, the omitted nodal points are generated at equal intervals along a straight line between the defined nodal points. The type of the nodal point, as well as XR, XZ, are set equal to zero.

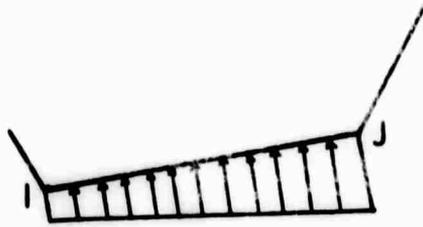
6. **Element Cards (6I5).** One card for each element will provide the following data.

<u>Information</u>	<u>Columns</u>
Number of element	1 - 5
Nodal point I	6 - 10
Nodal point J	11 - 15
Nodal point K	16 - 20
Nodal point L	21 - 25
Material type	26 - 30

Nodal points I, J, K, L are corners of each individual element in a counter-clockwise order for a right handed system of coordinates. For triangular elements set nodal point L same as nodal point K. The element cards must be in the numerical sequence. Any cards that are omitted will be automatically generated in the program by incrementing each of the I, J, K, and L nodal points by one. The material type will be taken the same as for the last element defined.

7. **Pressure Boundary Cards (2I5, 2F10.0).** One card for each boundary element which is subjected to a normal pressure will carry the following information:

<u>Information</u>	<u>Columns</u>
Nodal Point I	1 - 5
Nodal Point J	6 - 10
Normal Pressure at I	11 - 20
Normal Pressure at J	21 - 30



As shown in the sketch, the boundary element must be on the left as one progresses from I to J. Surface tensile forces is input as a negative pressure.

Output Information:

The following information is developed and printed by the program:

1. Reprint of input data
2. Nodal point displacements
3. Stresses at the center of each element

3.3 Program Listing

```
C      PLANE STRAIN ANALYSIS OF ELASTIC-PLASTIC MOHR-COULOMB MATERIALS
C
      IMPLICIT REAL*8(A-H,O-Z)
      COMMON ACELR,ACELZ,VOL,TEMP,SIG1(400,7),HED(18),E(4,12),RO(12),
      *R(500),Z(500),UR(500),UZ(500),CODE(500),T(500),PR(100,2),
      *NUMMAT,NUMPC,N,MTYPE,KKK,NUMNP,NUMEL,NNN,
      *IBC(100),JBC(100),MTAG(400)
      COMMON/ARG/ RR(5),ZZ(5),S(10,10),P(8),ST(3,10),C(3,3),SIG(7),EE(4)
      *,BO(1000),SR1,SR2,
      *RATIO(400),LM(4),IX(400,5),XC,YC
      COMMON /BANARG/ MBAND,NUMBLK,B(108),A(108,54)
      DIMENSION WORD(2)

C
      DATA WORD/ 6HSTART ,6HSTOP /
      CALL ERRSET(207,256,-1,1)
      CALL ERRSET(208,256,-1,1)

C
C
      5  READ (5,1006) WORD1
         IF (WORD1.EQ.WORD(1)) GO TO 50
         IF (WORD1.EQ.WORD(2)) STOP
         GO TO 5
      50 READ (5,1000) HED,NUMNP,NUMEL,NUMMAT,NUMPC,ACELR,ACELZ,NP
         WRITE(6,2000) HED,NUMNP,NUMEL,NUMMAT,NUMPC,ACELR,ACELZ,NP
         DO 55 M=1,NUMMAT
            READ (5,1001) MTYPE,RO(MTYPE)
            WRITE(6,2001) MTYPE,RO(MTYPE)
            READ (5,1002) (E(J,MTYPE),J=1,4)
            WRITE(6,2002) (E(J,MTYPE),J=1,4)
      55 CONTINUE

C
         WRITE (6,2003)
         L=0

C
      60 READ (5,1003) N,CODE(N),R(N),Z(N),UR(N),UZ(N),T(N)
         NL=L+1
         ZX=N-L
         DR=(R(N)-R(L))/ZX
         DZ=(Z(N)-Z(L))/ZX
         DT=(T(N)-T(L))/ZX
      70 L=L+1
         IF(N-L) 100,90,80
      80 CODE(L)=0.0
         R(L)=R(L-1)+DR
         Z(L)=Z(L-1)+DZ
         T(L)=T(L-1)+DT
         UR(L)=0.0
         UZ(L)=0.0
```

```

GO TO 70
90 WRITE (6,2004) (K, CODE(K), R(K), Z(K), UR(K), UZ(K), T(K), K=NL, N)
IF (NUMNP-N) 100, 110, 60
100 WRITE (6,2005) N
CALL EXIT
110 CONTINUE

```

C

```

WRITE (6,2006)
N=0
130 READ (5,1004) M, (IX(M,I), I=1,5), (SIGI(M,I), I=1,4)
ZX=M-N
DO 135 I=1,4
135 SIG(I)=(SIGI(M,I)-SIGI(N,I))/ZX
140 N=N+1
IF (M-N) 170, 170, 150
150 IX(N,1)=IX(N-1,1)+1
IX(N,2)=IX(N-1,2)+1
IX(N,3)=IX(N-1,3)+1
IX(N,4)=IX(N-1,4)+1
IX(N,5)=IX(N-1,5)
DO 160 I=1,4
160 SIGI(N,I)=SIGI(N-1,I)+SIG(I)
170 WRITE(6,2007) N, (IX(N,I), I=1,5), (SIGI(N,I), I=1,4)
IF (M-N) 180, 180, 140
180 IF (NUMEL-N) 190, 190, 130
190 CONTINUE

```

C

```

IF (NUMPC) 290, 310, 290
290 WRITE (6,2008)
DO 300 L=1, NUMPC
READ (5,1005) IBC(L), JBC(L), PR(L,1), PR(L,2)
300 WRITE (6,2009) IBC(L), JBC(L), PR(L,1), PR(L,2)
310 CONTINUE

```

C

```

J=0
DO 340 N=1, NUMEL
MTAG(N)=0
SIGI(N,5)=0.
SIGI(N,6)=0.
SIGI(N,7)=0.
DO 340 I=1,4
DO 325 L=1,4
KK=IABS(IX(N,I)-IX(N,L))
IF (KK-J) 325, 325, 320
320 J=KK
325 CONTINUE
340 CONTINUE
MBAND=2*J+2

```

```

WRITE(6,1007) MBAND
DO 350 N=1,NUMNP
BO(2*N-1)=0.
350 BO(2*N)=0.
C
SR1=1.0
SR2=0.0
DO 500 NNN=1,NP
KKK=0
CALL STIFF
CALL BANSOL
CALL STRESS
DO 400 N=1,NUMNP
NN=2*N
BO(NN-1)=BO(NN-1)+B(NN-1)
BO(NN)=BO(NN)+B(NN)
400 CONTINUE
WRITE(6,2010) (N,BO(2*N-1),BO(2*N),N=1,NUMNP)
IF(KKK.EQ.0) CALL EXIT
500 CONTINUE
GO TO 5
C
1000 FORMAT (18A4/4I5,2F10.2,2I5)
1001 FORMAT (1I5,1F10.0)
1002 FORMAT (6F10.0)
1003 FORMAT (I5,F5.0,5F10.0)
1004 FORMAT(6I5,4F10.0)
1005 FORMAT (2I5,2F10.0)
1006 FORMAT(A6)
1007 FORMAT( ' BAND WIDTH FOR THIS DATA = ' , I5 )
2000 FORMAT (1H1 18A4/
1 30H0 NUMBER OF NODAL POINTS----- I3 /
2 30H0 NUMBER OF ELEMENTS----- I3 /
3 30H0 NUMBER OF DIFF. MATERIALS--- I3 /
4 30H0 NUMBER OF PRESSURE CARDS---- I3 /
5 30H0 X-ACCELERATION----- E12.4/
6 30H0 Y-ACCELERATION----- E12.4/
7 30H0 NUMBER OF APPROXIMATIONS---- I12)
2001 FORMAT (17H0MATERIAL NUMBER# I3, 15H, MASS DENSITY# E12.4)
2002 FORMAT(16H0ELASTIC MODULUS 14X 2HNU BX 8HCOHESION 2X 14HFRICITION A
*NGLE/(2E16.5,2F16.5))
2003 FORMAT (111H1NODAL POINT TYPE X ORDINATE Y ORDINATE X LO
1AD OR DISPLACEMENT Y LOAD OR DISPLACEMENT PORE PRESSURE )
2004 FORMAT (I12,F12.2,2F12.3,2E24.7,F15.3)
2005 FORMAT (26H0NODAL POINT CARD ERROR N# I5)
2006 FORMAT(96H1ELEMENT NO. I J K L MATERIAL X-ST
*RESS Y-STRESS XY-STRESS Z-STRESS)
2007 FORMAT(I12,4I6,1I12,4F12.3)

```

2008 FORMAT (29HOPRESSURE BOUNDARY CONDITIONS/ 40H I J PRESS
*URE I PRESSURE J)
2009 FORMAT (2I6,2F12.3)
2010 FORMAT (12HIN.P. NUMBER 18X 2HUX 18X 2HUY / (1I12,2E20.7))
END

SUBROUTINE STIFF

```

C
  IMPLICIT REAL*8(A-H,O-Z)
  COMMON ACELR,ACELZ,VOL,TEMP, SIGI(400,7),HED(18),E(4,12),RO(12),
  *R(500),Z(500),UR(500),UZ(500),CODE(500),T(500),PR(100,2),
  *NUMMAT,NUMPC,N,MTYPE,KKK,NUMNP,NUMEL,NNN,
  *IBC(100),JBC(100),MTAG(400)
  COMMON/ARG/ RR(5),ZZ(5),S(10,10),P(8),ST(3,10),C(3,3),SIG(7),EE(4)
  *,BO(1000),SR1,SR2,
  *RATIO(400),LM(4),IX(400,5),XC,YC
  COMMON /BANARG/ MBAND,NUMBLK,B(108),A(108,54)

C
C
  REWIND 2
  NR=27
  ND=2*NB
  ND2=2*ND
  STOP=0.0
  NUMBLK=0
  DO 50 N=1,ND2
  B(N)=0.0
  DO 50 M=1,ND
50 A(N,M)=0.0

C
60 NUMBLK=NUMBLK+1
  NH=NB*(NUMBLK+1)
  NM=NH-NB
  NL=NM-NR+1
  KSHIFT=2*NL-2
  DO 210 N=1,NUMEL
  IF (IX(N,5)) 210,210,65
65 DO 80 I=1,4
  IF (IX(N,I)-NL) 80,70,70
70 IF (IX(N,I)-NM) 90,90,80
80 CONTINUE
  GO TO 210
90 CALL QUAD
  IX(N,5)=-IX(N,5)
  IF(VOL) 100,100,110
100 WRITE(6,2000) N
  STOP=1.0

C
110 MM=4
  IF(IX(N,3)-IX(N,4)) 130,120,130
120 MM=3
130 DO 140 I=1,MM
140 LM(I)=2*IX(N,I)-2
  DO 200 I=1,MM

```

```

DO 200 K=1,2
II=LM(I)+K-KSHIFT
KK=2*I-2+K
A(II)=B(II)+P(KK)
DO 200 J=1,MM
DO 200 L=1,2
JJ=LM(J)+L-1+1-KSHIFT
LL=2*J-2+L
IF(JJ) 200,200,175
175 IF(MD-JJ) 180,195,195
180 WRITE (6,2001) N
STOP=1.0
GO TO 210
195 A(II,JJ)=A(II,JJ)+S(KK,LL)
200 CONTINUE
210 CONTINUE

```

C

```

DO 220 N=NL,NM
K=2*N-KSHIFT
A(K)=B(K)+UZ(N)
B(K-1)=B(K-1)+UR(N)
T(N)=0.
UZ(N)=0.
220 UR(N)=0.

```

C

```

IF (NUMPC) 225,310,225
225 DO 300 L=1,NUMPC
I=IAC(L)
J=JAC(L)
DR=Z(I)-Z(J)
DZ=R(J)-R(I)
PP2=(PR(L,2)+PR(L,1))/6.
PP1=PP2+PR(L,1)/6.
PP2=PP2+PR(L,2)/6.
II=2*I-KSHIFT
JJ=2*J-KSHIFT
IF(II) 265,265,235
235 IF(II-MD) 240,240,265
240 B(II-1)=B(II-1)+PP1*DR
A(II)=A(II)+PP1*DZ
265 IF(JJ) 300,300,270
270 IF(JJ-MD) 275,275,300
275 B(JJ-1)=B(JJ-1)+PP2*DR
A(JJ)=A(JJ)+PP2*DZ
300 CONTINUE

```

C

```

310 DO 400 M=NL,NM
IF (M-NUMNP) 315,315,400

```

```

315 U=UR(M)
    N=2*N-1-KSHIFT
    IF (CODE(M)) 390,400,316
316 IF (CODE(M)-1.) 317,370,317
317 IF (CODE(M)-2.) 318,390,318
318 IF (CODE(M)-3.) 390,380,390
370 CALL MODIFY(A,B,ND2,MBAND,N,U)
    GO TO 400
380 CALL MODIFY(A,B,ND2,MBAND,N,U)
390 U=UZ(M)
    N=N+1
    CALL MODIFY(A,B,ND2,MBAND,N,U)
400 CONTINUE
C
    WRITE (2) (B(N),(A(N,M),M=1,MBAND),N=1,ND)
C
    DO 420 N=1,ND
    K=N+ND
    B(N)=B(K)
    B(K)=0.0
    DO 420 M=1,ND
    A(N,M)=A(K,M)
420 A(K,M)=0.0
C
    IF (NM-NUMNP) 60,480,480
480 CONTINUE
    ACELR=0.
    ACELZ=0.
    NUMPC=0
C
    IF (STOP) 490,500,490
490 CALL EXIT
500 RETURN
C
2000 FORMAT (26HNEGATIVE AREA ELEMENT NO. 14)
2001 FORMAT (29HOBAND WIDTH EXCEEDS ALLOWABLE 14)
    END

```

SUBROUTINE QUAD

C

```

IMPLICIT REAL*8(A-H,O-Z)
COMMON ACELR,ACELZ,VOL,TEMP, SIGI(400,7),MED(10),E(4,12),RO(12),
*R(500),Z(500),UR(500),UZ(500),CODE(500),T(500),PR(100,2),
*NUMMAT,NUMPC,N,MTYPE,KKK,NUMNP,NUMEL,NNN,
*IBC(100),JBC(100),MTAG(400)
COMMON/ARG/ RR(5),ZZ(5),S(10,10),P(8),ST(3,10),C(3,3),SIG(7),EE(4)
*.BN(1000),SR1,SR2,
*RATIO(400),LM(4),IX(400,5),XC,YC
COMMON /BANARG/ MBAND,NUMBLK,B(108),A(108,54)
DIMENSION U(3),V(3)

```

C

```

CALL STRSTR
DO 130 J=1,10
DO 120 I=1,3
120 ST(I,J)=0.
DO 130 I=1,10
130 S(I,J)=0.
DO 140 I=1,4
NPP=IX(N,I)
RR(I)=R(NPP)
140 ZZ(I)=Z(NPP)
XC=(RR(1)+RR(2)+RR(3)+RR(4))/4.
YC=(ZZ(1)+ZZ(2)+ZZ(3)+ZZ(4))/4.
RR(5)=XC
ZZ(5)=YC
K=5
J=1
I=4
LM(3)=9
NT=4
IF(IX(N,3)-IX(N,4)) 160,150,160
150 NT=1
LM(3)=5
I=1
K=3
J=2
XC=(RR(1)+RR(2)+RR(3))/3.
YC=(ZZ(1)+ZZ(2)+ZZ(3))/3.
RR(5)=RR(3)
ZZ(5)=ZZ(3)
160 DO 200 NN=1,NT
LM(1)=2*I-1
LM(2)=2*J-1
U(1)=ZZ(J)-ZZ(K)
U(2)=ZZ(K)-ZZ(I)
U(3)=ZZ(I)-ZZ(J)

```

```

V(1)=RR(K)-RR(J)
V(2)=RR(I)-RR(K)
V(3)=RR(J)-RR(I)
AREA=(RR(J)*U(2)+RR(I)*U(1)+RR(S)*U(3))/2.
VOL=VOL+AREA
COMM=.25/AREA
XNT=NT
COM=2.0/XNT
COM=COM*COMM
C
DO 180 L=1,3
II=LM(L)
ST(1,II)=ST(1,II)+U(L)*COM
ST(2,II+1)=ST(2,II+1)+V(L)*COM
ST(3,II)=ST(3,II)+V(L)*COM
ST(3,II+1)=ST(3,II+1)+U(L)*COM
DO 180 M=1,3
JJ=LM(M)
S(II,JJ)=S(II,JJ)+(U(L)*C(1,1)*U(M)+V(L)*C(3,3)*V(M)+V(L)*C(1,3)*U
I(M)+U(L)*C(1,3)*V(M))*COMM
S(II,JJ+1)=S(II,JJ+1)+(U(L)*C(1,2)*V(M)+V(L)*C(3,3)*U(M)+V(L)*C(2,
13)*V(M)+U(L)*C(1,3)*U(M))*COMM
S(II+1,JJ+1)=S(II+1,JJ+1)+(V(L)*C(2,2)*V(M)+U(L)*C(3,3)*U(M)+U(L)*
1C(2,3)*V(M)+V(L)*C(2,3)*U(M))*COMM
S(JJ+1,II)=S(II,JJ+1)
180 CONTINUE
I=J
J=J+1
200 CONTINUE
C
IF(IX(N,3)-IX(N,4)) 220,250,220
220 DO 240 I=1,2
KK=10-I
DO 240 K=1,KK
CC=S(KK+1,K)/S(KK+1,KK+1)
DO 230 J=1,3
230 ST(J,K)=ST(J,K)-CC*ST(J,KK+1)
DO 240 J=1,KK
240 S(J,K)=S(J,K)-CC*S(J,KK+1)
250 CONTINUE
C
II=C
IF(NNN.EQ.1) GO TO 260
II=4
260 SIG(1)=-SIGI(N,II+1)*TEMP
SIG(2)=-SIGI(N,II+2)*TEMP
SIG(3)=-SIGI(N,II+3)
DO 520 I=1,8

```

```
P(1)=0.0  
DO 510 J=1,3  
510 P(1)=P(1)+ST(J,1)*SIG(J)  
520 P(1)=P(1)*VOL
```

```
C  
DX=VOL*ACELR*RO(MTYPE)/4.  
DY=VOL*ACELZ*RO(MTYPE)/4.  
DO 530 I=1,4  
P(2*I)=P(2*I)+DY  
530 P(2*I-1)=P(2*I-1)+DX  
RETURN  
END
```

SUBROUTINE STRSTR

C
 IMPLICIT REAL*8(A-H,O-Z)
 COMMON ACELR,ACELZ,VOL,TEMP,SIG1(400,7),MED(18),E(4,12),RO(12),
 *R(500),Z(500),UR(500),UZ(500),CODE(500),T(500),PR(100,2),
 *NUMMAT,NUMPC,N,MTYPE,KKK,NUMNP,NUMEL,NNN,
 *IBC(100),JBC(100),MTAG(400)
 COMMON/ARG/ RR(5),ZZ(5),S(10,10),P(8),ST(3,10),C(3,3),SIG(7),EE(4)
 *,BO(1000),SR1,SR2,
 *RATIO(400),LM(4),IX(400,5),XC,YC
 COMMON /BANARG/ MBAND,NUMBLK,B(108),A(108,54)

C
 I=IX(N,1)
 J=IX(N,2)
 K=IX(N,3)
 L=IX(N,4)
 MTYPE=IX(N,5)
 VOL=0.

C
 TFMP=(T(I)+T(J)+T(K)+T(L))/4.0
 DO 50 KK=1,4
 50 EE(KK)=E(KK,MTYPE)
 IF(MTAG(N)) 60,60,70
 60 CC=EE(2)/1.-EE(2)
 RB=EE(1)/1.-EE(2)**2
 COMM=RB/1.-CC**2
 C(1,1)=COMM
 C(1,2)=COMM*CC
 C(1,3)=0.
 C(2,1)=C(1,2)
 C(2,2)=C(1,1)
 C(2,3)=0.
 C(3,1)=0.
 C(3,2)=0.
 C(3,3)=.5*COMM*(1.-CC)
 CC=DTAN(EE(4)/57.296)
 RB=DSQRT(9.0+12.0*CC*CC)
 EE(4)=CC/RB
 EE(3)=3.*EE(3)/RB
 GO TO 50
 70 CC=DTAN(EE(4)/57.296)
 RB=DSQRT(9.0+12.0*CC*CC)
 EE(4)=CC/RB
 EE(3)=3.*EE(3)/RB
 CC=2.*(1.+EE(2))/13.-6.*EE(2)
 DD=(SIG1(N,1)-SIG1(N,2))/2.
 BJ2=(DD*DD+SIG1(N,3)**2)/(1.-3.*(EE(4)**2))
 BJ2=DSQRT(BJ2)

```

BJ1=1.5*(SIGI(N,1)+SIGI(N,2))-3.*EE(4)*BJ2
DD=BJ1/BJ2
BB=1.+9.*(EE(4)**2)*CC
CC=3.*EE(4)*CC-DD/3.
DD=EE(4)-DD/6.
H1=.5*CC/(BB*BJ2)
H2=DD*CC/BB-EE(2)*EE(3)/(BB*BJ2*(1.-2.*EE(2)))
H3=.5/(BB*BJ2*BJ2)

```

C

```

BB=EE(1)/(1.+EE(2))
C(1,1)=BB*(1.-H2-2.*H1*SIGI(N,1)-H3*(SIGI(N,1)**2))
C(1,2)=-BB*(H2+H1*(SIGI(N,1)+SIGI(N,2))+H3*SIGI(N,1)*SIGI(N,2))
C(1,3)=-BB*(H1*SIGI(N,3)+H3*SIGI(N,1)*SIGI(N,3))
C(2,1)=C(1,2)
C(2,2)=BB*(1.-H2-2.*H1*SIGI(N,2)-H3*(SIGI(N,2)**2))
C(2,3)=-BB*(H1*SIGI(N,3)+H3*SIGI(N,2)*SIGI(N,3))
C(3,1)=C(1,3)
C(3,2)=C(2,3)
C(3,3)=BB*(.5-H3*(SIGI(N,3)**2))
500 RETURN
END

```

```
SUBROUTINE MODIFY(A,B,NEQ,MBAND,N,U)
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION A(108,54),B(108)
  DO 250 M=2,MBAND
    K=N-M+1
    IF(K) 235,235,230
230  H(K)=B(K)-A(K,M)*U
    A(K,M)=0.0
235  K=N+M-1
    IF(NEQ-K) 250,240,240
240  B(K)=B(K)-A(N,M)*U
    A(N,M)=0.0
250  CONTINUE
    A(N,1)=1.0
    B(N)=U
  RETURN
  END
```

```
SUBROUTINE BANSOL
IMPLICIT REAL*8(A-H,O-Z)
COMMON /BANARG/ MM,NUMBLK,R(108),A(108,54)
```

C

```
NN=54
NL=NN+1
NH=NN+NN
REWIND 1
REWIND 2
NB=0
GO TO 150
```

C

```
100 NR=NB+1
DO 125 N=1,NN
NM=NN+N
B(N)=B(NM)
R(NM)=0.0
DO 125 M=1,MM
A(N,M)=A(NM,M)
125 A(NM,M)=0.0
```

C

```
IF (NUMBLK-NR) 150,200,150
150 READ (2) (B(N),(A(N,M),M=1,MM),N=NL,NH)
IF (NB) 200,100,200
```

C

```
200 DO 300 N=1,NN
IF (A(N,1)) 225,300,225
225 B(N)=B(N)/A(N,1)
DO 275 L=2,MM
IF (A(N,L)) 230,275,230
230 C=A(N,L)/A(N,1)
I=N+L-1
J=0
DO 250 K=L,MM
J=J+1
250 A(I,J)=A(I,J)-C*A(N,K)
B(I)=B(I)-A(N,L)*B(N)
A(N,L)=C
275 CONTINUE
300 CONTINUE
```

C

```
IF (NUMBLK-NB) 375,400,375
375 WRITE (1) (B(N),(A(N,M),M=2,MM),N=1,NN)
GO TO 100
```

C

```
400 DO 450 M=1,NN
N=NN+1-M
DO 425 K=2,MM
```

```
L=N+K-1
425 R(N)=B(N)-A(N,K)*B(L)
NM=N+NN
R(NM)=C(N)
450 A(NM,NB)=R(N)
NB=NH-1
IF (NB) 475,500,475
475 BACKSPACE 1
READ (1) (R(N),(A(N,M),M=2,MM),N=1,NN)
BACKSPACE 1
GO TO 400
```

C

```
500 K=0
UN 600 NB=1,NUMBLK
DN 600 N=1,NN
K=K+1
NM=N+NN
600 R(K)=A(NM,NB)
```

C

```
RETURN
END
```

```

SUBROUTINE STRESS
  IMPLICIT REAL*8(A-H,O-Z)
  COMMON ACFLK,ACELZ,VOL,TEMP, SIG(400,7),MED(18),E(4,12),RO(12),
  *K(500),Z(500),UR(500),U7(500),CODE(500),T(500),PRI(10,2),
  *NUMMAT,NUMPC,N,MTYPE,KKK,NUMNP,NUMEL,NNN,
  *IHC(100),JRC(100),MTAG(400)
  COMMON/ARG/ RR(5),ZZ(5),S(10,10),PIR(1,ST(3,10)),C(3,3),SIG(7),EE(4)
  *,RO(1000),SR1,SR2,
  *RATIO(400),LMI(4),IX(400,5),XC,YC
  COMMON /BANARG/ MBRAND,NUMRLK,B(100),A(100,54)

```

C

```

  TOL=0.
  SR=1.
  NUMR=0
  MPRINT=0
  KK=0
  DO 300 N=1,NUMEL
  RATIO(N)=1.
  IX(N,5)=IABS(IX(N,5))
  MTYPE=IX(N,5)
  CALL QUAD
  MM=4
  IF(IX(N,3)-IX(N,4)) 170,160,170
160 MM=3
170 DO 180 I=1,3
  RR(I)=0.
  DO 180 J=1,MM
  II=2*J
  JJ=2*IX(N,J)
180 RR(II)=RR(II)+ST(II,II)*B(JJ)+ST(II,II-1)*B(JJ-1)
  DO 190 I=1,3
  SIG(II)=0.
  DO 190 J=1,3
190 SIG(II)=SIG(II)+C(II,J)*RR(J)
  DO 195 I=1,3
  II=I+4
195 SIG(N,II)=SIG(N,II)+SIG(II)
  DD=(SIG(N,5)-SIG(N,6))/2.
  AJ2=(DD*DD+SIG(N,7)**2)/11.-3.*(EE(4)**2)
  AJ2=DSQRT(AJ2)
  AJ1=1.5*(SIG(N,5)+SIG(N,6))-3.*EE(4)*AJ2
  FAIL=AJ2+EE(4)*AJ1
  IF(MTAG(N).EQ.0) GO TO 200
  IF(MTAG(N).EQ.2) GO TO 300
  DD=DABS(FAIL-EE(3))
  CHECK=.02*EE(3)
  IF(DD-CHECK) 300,300,196
196 KKK=1

```

```

CR=CHECK/ND
IF((CR-SM) 147,300,300
197 SM=CH
GO TO 300
200 CONTINUE
IF(FAIL.LT.EE(3)) GO TO 300
KK=KK+1
ND=(SIGI(N,1)-SIGI(N,2))/2.
BJ2=(DD*ND+SIGI(N,3)**2)/(1.-3.*(EE(4)**2))
BJ2=DSQRT(BJ2)
RJ1=1.5*(SIGI(N,1)+SIGI(N,2))-3.*EE(4)*BJ2
AAA=BJ2*BJ2
BBB=AJ2*AJ2
FFF=1.-3.*EE(4)*EE(4)
CCC=(SIGI(N,1)-SIGI(N,2))*(SIGI(N,5)-SIGI(N,6))/4.+SIGI(N,3)*SIGI
*(N,7)
DDD=SIGI(N,1)+SIGI(N,2)
GGG=SIGI(N,5)+SIGI(N,6)
FF=FFF*FFF
GG=GGG*GGG
DD=DDD*DDD
EF=2.25*EE(4)*EE(4)
AA=AAA*FF-EF*DD
RR=BBB*FF-EF*GG
CC=CCC*FFF-EF*DDD*GGG
DD=1.5*EE(4)*EE(3)*DDD
FF=1.5*EE(4)*EE(3)*GGG
GG=EE(3)*EE(3)
AAA=AA+BB-2.*CC
RRR=AA-CC+DD-FF
CCC=2.*DD-GG+AA
GGG=BBB*BBB-AAA*CCC
GGG=DSQRT(GGG)
IF(AAA) 220,210,220
210 RATIO(N)=.5*CCC/BBB
GO TO 300
220 AA=RRR/AAA
BB=DARS(GGG/AAA)
RATIO(N)=AA-BB
IF(RATIO(N).LT.0.) RATIO(N)=AA+BB
IF(RATIO(N).GE.1.) RATIO(N)=.99999
IF(RATIO(N).LT.0.) RATIO(N)=0.
300 CONTINUE
C
IF(KK.EQ.0) GO TO 420
DO 350 N=1,NUMEL
IF(MTAG(N).GT.0) GO TO 350
ND=RATIO(N)

```

```

      IF(DI)-SR) 305,350,350
305 SR=DU
      NUMR=N
      KKK=1
350 CONTINUE
      IF(SR.LT.0.1) SR=0.1
      IF(NUMR.EQ.0) GO TO 420
C
260 DO 370 N=1,NUMEL
      IF(MTAG(N).GT.0) GO TO 370
      IF(RATIO(N).LE.SR) GO TO 355
      DD=RATIO(N)-SR
      IF(DD-.05) 355,355,370
355 MTAG(N)=1
370 CONTINUE
C
420 CONTINUE
      DO 410 N=1,NUMNP
      II=2*N-1
      R(II)=B(II)*SR
410 R(II+1)=B(II+1)*SR
      DO 600 N=1,NUMEL
      I=IX(N,1)
      J=IX(N,2)
      K=IX(N,3)
      L=IX(N,4)
      MTYPE=IX(N,5)
      IF(K.EQ.L) GO TO 460
      XC=(R(I)+R(J)+R(K)+R(L))/4.
      YC=(Z(I)+Z(J)+Z(K)+Z(L))/4.
      GO TO 470
460 XC=(R(I)+R(J)+R(K))/3.
      YC=(Z(I)+Z(J)+Z(K))/3.
470 CONTINUE
      DO 450 I=1,3
      II=I+4
      SIG(I)=SIG(N,II)-SIG(N,I)
      SIG(N,II)=SIG(I)*(1.-SR)
      SIG(N,II)=-SIG(N,II)
      SIG(I)=SIG(I)*SR+SIG(N,I)
450 SIG(N,I)=SIG(I)
      SIG(7)=EF(2)*(SIG(1)+SIG(2))
      DO 455 I=1,4
455 EE(I)=E(I,MTYPE)
      IF(MTAG(N).EQ.0) GO TO 500
      CC=DTAN(EE(4)/57.296)
      BB=DSQRT(9.0+12.0*CC*CC)
      EE(4)=CC/BB

```

```

DD=(SIG(1)-SIG(2))/2.
BJ2=(DD*DD+SIG(3)**2)/(1.-3.*(EE(4)**2))
BJ2=DSQRT(BJ2)
SIG(7)=.5*(SIG(1)+SIG(2))-3.*EE(4)*BJ2
500 SIG(N,4)=SIG(7)
CC=(SIG(1)+SIG(2))/2.
BB=(SIG(1)-SIG(2))/2.
CR=DSQRT(BB**2+SIG(3)**2)
SIG(4)=CC+CR
SIG(5)=CC-CR
SIG(6)=0.0
IF ((BB.EQ.0.0).AND.(SIG(3).EQ.0.0)) GO TO 510
SIG(6)=28.698*DATAN2(SIG(3),BB)
510 CONTINUE
IF(MPRINT) 520,520,550
520 WRITE(6,2000) NNN
MPRINT=50
550 MPRINT=MPRINT-1
WRITE(6,2001) N, XC, YC, (SIG(I), I=1,7), MTAG(N)
IF(MTAG(N).EQ.0) GO TO 560
DD=(SIG(N,5)-SIG(N,6))/2.
BJ2=(DD*DD+SIG(N,7)**2)/(1.-3.*(EE(4)**2))
BJ2=DSQRT(BJ2)
TOL=TOL+BJ2
GO TO 600
560 SIG(7)=EE(2)*(SIG(N,5)+SIG(N,6))
BJ2= DSQRT(((SIG(N,5)-SIG(N,6))**2+(SIG(N,6)-SIG(7))**2+(SIG(7)
*-SIG(N,5))**2)/6.0+SIG(N,7)**2)
TOL=TOL+BJ2
600 CONTINUE
SR2= (SR1*SR) + SR2
SR1= (1.0-SR) * SR1
WRITE(6,2002) TOL,SR ,NUMR, KK, SR2
IF(TOL-1.) 660,660,650
650 KKK=1
GO TO 700
660 KKK=0
700 CONTINUE
II=0
DO 710 I=1, NUMEL
710 IF(MTAG(I).GE.1) II=II+1
IF(KK.GE.NUMEL) CALL EXIT
800 RETURN
C
2000 FORMAT(1H1/
*36H STRESSES AFTER APPROXIMATION NUMBER 14////
*7H EL.NO. 7X 1HX 7X 1HY 4X 8HX-STRESS 4X 8HY-STRESS 3X 9HXY-STRESS
* 2X 10HMAX-STRESS 2X 10HMIN-STRESS 7H ANGLE 4X 8HZ-STRESS 5X 7HPL

```

*ASTIC)

2001 FORMAT (I7,2F8.2,1P5E12.4,0P1F7.2,1PE12.4,1I2)

2002 FORMAT(39H0THE UNBALANCED LOAD AT THIS STAGE IS E14.5//

*47H THE RATIO FOR CORRECTION OF STORED STRESSES IS F10.4//

*31H THE NEXT ELEMENT YIELDING IS I4/

*91H AND THE TOTAL NUMBER OF ELEMENTS THAT CAN YIELD WITH THE LINEA

*R ADDITION OF TOTAL LOAD IS I4/

*50H LOAD UP TO THIS STAGE AS A FRACTION OF TOTAL IS F20.5)

C

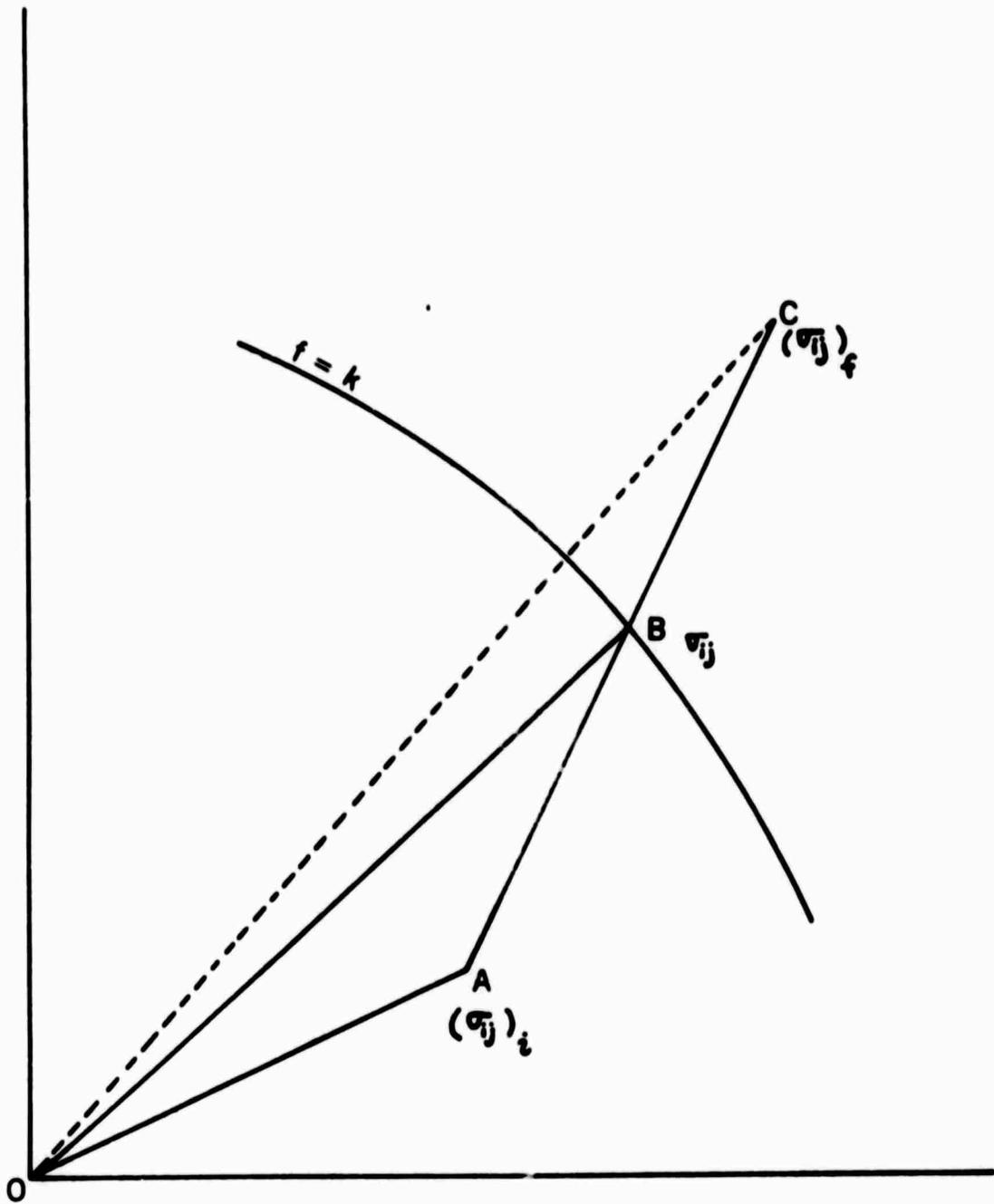
END

3.4 Example

Naghdi (1957) solved the problem of an elastic-perfectly plastic wedge under uniform loading on one face (Fig. 3.2). Plane strain conditions were considered. The wedge material was assumed to yield according to Von Mises' yield criterion. This type of material is a special case of Mohr-Coulomb material having the angle of internal friction $\phi = 0$.

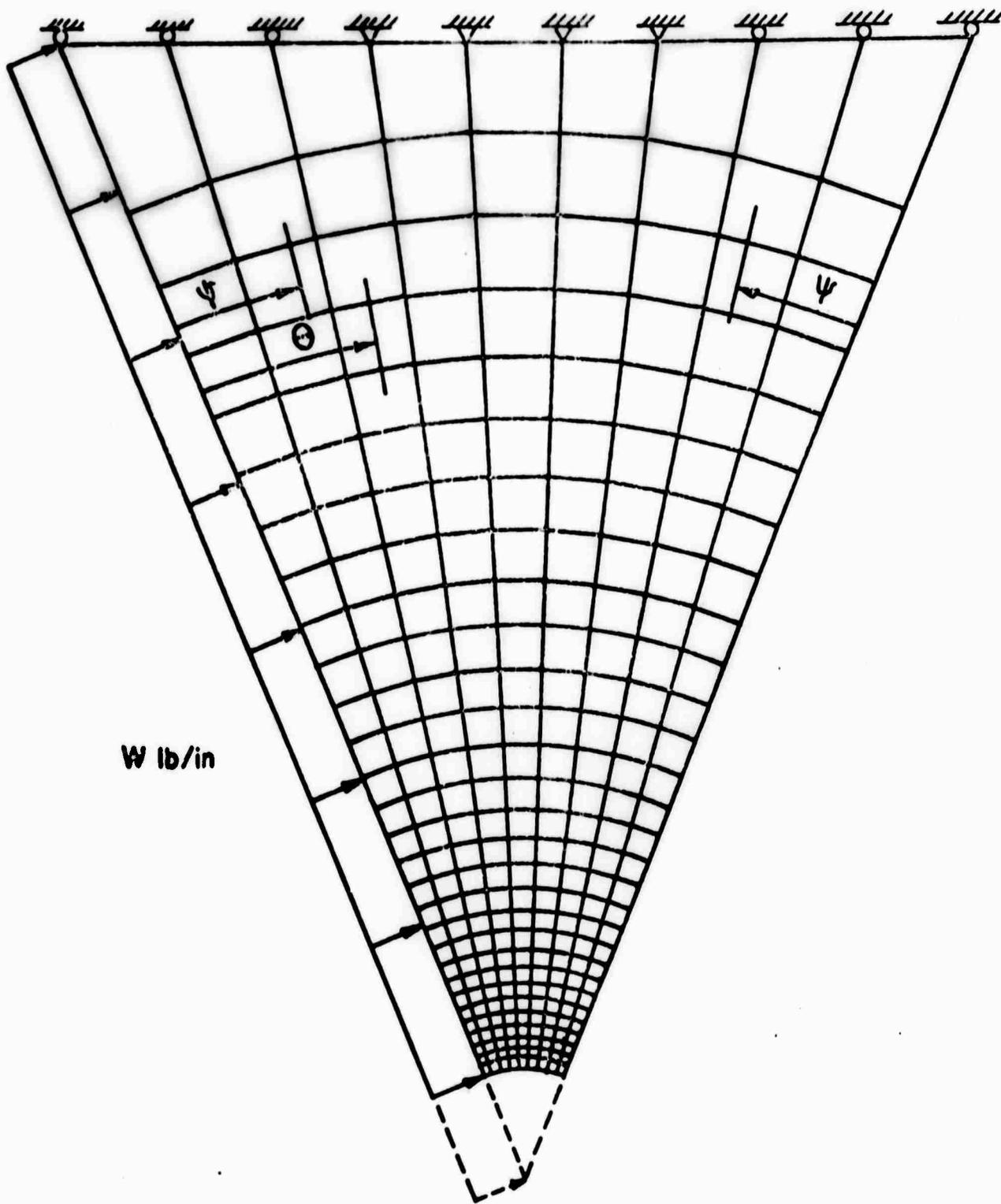
Figures 3.3 and 3.4 show the theoretical and computed results for the distribution of radial and circumferential stress at various stages of loadings. The angle ψ denotes the angle upto which the yielding has progressed from the boundaries. Fig. 3.5 shows the radial strain distribution at various stages.

Generally the agreement between results computed by the method outlined and the exact analysis were found to be good.



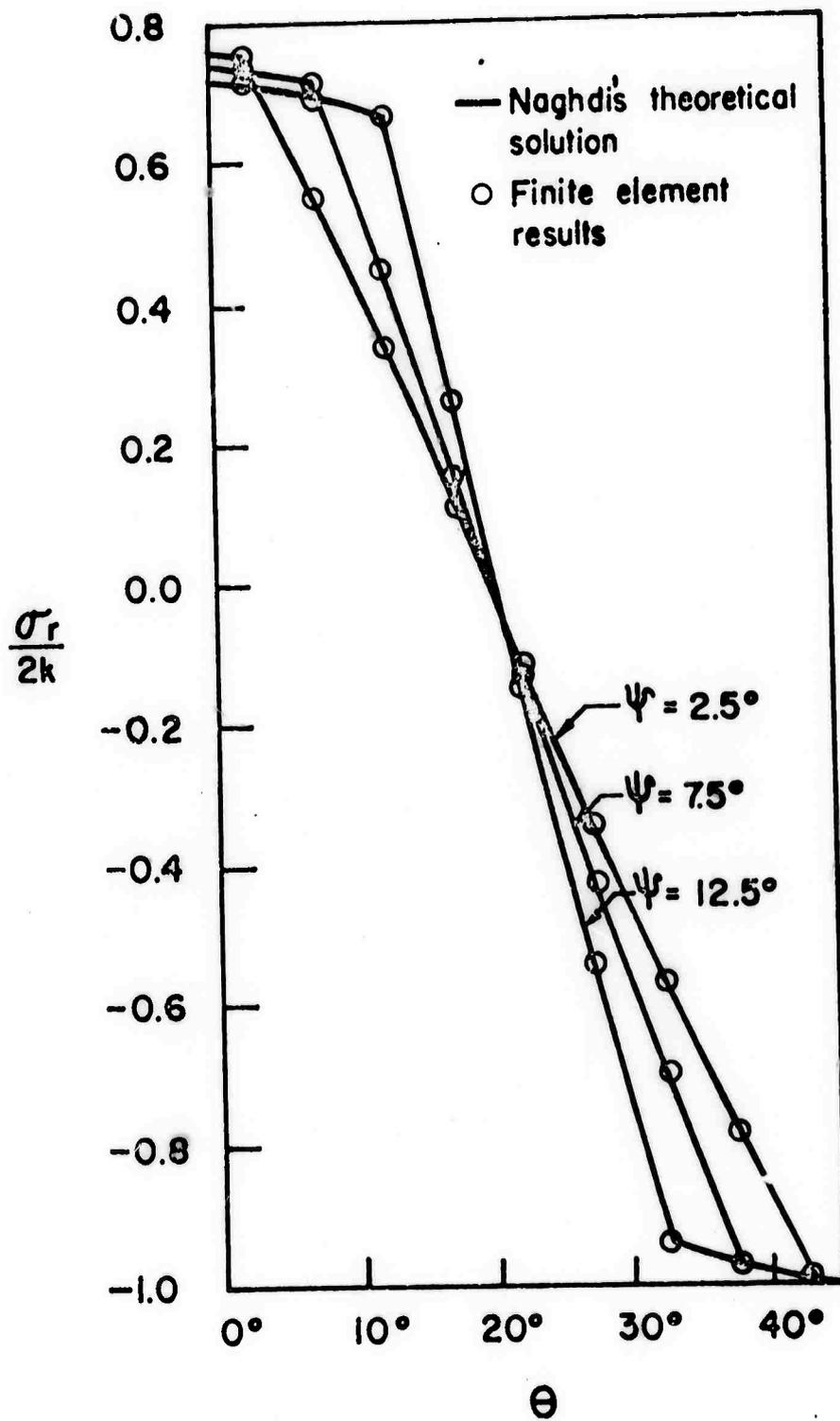
CALCULATION OF STRESS RATIO

FIG 3.1



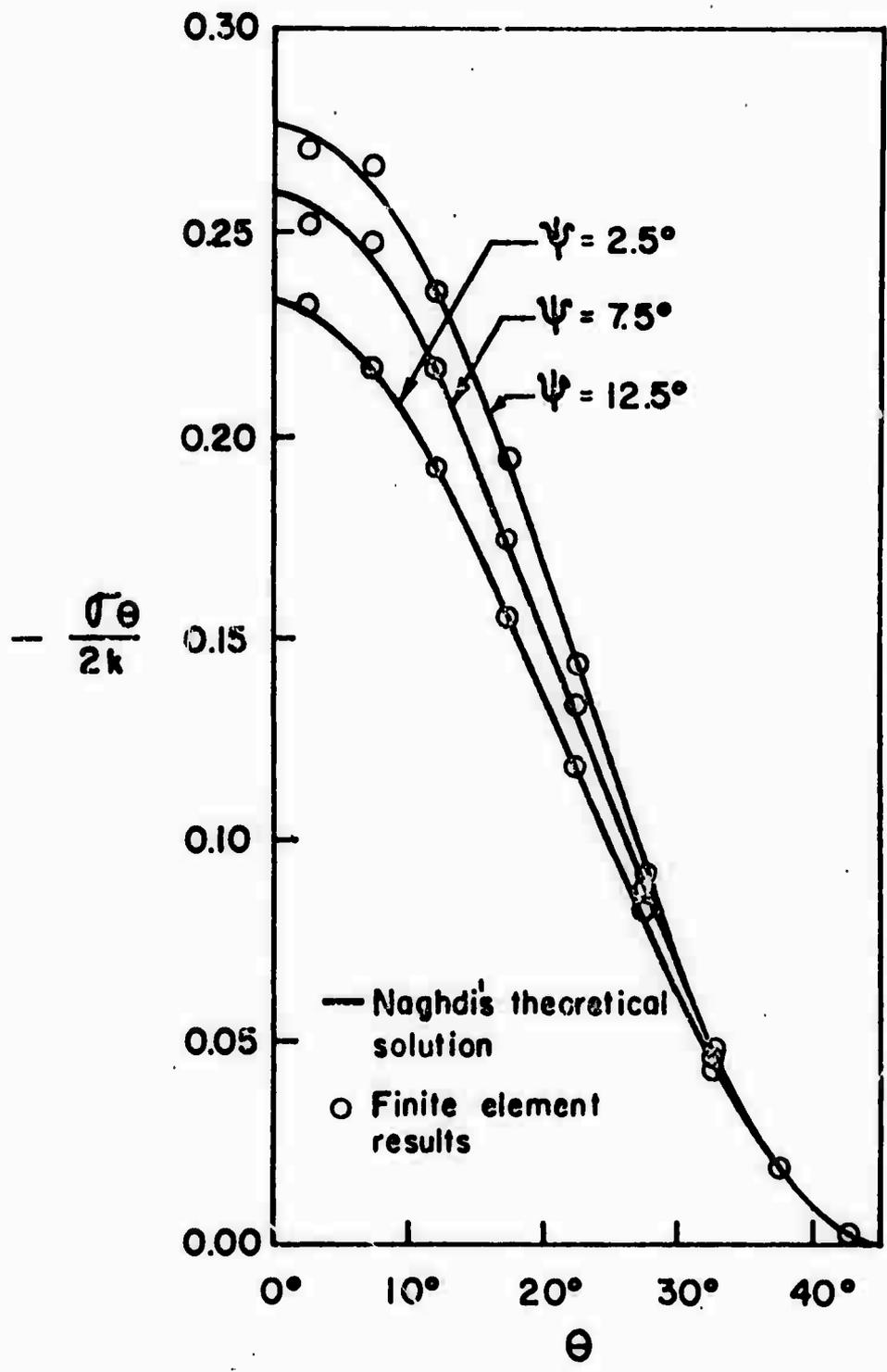
FINITE ELEMENT IDEALIZATION FOR ELASTIC-PLASTIC WEDGE

FIG 3.2



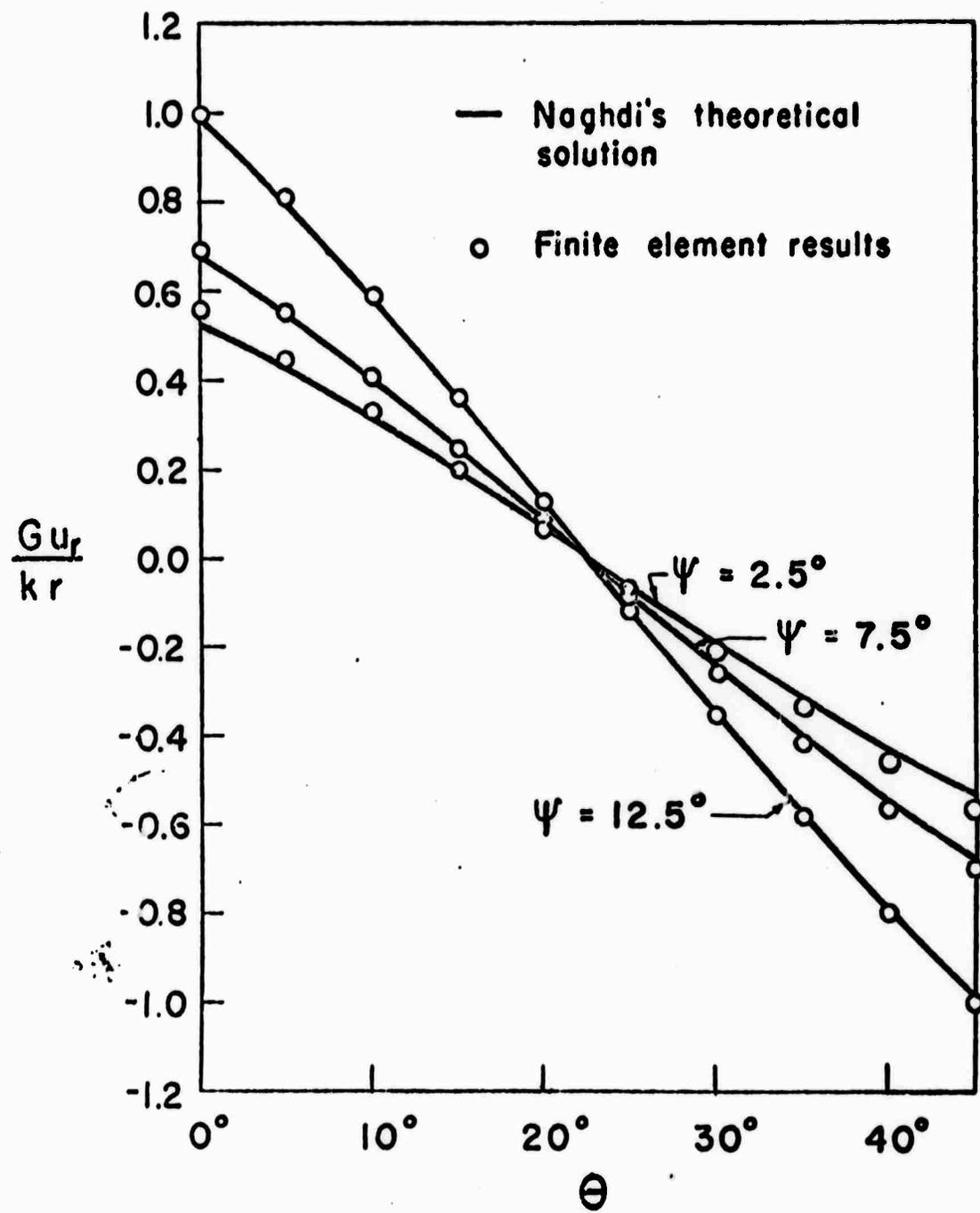
DISTRIBUTION OF RADIAL STRESS FOR WEDGE AT VARIOUS LOADS

FIG 3.3



DISTRIBUTION OF CIRCUMFERENTIAL STRESS FOR WEDGE AT VARIOUS LOADS

FIG 3.4



DISTRIBUTION OF RADIAL DISPLACEMENTS

FIG 3.5

CHAPTER IV

COMPUTER PROGRAM FOR ANALYSIS OF JOINTED ROCK

Chapter IV. Computer Program for Analysis of Jointed Rock

4.1 Organization

The computer program described here is based on the theory described in Chapter I and II. The rock mass is considered as a linear elastic material in the direction of compressive stresses and is assumed to have no resistance to deformation in the direction of principal stresses. The program corrects the discrepancy in the method presented by Zienkiewicz et al (1968). This was pointed out towards the end of Chapter I. It makes allowance for the boundary conditions, residual stresses, stresses due to temperature change, and varying pressure boundaries. This program also uses the quadrilateral elements and generates stiffness the same way as that described in Chapter III.

The principal program called MAIN controls all the data input and control information. It does the system initialization, prints the control data, geometrical and material properties. MAIN calls the subroutines for stiffness, solution of equations and stress calculations.

4.12. Stiffness Matrix

Stiffness matrix for the analysis is computed in blocks by the subroutine STIFF. For element stiffness, it calls QUAD for triangular and quadrilateral elements and ONED for bar elements. Direct stiffness technique was used to get the total stiffness. Equations are modified for displacement boundary conditions.

4.13. Load Matrix

Load matrix for the analysis is computed in LOAD subroutine. The varying pressure boundary is taken into account. The load matrix is modified in the LOAD subroutine for each iteration performed. This accounts for the nonlinearity introduced by progressive cracking by considering the change in element stiffness as a pseudo load.

4.14. Calculations of Displacements

After the stiffness and load matrices for a stage have been computed, the resulting equations are solved by calling subroutine SYMBAN. This subroutine was developed to take the advantage of the fact that stiffness matrix remains the same throughout. Gaussian elimination technique is used to solve the equations. The elimination is done once for all and the reduced matrix stored on auxiliary units. Solution for each iteration consists of back-substitution only. This approach results in considerable economy in machine time.

4.15. Calculations of Stresses

After the displacements have been computed, the stresses are computed using the constitutive law. A change in the minor principal stress corresponding to Poisson's effect has been introduced when one of the principal stresses was compressive and others tensile.

4.2. Input Data Preparation

1. Control Card (A6). This card will carry the characters START in columns 1-5. This will start the processing of the data deck which consists of the following set of cards.
2. Job Title (72H). This card will give the descriptive identification for the job.
3. Control Information (4I5, 3F10.2, I5, E15.4)

<u>Information</u>	<u>Columns</u>
Total number of nodal points	1 - 5
Total number of elements	6 - 10
Number of different materials	11 - 15
Number of pressure boundary cards	16 - 20
Body Force in X-direction	21 - 30
Body Force in Y-direction	31 - 40
Reference (stress-free) temperature	41 - 50
Number of Iterations	51 - 55
Tolerance to Convergence	56 - 70

4. Material Property Cards. One set of cards must be provided for each material.

In each set:

- a. First card (2I5, F10.3, I5) will give the following information:

Material identification number	1 - 5
Number of temperature cards (8 maximum)	6 - 10
Mass density of the material	11 - 20

Material code to designate materials which cannot take tension 21 - 26

code = 1 for materials which cannot take tension
0 for materials which can take tension.

b. Subsequent cards, one for each temperature, the number being defined in columns 6-10 of the first card, will carry the following information (4F10.0):

<u>Information</u>	<u>Columns</u>
Temperature	1 - 10
Elastic modulus	11 - 20
Poisson's ratio	21 - 30
Coefficient of thermal expansion	31 - 40

5. Nodal Point Cards (I5, F5.0, 5F10.0)

One card for each nodal point with the following information:

Nodal point number	1 - 5
Type of nodal point	6 - 10
X-ordinate	11 - 20
Y-ordinate	21 - 30
XR	31 - 40
XZ	41 - 50
Temperature	51 - 60

If the number in columns 6-10 is

Zero XR is the specified X-load and XZ is the specified Y-load

1 XR is the specified X-displacement and XZ is the specified Y-load

2 XR is the specified X-load and XZ is the specified Y-displacement

3 XR is the specified X-displacement and XZ is the specified Y-displacement

All loads are considered to be total forces acting on an element of unit thickness. Nodal point cards must be in numerical sequence. If cards are omitted, the omitted nodal points are generated at equal intervals along a straight line between the defined nodal points. The necessary temperatures are determined by linear interpolation. The type of the nodal point, as well as XR, XZ, are set equal to zero.

6. Element Material Cards (12I5)

These cards shall carry the material type of all the elements. Each card shall have material types for 12 elements in sequence. The material type for each element must be read in as no interpolation has been provided for.

7. Element Cards (5I5, 5X, 3 F10.0)

One card for each element will provide the following data.

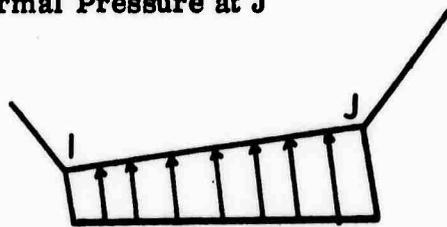
<u>Information</u>	<u>Columns</u>
Number of element	1 - 5
Nodal point I	6 - 10
Nodal point J	11 - 15
Nodal point K	16 - 20
Nodal point L	21 - 25
Initial stresses:	
(i) component in x-direction	31 - 40
(ii) component in y-direction	41 - 50
(iii) shearing stress on x-y planes	51 - 60

Nodal points I, J, K, L are corners of each individual element in a counter-clockwise order for a right handed system of coordinates. For triangular elements set nodal point L same as nodal point K. The element cards must be in the numerical sequence. Any cards that are omitted will be automatically generated in the program by incrementing each of the I, J, K, and L nodal points by one.

8. Pressure Boundary Cards (2I5, F10.0)

One card for each boundary element which is subjected to a normal pressure will carry the following information:

<u>Information</u>	<u>Columns</u>
Nodal Point I	1 - 5
Nodal Point J	6 - 10
Normal Pressure at I	11 - 20
Normal Pressure at J	21 - 30



As shown in the sketch, the boundary element must be on the left as one progresses from I to J. Surface tensile force is input as a negative pressure.

Output Information:

The following information is developed and printed by the program:

1. Reprint of input data
2. Nodal point displacements
3. Stresses at the center of each element

4.3 Program Listing

```

*****
* TWO DIMENSIONAL ANALYSIS OF A NO TENSION SYSTEM *
*****
DOUBLE PRECISION S,C,B,A,P,ST,SIG,U,V,CC,OB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FDR,TOL,UR,UZ ,SIGI,VOL,COM,WORD(2),WORD1,E,COSA,SINA
2,CU,ZZ,EL,XC,YC,PR,ACELR,ACELZ,TEMP,T, R,Q,RD,Z
3,DR,DZ,DT,ZX
COMMON ACELR,ACELZ,TEMP,Q,RD(12),R(900),Z(900),T(900),PR(200,2),
1NUMNP,NUMEL,NUMMAT,NUMPC,MTYPE,LLL,N,HED(18),NTC(10),CODE(900),
2MTC(10),IBC(200),JBC(200),NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIGI(900,6),CU(1800),TOL,VOL
1,F(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,FE(3),LM(4),IX(800,5)
COMMON /BANARG/ MBAND,NUMBLK,B(1800),A(108,54),KKK,JA
DATA WORD/6HSTART ,6HSTOP /
DEFINE FILE 1(50,1500,U,NRK),2(800,256,U,10)
CALL FPRSET(208,256,-1,1)
10 READ(5,1006) WORD1
IF(WORD1.EQ.WORD(1)) GO TO 30
IF(WORD1.EQ.WORD(2)) STOP
GO TO 10
30 READ (5,1000) HED,NUMNP,NUMEL,NUMMAT,NUMPC,ACELR,ACELZ,Q,NP,TOL
WRITE(6,2000) HED,NUMNP,NUMEL,NUMMAT,NUMPC,ACELR,ACELZ,Q,NP,TOL
40 DO 50 M=1,NUMMAT
READ (5,1001) MTYPE,NTC(MTYPE),RD(MTYPE),MTC(MTYPE)
WRITE(6,2001) MTYPE,NTC(MTYPE),RD(MTYPE),MTC(MTYPE)
NUMTC=NTC(MTYPE)
READ (5,1002) ((E(I,J,MTYPE),J=1,4),I=1,NUMTC)
WRITE (6,2002) ((E(I,J,MTYPE),J=1,4),I=1,NUMTC)
50 CONTINUE
WRITE (6,2003)
L=0
R(L)=0.
Z(L)=0.
T(L)=0.
60 READ (5,1003) N,CODE(N),R(N),Z(N),UR(N),UZ(N),T(N)
NL=L+1
ZX=N-L
DR=(R(N)-R(L))/ZX
DZ=(Z(N)-Z(L))/ZX
DT=(T(N)-T(L))/ZX
70 L=L+1
IF(N-L) 100,90,80
80 CODE(L)=0.0
R(L)=R(L-1)+DR
Z(L)=Z(L-1)+DZ

```

```

T(L)=T(L-1)+DT
UR(L)=C.0
UZ(L)=C.0
GO TO 70
90 WRITE (6,2004) (K,CODE(K),R(K),Z(K),UR(K),UZ(K),T(K),K=NL,N)
IF(NUMNP-N) 100,110,60
100 WRITE (6,2005) N
CALL EXIT
110 CONTINUE
READ(5,1007) (IX(N,5),N=1,NUMEL)
WRITE (6,2006)
N=0
130 READ(5,1004) M,(IX(M,I),I=1,4),(SIGI(M,I),I=1,3)
IF(M.EQ.1) GO TO 140
ZX=M-N
DO 135 I=1,3
135 SIG(I)=(SIGI(M,I)-SIGI(N,I))/ZX
140 N=N+1
IF (M-N) 170,170,150
150 IX(N,1)=IX(N-1,1)+1
IX(N,2)=IX(N-1,2)+1
IX(N,3)=IX(N-1,3)+1
IX(N,4)=IX(N-1,4)+1
DO 160 I=1,3
160 SIGI(N,I)=SIGI(N-1,I)+SIG(I)
170 WRITE(6,2007) N,(IX(N,I),I=1,5),(SIGI(N,I),I=1,3)
IF (M-N) 180,180,140
180 IF (NUMEL-N) 190,190,130
190 CONTINUE
IF (NUMPC) 290,310,290
290 WRITE (6,2008)
DO 300 L=1,NUMPC
READ(5,1005) IBC(L),JBC(L),PR(L,1),PR(L,2)
300 WRITE(6,2009) IBC(L),JBC(L),PR(L,1),PR(L,2)
310 CONTINUE
DO 440 N=1,NUMNP
NN=2*N
CU(NN-1)=0.
440 CU(NN)=0.
J=0
DO 340 N=1,NUMEL
DO 340 I=1,4
DO 325 L=1,4
KK=IABS(IX(N,I)-IX(N,L))
IF (KK-J) 325,325,320
320 J=KK
325 CONTINUE
340 CONTINUE

```

```

MHAND=2*J+2
WRITE(6,2012) MHAND
350 CALL STIFF
NCHECK=1
KKK=1
CALL SYMBAN
KKK=2
DO 500 LLL=1,NP
CALL LOAD
CALL SYMBAN
DO 400 N=1,NUMNP
NN=2*N
CU(NN-1)=CU(NN-1)+B(NN-1)
400 CU(NN)=CU(NN)+B(NN)
WRITE(6,2013) LLL
WRITE(6,2010) (N,B(2*N-1),B(2*N),CU(2*N-1),CU(2*N), N = 1,NUMNP)
CALL STRESS
IF(NCHECK.EQ.0) GO TO 600
500 CONTINUE
GO TO 700
600 WRITE(6,2011) LLL
700 GO TO 10
999 FORMAT(I5)
1000 FORMAT (18A4/4I5,3F10.2,I5,E15.4,I5)
1001 FORMAT (2I5,1F10.3,I5)
1002 FORMAT (5F10.3)
1003 FORMAT (I5,F5.1,5F10.4)
1004 FORMAT (5I5,5X,3F10.4)
1005 FORMAT(2I5,2F10.3)
1006 FORMAT(A6)
1007 FORMAT(12I5)
2000 FORMAT (1H1 18A4/
1 30H0 NUMBER OF NODAL POINTS----- 13 /
2 30H0 NUMBER OF ELEMENTS----- 13 /
3 30H0 NUMBER OF DIFF. MATERIALS--- 13 /
4 30H0 NUMBER OF PRESSURE CARDS---- 13 /
5 30H0 X-ACCELERATION----- E12.4/
6 30H0 Y-ACCELERATION----- E12.4/
7 30H0 REFERENCE TEMPERATURE----- E12.4/
8 30H0 NO. OF APPROXIMATIONS----- 15/
9 30H0 TOLERANCE FOR CONVERGENCE--- E12.4)
2001 FORMAT (17H0 MATERIAL NUMBER= 13, 30H, NUMBER OF TEMPERATURE CARDS=
1 13, 15H, MASS DENSITY= E12.4,16H, MATERIAL CODE= 15)
2002 FORMAT (15H0 TEMPERATURE 10X 5HEC 9X 6HNU 10X 5HALPHA/
1(F15.2,3E15.5))
2003 FORMAT (108H1NODAL POINT TYPE X ORDINATE Y ORDINATE X LO
1AD OR DISPLACEMENT Y LOAD OR DISPLACEMENT TEMPERATURE )
2004 FORMAT (112,F12.2,2F12.5,2E24.7,F12.3)

```

```

2005 FORMAT (26HONODAL POINT CARD ERROR N= 15)
2006 FORMAT (29H1ELEMENT NO.      I      J      K      L      MATERIAL
1SIGIXX      SIGIYY      SIGIXY
2007 FORMAT (11I3,4I6,11I2,3F12.3)
2008 FORMAT(29HOPRESSURE BOUNDARY CONDITIONS/40H      I      J      PRESSUR
1F I      PRESSURE J)
2009 FORMAT(2I6,2F14.3)
2010 FORMAT(12HON.P.NUMBER 17X 3HDUX 17X 3HDUY 18X 2HUX 18X 2HUY/
1(11I2,4F20.7))
2011 FORMAT(35H0 NUMBER OF CYCLES TO CONVERGENCE = 15)
2012 FORMAT(28H BAND WIDTH FOR THIS DATA = 15)
2013 FORMAT(30H1 RESULTS OF ITERATION NO. 15//)
END

```

SUBROUTINE STIFF

```

C
DOUBLE PRECISION S,C,B,A,P,SI,SIG,U,V,CC,BB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FOR,TOL,UR,UZ ,SIGI,VOL,COM,E,COXA,SINA
2,CU,ZZ,EE,XC,YC,PR,ACELR,ACELZ,TEMP,T, R,Q,RO,Z
COMMON ACELR,ACELZ,TEMP,Q,RO(12),R(900),Z(900),T(900),PR(200,2),
1NUMNP,NUMEL,NUMMAT,NUMPC,MTYPE,LLL,N,HED(18),NTC(10),CODE(900),
2MTC(10),IBC(200),JBC(200),NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIGI(900,6),CU(1800),TOL,VOL
1,E(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,EE(3),LM(4),IX(800,5)
COMMON /BANARG/ MBAND,NUMBLK,B(1800),A(108,54),KKK,JA
DEFINE FILE 1(50,1500,U,NBK),2(800,256,U,ID)
NB=27
ND=2*NB
ND2=2*ND
STOP=0.0
NUMBLK=0
JA=ND2*(MBAND+1)/1500+1
NBK=1
DO 50 N=1,ND2
DO 50 M=1,ND
50 A(N,M)=0.0
60 NUMBLK=NUMBLK+1
NH=NB*(NUMBLK+1)
NM=NH-NB
NL=NM-NB+1
KSHIFT=2*NL-2
DO 210 N=1,NUMEL
IF (IX(N,5)) 210,210,65
65 DO 80 I=1,4
IF (IX(N,I)-NL) 80,70,70
70 IF (IX(N,I)-NM) 90,90,80
80 CONTINUE
GO TO 210
90 IF(IX(N,3)-IX(N,2)) 95,85,95
85 CALL ONED
IX(N,5)=-IX(N,5)
MM=2
GO TO 130
95 CALL QUAD
IX(N,5)=-IX(N,5)
ID=N
WRITE(2*ID) ((C(IJK,JJK),JJK=1,3),EE(IJK),IJK=1,3)
1 ,((S(JJI,KKI),KKI=1,8),JJI=1,8),((ST(IKK,JKK),JKK=1,8),IKK=1,3)
2 ,(RR(JII),ZZ(JII),JII=1,4),XC,YC,TEMP,VOL
IF(VOL) 100,100,110

```

```

100 WRITE(6,2000) N
    STOP=1.0
110 MM=4
    IF(IX(N,3)-IX(N,4)) 130,120,130
120 MM=3
130 DO 140 I=1,MM
140 LM(I)=2*IX(N,I)-2
    DO 200 I=1,MM
    DO 200 K=1,2
    II=LM(I)+K-KSHIFT
    KK=2*I-2+K
    DO 200 J=1,MM
    DO 200 L=1,2
    JJ=LM(J)+L-II+1-KSHIFT
    LL=2*J-2+L
    IF(JJ) 200,200,175
175 IF(NJ)-JJ) 180,195,195
180 WRITE(6,2001) N
    STOP=1.0
    GO TO 210
195 A(II,JJ)=A(II,JJ)+S(KK,LL)
200 CONTINUE
210 CONTINUE
    DO 400 M=NL,NH
    IF(M-NUMNP) 315,315,400
315 U=UR(M)
    N=2*M-1-KSHIFT
    IF(CODE(M)) 390,400,316
316 IF(CODE(M)-1.) 317,370,317
317 IF(CODE(M)-2.) 318,390,318
318 IF(CODE(M)-3.) 390,380,390
370 CALL MODIFY(A,ND2,MBAND,N)
    GO TO 400
380 CALL MODIFY(A,ND2,MBAND,N)
390 U=UZ(M)
    N=N+1
    CALL MODIFY(A,ND2,MBAND,N)
400 CONTINUE
    WRITE(1*NRK) ((A(N,M),M=1,MBAND),N=1,ND)
    NRK=NRK+JA
    DO 420 N=1,ND
    K=N+ND
    DO 420 M=1,ND
    A(N,M)=A(K,M)
420 A(K,M)=0.0
    IF(NM.LT.NUMNP) GO TO 60
    DO 440 N=1,NUMEL
    IX(N,5)=IABS(IX(N,5))

```

480 CONTINUE
IF (STOP) 490,500,490
490 CALL EXIT
500 RETURN
2000 FORMAT (26HNEGATIVE AREA ELEMENT NO. 14)
2001 FORMAT (29HORAND WIDTH EXCEEDS ALLOWABLE 14)
FND

SUBROUTINE UNED

C

```
DOUBLE PRECISION S,C,B,A,P,ST,SIG,U,V,CC,RR,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FOR,TOL,UR,UZ ,SIGI,VOL,COM,E,COSA,SINA
2,CU,ZZ,EE,XC,YC,PR,ACELR,ACELZ,TEMP,T, R,Q,RO,Z
COMMON ACELR,ACELZ,TEMP,Q,RO(12),R(900),Z(900),T(900),PR(200,2),
INUMNP,NUMEL,NUMMAT,NUMPC,MTYPE,LLL,N,HED(18),NTC(10),CODE(900),
2MTC(10),IBC(200),JBC(200),NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIGI(900,6),CU(1800),TOL,VOL
1,E(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,EE(3),LM(4),IX(800,5)
COMMON /BANARG/ MBAND,NUMBLK,B(1800),A(108,54),KKK,JA
```

C

```
DO 100 I=1,8
P(I)=0.0
DO 100 J=1,8
100 S(I,J)=0.0
MTYPE=IX(N,5)
I=IX(N,1)
J=IX(N,2)
DX=R(J)-R(I)
DY=Z(J)-Z(I)
XL=DSQRT(DX**2+DY**2)
COSA=DX/XL
SINA=DY/XL
COMM=E(1,2,MTYPE)*E(1,4,MTYPE)/XL
```

C

```
S(1,1)=COSA*COSA*COMM
S(1,2)=COSA*SINA*COMM
S(1,3)=-S(1,1)
S(1,4)=-S(1,2)
S(2,1)=S(1,2)
S(2,2)=SINA*SINA*COMM
S(2,3)=-S(1,2)
S(2,4)=-S(2,2)
S(3,1)=S(1,3)
S(3,2)=S(2,3)
S(3,3)=S(1,1)
S(3,4)=S(1,2)
S(4,1)=S(1,4)
S(4,2)=S(2,4)
S(4,3)=S(3,4)
S(4,4)=S(2,2)
```

C

C

C

RETURN

```

SUBROUTINE QUAD
DOUBLE PRECISION S,C,B,A,P,ST,SIG,U,V,CC,BB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FOR,TOL,UR,UZ,SIG1,VOL,COM,E,COSA,SINA
2,CU,ZZ,EE,XC,YC,PR,ACELR,ACELZ,TEMP,T,R,Q,RO,Z
3,U,V,XT,XS,RATIO,DEN,XNT
COMMON ACELR,ACELZ,TEMP,Q,RO(12),R(900),Z(900),T(900),PR(200,2),
1NUMNP,NUMEL,NUMMAT,NUMPC,MTYPE,LLL,N,HED(18),NTC(10),CODE(900),
2MTC(10),IBC(200),JBC(200),NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIG1(900,6),CU(1800),TOL,VOL
1,E(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,EE(3),LM(4),IX(800,5)
COMMON /BANARG/ MBRAND,NUMBLK,B(1800),A(108,54),KKK,JA

```

C.

```

DIMENSION U(3),V(3)
I=IX(N,1)
J=IX(N,2)
K=IX(N,3)
L=IX(N,4)
MTYPE=IX(N,5)
VOL=0.
TEMP=(T(I)+T(J)+T(K)+T(L))/4.C
RATIO=0.0
NUMTC=NTC(MTYPE)
IF (NUMTC.EQ.1) GO TO 100
DO 50 M=2,NUMTC
IF (E(M,1,MTYPE)-TEMP) 50,60,60
50 CONTINUE
60 DEN=E(M,1,MTYPE)-E(M-1,1,MTYPE)
IF(DEN) 70,80,70
70 RATIO=(TEMP-E(M-1,1,MTYPE))/DEN
80 DO 90 KK=1,3
90 EF(KK)=E(M-1,KK+1,MTYPE)+RATIO*(E(M,KK+1,MTYPE)-E(M-1,KK+1,MTYPE))
GO TO 110
100 DO 105 KK=1,3
105 EF(KK)=E(1,KK+1,MTYPE)
110 CONTINUE
COMM=EE(1)/(1.-EE(2)**2)
C(1,1)=COMM
C(1,2)=COMM*EF(2)
C(1,3)=0.
C(2,1)=C(1,2)
C(2,2)=C(1,1)
C(2,3)=0.
C(3,1)=0.
C(3,2)=0.
C(3,3)=.5*COMM*(1.-EE(2))
DO 130 J=1,10

```

```

      DO 120 I=1,3
120  ST(I,J)=0.
      DO 130 I=1,10
130  S(I,J)=0.
      DO 140 I=1,4
      NPP=IX(N,I)
      RR(I)=R(NPP)
140  ZZ(I)=Z(NPP)
      IF(IX(N,3)-IX(N,4)) 145,150,145
145  XC=(RR(1)+RR(2)+RR(3)+RR(4))/4.
      YC=(ZZ(1)+ZZ(2)+ZZ(3)+ZZ(4))/4.
      RR(5)=XC
      ZZ(5)=YC
      K=5
      J=1
      I=4
      LM(3)=9
      NT=4
      GO TO 160
150  NT=1
      LM(3)=5
      I=1
      K=3
      J=2
      XC=(RR(1)+RR(2)+RR(3))/3.
      YC=(ZZ(1)+ZZ(2)+ZZ(3))/3.
      RR(5)=RR(3)
      ZZ(5)=ZZ(3)
160  DO 200 NN=1,NT
      LM(1)=2*I-1
      LM(2)=2*J-1
      U(1)=ZZ(J)-ZZ(K)
      U(2)=ZZ(K)-ZZ(I)
      U(3)=ZZ(I)-ZZ(J)
      V(1)=RR(K)-RR(J)
      V(2)=RR(I)-RR(K)
      V(3)=RR(J)-RR(I)
      AREA=(RR(J)*U(2)+RR(I)*U(1)+RR(5)*U(3))/2.
      VOL=VOL+AREA
      COMM=.25/AREA
      XNT=NT
      COM=2./XNT
      COM=COM*COMM
      DO 180 L=1,3
      II=LM(L)
      ST(1,II)=ST(1,II)+U(L)*COM
      ST(2,II+1)=ST(2,II+1)+V(L)*COM
      ST(3,II)=ST(3,II)+V(L)*COM

```

```

ST(3,II+1)=ST(3,II+1)+U(L)*COM
DO 180 M=1,3
JJ=L*(M)
S(II,JJ)=S(II,JJ)+(U(L)*C(1,1)*U(M)+V(L)*C(3,3)*V(M))*COMM
S(II,JJ+1)=S(II,JJ+1)+(U(L)*C(1,2)*V(M)+V(L)*C(3,3)*U(M))*COMM
S(II+1,JJ+1)=S(II+1,JJ+1)+(V(L)*C(1,1)*V(M)+U(L)*C(3,3)*U(M))*COMM
S(JJ+1,II)=S(II,JJ+1)
180 CONTINUE
I=J
J=J+1
200 CONTINUE
IF(IX(N,3)-IX(N,4)) 220,250,220
220 DO 240 I=1,2
KK=10-I
DO 240 K=1,KK
CC=S(KK+1,K)/S(KK+1,KK+1)
DO 230 J=1,3
230 ST(J,K)=ST(J,K)-CC*ST(J,KK+1)
DO 240 J=1,KK
240 S(J,K)=S(J,K)-CC*S(J,KK+1)
250 CONTINUE
RETURN
END

```

SUBROUTINE STRESS

```

DOUBLE PRECISION S,C,B,A,P,ST,SIG,U,V,CC,BB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FOR,TOL,UR,UZ ,SIGI,VOL,COM,E,COSA,SINA
2, CU, Z, EE, XC, YC, PK, ACELR, ACELZ, TEMP, T, R, Q, RO, Z
3, CC, BB, CK, SS, SC, S2, C2, EPS, DT
COMMON ACELR, ACELZ, TEMP, Q, RO(12), R(900), Z(900), T(900), PR(200,2),
1NUMMP, NUMEL, NUMMAT, NUMPC, MTYPE, LLL, N, HED(18), NTC(10), CODE(900),
2MTC(10), IRC(200), JBC(200), NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIGI(900,6),CU(1800),TOL,VOL
1,E(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,EE(3),LM(4),IX(800,5)
COMMON /BANARG/ MBAND,NUMBLK,B(1800),A(108,54),KKK,JA
DEFINE FILE 1(50,1500,U,NBK),2(800,256,U,ID)
FOR = 0.0
MPRINT=0
DO 600 M=1,NUMEL
ID=M
FIND(2>ID)
N=M
IX(N,5)=IARS(IX(N,5))
MTYPE=IX(N,5)
SIGI(N,4)=0.
SIGI(N,5)=0.
SIGI(N,6)=0.
IF(IX(N,3)-IX(N,2)) 90,60,90
60 I=IX(N,1)
J=IX(N,2)
XC=(R(I)+R(J))/2.0
YC=(Z(I)+Z(J))/2.0
DX=R(J)-R(I)
DY=Z(J)-Z(I)
XL=DSQRT(DX**2+DY**2)
DU=B(2*J-1)-B(2*I-1)
DV=B(2*J)-B(2*I)
DI=DV*DY/XL+DU*DX/XL
SIG(1)=E(1,4,MTYPE)*DL*E(1,2,MTYPE)/XL+SIGI(N,1)*E(1,4,MTYPE)
IF(SIG(1).GT.0.) GO TO 100
SIGI(N,1)=SIG(1)
GO TO 500
100 SIGI(N,4)=E(1,2,MTYPE)*DL/XL+SIGI(N,1)
SIGI(N,1)=0.
GO TO 420
90 READ(2>ID) ((C(IJK,JJK),JJK=1,3),EE(IJK),IJK=1,3)
1,((S(JJI,KKI),KKI=1,8),JJI=1,8),((ST(IKK,JKK),JKK=1,8),IKK=1,3)
2,((RR(JII),ZZ(JII),JII=1,4),XC,YC,TEMP,VOL
MM=4

```

```

IF(IX(N,3)-IX(N,4)) 170,160,170
160 MM=3
170 DO 180 I=1,3
RR(I)=C.
DO 180 J=1,MM
II=2*J
JJ=2*IX(N,J)
180 RR(I)=RR(I)+ST(I,II)*B(JJ)+ST(I,II-1)*B(JJ-1)
IF(LLI.GT.1) GO TO 182
DT=TFMP-Q
DX=EE(3)*DT
DY=DX
SIG(1)=-C(1,1)*DX-C(1,2)*DY +SIGI(N,1)
SIG(2)=-C(2,1)*DX-C(2,2)*DY +SIGI(N,2)
SIG(3)=SIGI(N,3)
GO TO 184
182 DO 183 I=1,3
183 SIG(I)=0.C
184 CONTINUE
DO 190 I=1,3
DO 185 J=1,3
185 SIG(I)=SIG(I)+C(I,J)*RR(J)
190 CONTINUE
IF(LLI.EQ.1) GO TO 195
DO 192 I=1,3
192 SIG(I)= SIG(I)+SIGI(N,I)
195 CONTINUE
CC=(SIG(1)+SIG(2))/2.0
BB=(SIG(1)-SIG(2))/2.
CR=DSORT(BB**2+SIG(3)**2)
SIG(4)=CC+CR
SIG(5)=CC-CR
SIG(6)=C.0
IF((BB.FQ.C.).AND.(SIG(3).EQ.0.)) GO TO 200
SIG(6)=28.648*DATAN2(SIG(3),BB)
DX=C.0
200 SIGI(N,1)=SIG(1)
SIGI(N,2)=SIG(2)
SIGI(N,3)=SIG(3)
IF((SIG(4).LE.0.C0).OR.(MTC(MTYPE).EQ.0)) GO TO 500
IF(SIG(5).GE.0.C0001) GO TO 370
EPS=SIG(6)/57.296
CC=DCOS(EPS)
SS=DSIN(EPS)
C2=CC*CC
S2=SS*SS
SC=SS*CC

```

C

```

DX= EF(2)*SIG(4)
SIGI(N,4)= SIG(4)*C2+DX*S2
SIGI(N,5)= SIG(4)*S2+DX*C2
SIGI(N,6)=SIG(4)*SC-DX*SC

```

C
C

```

GO TO 400
370 SIGI(N,4)=SIG(1)
   SIGI(N,5)=SIG(2)
   SIGI(N,6)=SIG(3)
400 SIGI(N,1)=SIG(1)-SIGI(N,4)
   SIGI(N,2)=SIG(2)-SIGI(N,5)
   SIGI(N,3)=SIG(3)-SIGI(N,6)
420 DX=SIGI(N,4)**2+SIGI(N,5)**2+SIGI(N,6)**2
   DX=DSQRT(DX)
   IF(DX.LE.FOR) GO TO 450
   IJK=N
   FOR=DX
450 CONTINUE
500 IF(MPRINT) 550,520,550
520 WRITE(6,2000)
   MPRINT=50
550 MPRINT=MPRINT-1
   WRITE(6,2001)N,XC,YC,(SIG(I),I=1,6),DX
600 CONTINUE
   WRITE(6,2002)FOR,IJK
   IF(FOR.LE.TOL) NCHECK = 0
   RETURN
2000 FORMAT (7H1EL.NO. 7X 1HX 7X 1HY 4X 8HX-STRESS 4X 8HY-STRESS 3X
1 9HXY-STRESS 2X 10HMAX-STRESS 2X 10HMIN-STRESS 7H ANGLE 2X 17HUNB
2ALANCED FORCE )
2001 FORMAT (17,2F8.2,1P5F12.4,0P1F7.2,1PE20.4)
2002 FORMAT(30HMAXIMUM UNBALANCED FORCE = E12.5,16H IN ELEMENT NO,
1 15)
END

```

```
SUBROUTINE MODIFY(A,NEQ,MBAND,N)
DOUBLE PRECISION A
DIMENSION A(108,54)
DO 250 M=2,MBAND
K=N-M+1
IF(K.LE.0) GO TO 235
A(K,M)=0.0
235 K=N+M-1
IF(NEQ.LT.K) GO TO 250
A(N,M)=0.0
250 CONTINUE
A(N,1)=1.0
RETURN
END
```

```

SUBROUTINE LOAD
DOUBLE PRECISION S,C,B,A,P,ST,SIG,U,V,CC,BB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XI,RR,FOR,TOL,UR,UZ ,SIGI,VOL,COM,E,COSA,SINA
2, CU,ZZ,FE,XC,YC,PR,ACELR,ACELZ,TEMP,T, R,Q,RO,Z
3, PP1,PP2,DR,DZ,EP,XMM,DT
COMMON ACFLP,ACELZ,TEMP,Q,RO(12),R(900),Z(900),T(900),PR(200,2),
1NUMNP,NUMEL,NUMMAT,NUMPC,MTYPE,LLL,N,HED(18),NTC(10),CODE(900),
2MTC(10),IBC(200),JBC(200),NCHECK
COMMON/SYSARG/UR(900),UZ(900),SIGI(900,6),CU(1800),TOL,VOL
1,E(8,4,12)
COMMON /ARG/C(3,3),S(10,10),SIG(6),P(8),ST(3,10),RR(5),ZZ(5),
1XC,YC,EF(3),LM(4),IX(800,5)
COMMON /BANARG/ MBRAND,NUMBLK,B(1800),A(108,54),KKK,JA
DO 50 N=1,NUMNP
B(2*N-1)=UR(N)
B(2*N)=UZ(N)
UR(N)=0.
UZ(N)=0.
50 CONTINUE
IF((NUMPC.EQ.0).OR.(LLL.GT.1)) GO TO 300
DO 200 L=1,NUMPC
I=IHC(L)
J=JHC(L)
DR=Z(I)-Z(J)
DZ=R(J)-R(I)
PP2=(PR(L,2)+PR(L,1))/6.
PP1=PP2+PR(L,1)/6.
PP2=PP2+PR(L,2)/6.
II=2*I
JJ=2*J
B(II-1)=B(II-1)+PP1*DR
B(II)=B(II)+PP1*DZ
B(JJ-1)=B(JJ-1)+PP2*DR
B(JJ)=B(JJ)+PP2*DZ
200 CONTINUE
300 DO 700 N=1,NUMEL
I=IX(N,1)
J=IX(N,2)
K=IX(N,3)
L=IX(N,4)
MTYPE=IX(N,5)
IF(LLL.EQ.1) GO TO 330
IF(MTC(MTYPE).EQ.0) GO TO 700
IF(SIGI(N,4).NE.0.) GO TO 320
IF(SIGI(N,5).NE.0.) GO TO 320
IF(SIGI(N,6).NE.0.) GO TO 320
GO TO 700
320 CONTINUE

```

```

330 IF(J.EQ.K) GO TO 500
    ID=N
    RFAD(2*ID) ((C(IIK,JJK),JJK=1,3),EE(IIK),IIK=1,3)
    1 ,(S(JJI,KKI),KKI=1,8),JJI=1,8),(ST(IKK,JKK),JKK=1,8),IKK=1,3)
    2 ,(RK(JII),ZZ(JII),JII=1,4),XC,YC,TEMP,VOL
    IF(LLL.EQ.1) GO TO 400
    SIG(1)=-SIGI(N,4)
    SIG(2)=-SIGI(N,5)
    SIG(3)=-SIGI(N,6)
    GO TO 450
400 DT=TEMP-Q
    DX=FF(3)*DT
    DY=EE(3)*DT
    SIG(1)=-C(1,1)*DX-C(1,2)*DY +SIGI(N,1)
    SIG(2)=-C(2,1)*DX-C(2,2)*DY +SIGI(N,2)
    SIG(3)=SIGI(N,3)
450 DO 520 I=1,8
    P(I)=0.0
    DO 510 J=1,3
510 P(I)=P(I)-ST(J,I)*SIG(J)
520 P(I)=P(I)*VOL
    IF(LLL.EQ.1) GO TO 540
    DO 530 I=1,3
530 SIG(I)=0.0
    GO TO 600
540 MM=4
    IF(IX(N,3).EQ.IX(N,4)) MM=3
    XMM=MM
    DY=VOL*ACELZ*RO(MTYPE)/XMM
    DX=VOL*ACELR*RO(MTYPE)/XMM
    DO 550 I=1,MM
    P(2*I)=P(2*I)+DY
550 P(2*I-1)=P(2*I-1)+DX
    GO TO 600
500 CALL ONED
    DX=R(J)-R(I)
    DY=Z(J)-Z(I)
    EP=-SIGI(N,4)/E(1,2,MTYPE)
    DX=DX*EP
    DY=DY*EP
    P(1)=S(1,1)*DX+S(1,2)*DY
    P(2)=S(2,1)*DX+S(2,2)*DY
    P(3)=-P(1)
    P(4)=-P(2)
600 DO 620 II=1,4
620 LM(II)=2*IX(N,II)-1
    DO 650 JJ=1,4
    II=LM(JJ)

```

```
B(II)=B(II)+P(2*JJ-1)
650 B(II+1)=B(II+1)+P(2*JJ)
700 CONTINUE
DO 750 N=1,NUMNP
IF(CODE(N).EQ.0.) GO TO 750
IF((CODE(N).EQ.1.).OR.(CODE(N).EQ.3.)) B(2*N-1)=0.
IF((CODE(N).EQ.2.).OR.(CODE(N).EQ.3.)) B(2*N)=0.0
750 CONTINUE
RETURN
END
```

```

SUBROUTINE SYMBAN
DOUBLE PRECISION S,C,R,A,P,ST,SIG,U,V,CC,BB,CR,AREA,COMM,DU,DV,DL,
1DX,DY,XL,RR,FUR,TOL,IJR,UZ ,SIGI,VOL,COM,E,COSA,SINA
2,CU,ZZ,FE,XC,YC,PR,ACELR,ACELZ,TEMP,T,DT,Q,RO
COMMON /BANARG/ MBAND,NUMBLK,B(1800),A(108,54),KKK,JA
OFFINE FILE 1(50,1500,U,NRK),2(800,256,U,10)

```

C
C

```

NN=54
NL=NN+1
NH=NN+NN
NR=C
NNK=0
NRK=1
FIND(1*1)
IF(KKK.GT.1) GO TO 2000
GO TO 150
100 NRK=NRK+JA
FIND(1*NRK)
NB=NB+1
DO 125 N=1,NN
NM=NN+N
DO 125 M=1,MBAND
A(N,M)=A(NM,M)
125 A(NM,M)=0.
C
IF(NUMBLK-NB) 150,200,150
150 READ(1*NRK)((A(N,M),M=1,MBAND),N=NL,NH)
NNK=NRK
IF(NR) 200,100,200
200 DO 300 N=1,NN
IF(A(N,1)) 225,300,225
225 DO 275 L=2,MBAND
IF(A(N,L)) 230,275,230
230 C=A(N,L)/A(N,1)
I=N+L-1
J=0
DO 250 K=L,MBAND
J=J+1
250 A(I,J)=A(I,J)-C*A(N,K)
A(N,L)=C
275 CONTINUE
300 CONTINUE
NRK=NRK-JA
WRITE(1*NRK)((A(N,M),M=1,MBAND),N=1,NN)
IF(NUMBLK.EQ.NB) RETURN
GO TO 100
2000 NQ=0

```

```

NEQ=1800
GO TO 450
400 NB=NB+1
NRK=NNK+JA
FIND(1,NBK)
DO 425 N=1,NN
NM=NN+N
DO 425 M=1,MBAND
A(N,M)=A(NM,M)
425 A(NM,M)=0.
IF(NUMBLK-NB) 450,500,450
450 REAL(1,NBK)((A(N,M),M=1,MBAND),N=NL,NH)
NNK=NBK
IF(NB) 500,400,500
500 DO 550 N=1,NN
J=NQ+N
DO 540 L=2,MBAND
I=J+L-1
IF(NEQ-I) 545,540,540
540 R(I)=R(I)-A(N,L)*B(J)
545 IF(A(N,1).EQ.0.) A(N,1)=1.
550 R(J)=R(J)/A(N,1)
NRK=NNK-JA
IF(NUMBLK.EQ.NB) GO TO 700
600 NQ=NQ+NN
GO TO 400
700 DO 750 M=1,NN
N=NN+1-M
J=NQ+N
DO 750 L=2,MBAND
IF(A(N,L)) 740,750,740
740 I=J+L-1
IF(NEQ-I) 750,745,745
745 R(J)=R(J)-A(N,L)*R(I)
750 CONTINUE
NB=NB-1
IF(NB.EQ.C) RETURN
FIND(1,NBK)
DO 800 N=1,NN
NM=NN+N
DO 800 M=1,MBAND
A(NM,M)=A(N,M)
800 A(N,M)=0.
READ(1,NBK)((A(N,M),M=1,MBAND),N=1,NN)
NRK=NRK-JA
NQ=NQ-NN
GO TO 700
END

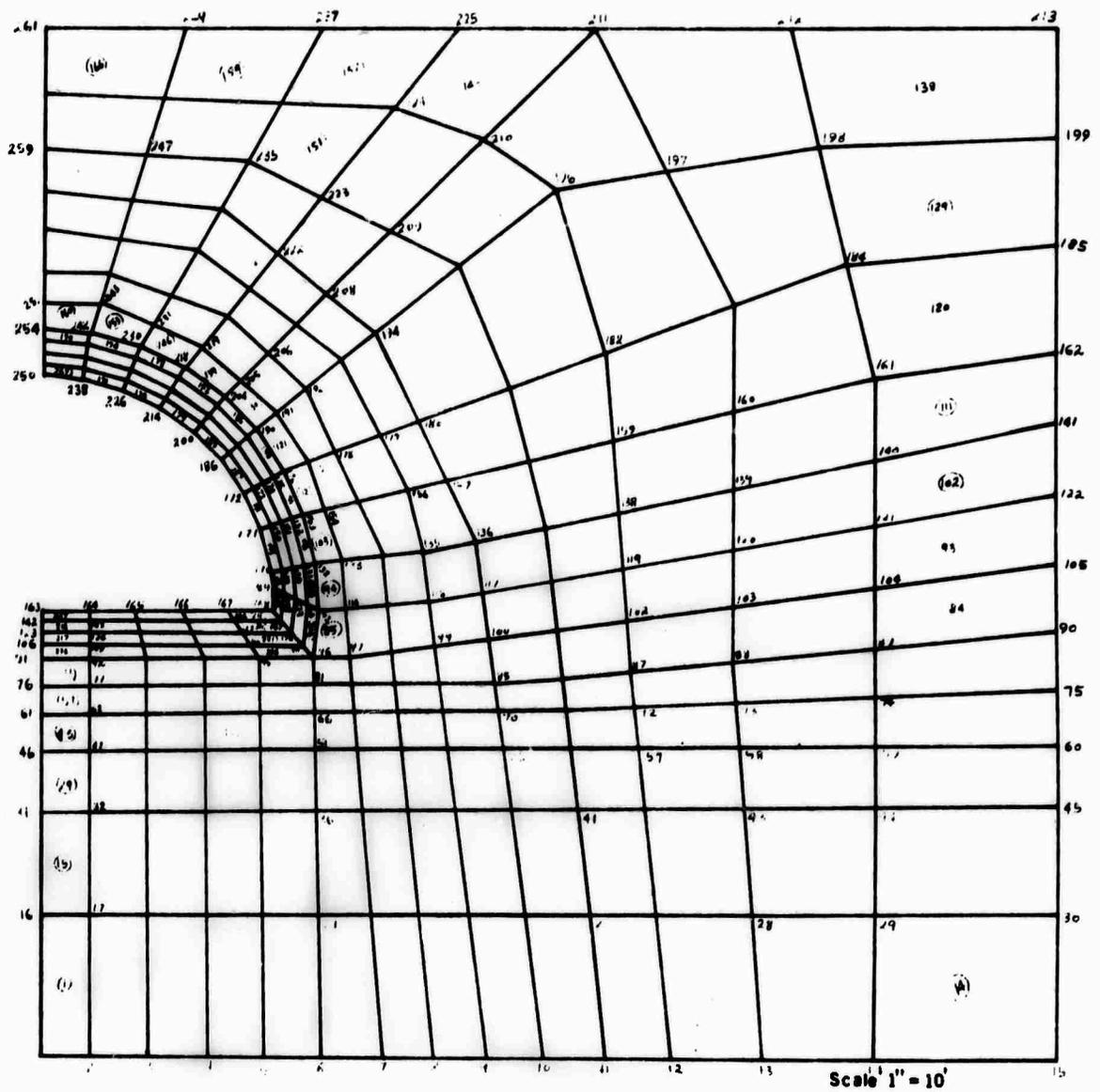
```

4.4 Example

A tunnel whose configuration and finite element idealization is shown in Fig. 4-1 was solved assuming rock to crack under tension. Lining was assumed to be capable of resisting tensile loads. Figs. 4-2 and 4-3 show the initial elastic and final 'no tension' solutions. The tensile stresses are indicated by arrows. Comparison of Fig. 4-2 and 4-3 shows the redistribution of stresses caused by the inability of rock to withstand tension. In the example, the excavation and lining of tunnel is assumed to be a single step ignoring the effect of sequential operations. This is unrealistic and further development will allow for actual sequence of construction. However, the example illustrates use of the computer program and is similar to the one used by Zienkiewicz, Valliappan and King (1968).

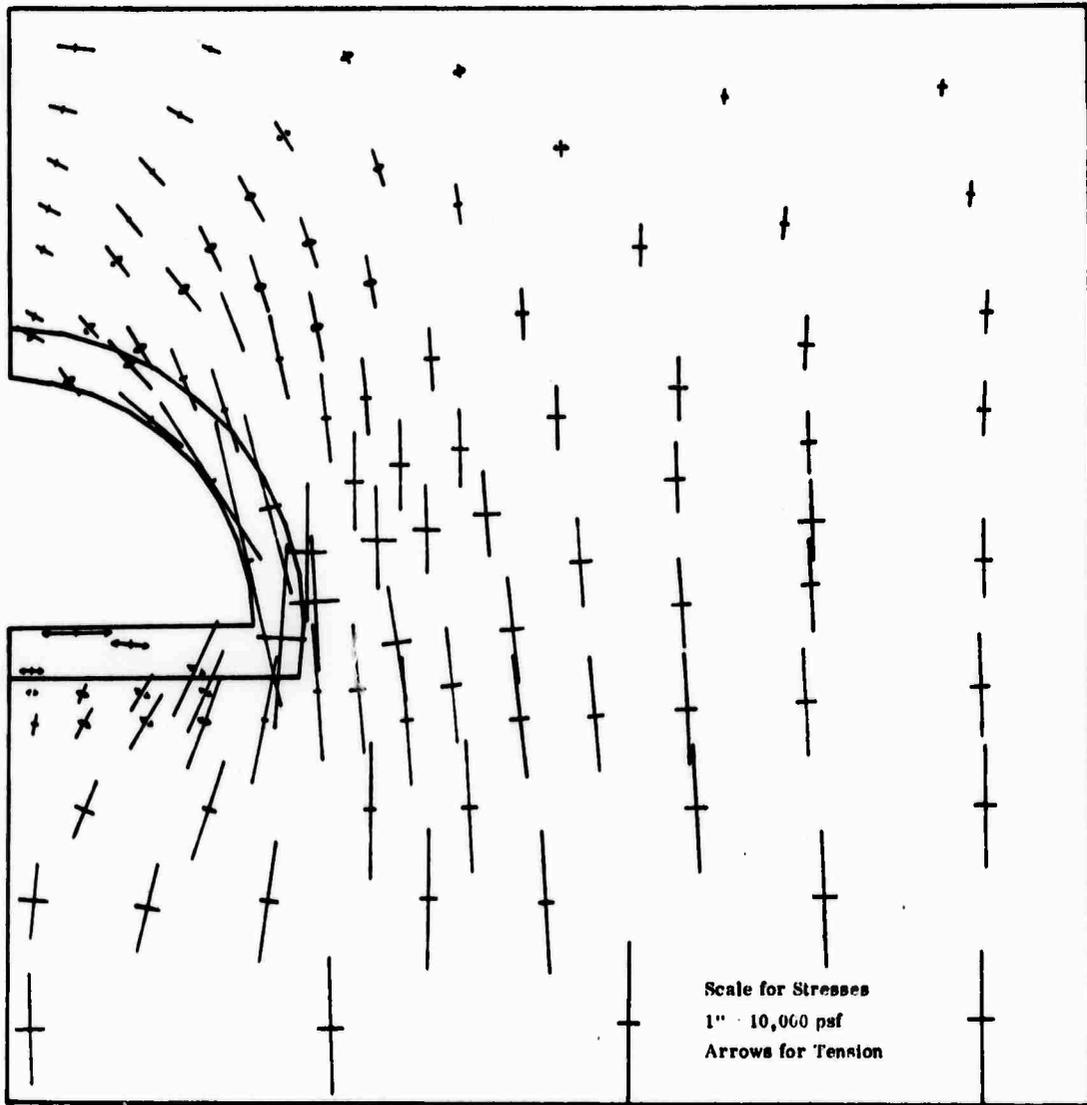
The following material constants were used for the solution:

Lining:	$E = 2 \times 10^6$ psi
	$\nu = 0.15$
	$\gamma = 150$ lbs./ft. ³
Rock:	$E = 1 \times 10^6$ psi
	$\nu = 0.2$
	$\gamma = 150$ lbs./ft. ³



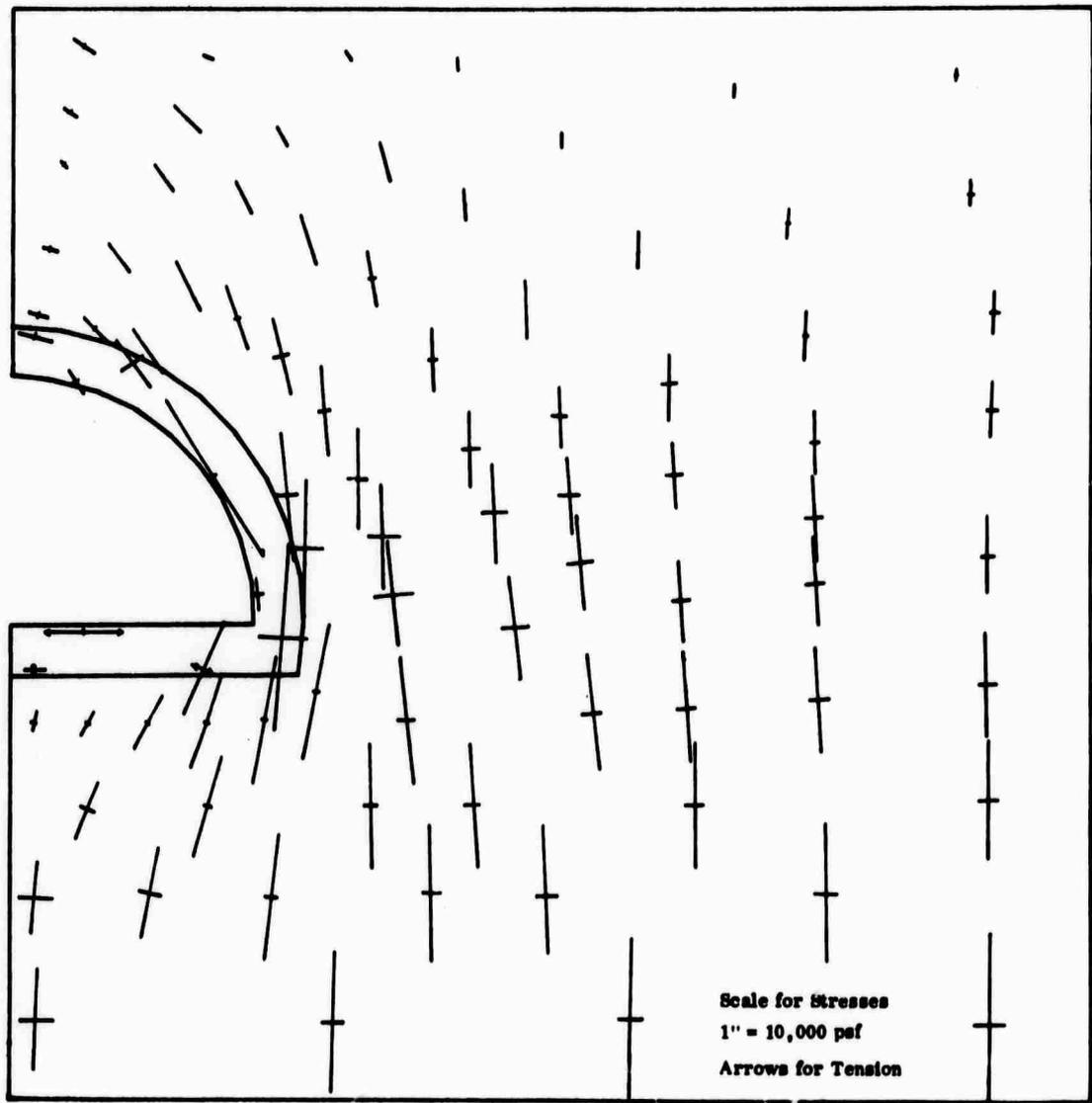
FINITE ELEMENT IDEALIZATION OF LINED TUNNEL

FIG. 4.1



ELASTIC SOLUTION FOR THE LINED TUNNEL

FIG. 4.2



FINAL NO-TENSION SOLUTION FOR THE LINED TUNNEL

FIG. 4.3

CHAPTER V
ADDITIONAL COMMENTS

Chapter V. Additional Comments

The work reported is continuing and significant changes may be made before the finite element computer programs reach their final form. In the case of elastic-plastic materials following Mohr-Coulomb Theory non-monotonic loads have to be allowed for and also alternative numerical solution procedures have to be examined. For the case of jointed rock the program included in the report represents a necessary first step. Modification to include the Griffith failure criterion is being done. For Griffith rupture caused by a tensile stress field acting on pre-existing flaws, the solution scheme appears to be fairly straight forward. However, the case of fracture under compressive stress fields may require development of new methods.

List of References

1. Baker, L. E., Sandhu, R. S. and Shieh, W. Y. (1969), "Application of Elasto-Plastic Analysis in Rock Mechanics by Finite Element Method," Proc. Eleventh Symposium on Rock Mechanics, Berkeley, California.
2. Bieniawski, Z. T. (1967), "Mechanism of Brittle Fracture of Rock," Parts I, II, III, Int. J. Rock Mech. Min. Sci., V. 4, p. 395-525.
3. Brace, W. F., (1964), "Brittle Fracture of Rocks," p. III, in State of Stress in the Earth's Crust, ed. W. R. Judd, Am. Elsevier Publishing Co, Inc. New York.
4. Brace, W. F. (1966), "Dilatancy in the Fracture of Crystalline Rocks," J. Geoph. Res., V. 71, No. 16, 3939-53.
5. Brady, B. T. (1970), "A Mechanical Equation of State for Brittle Rock, Part I, The Pre-Failure Behavior of Brittle Rock," Int. J. Rock Mech. Min. Sci., V. 7, 385-421.
6. Brady, B. T. (1969a), "The Nonlinear Mechanical Behavior of Brittle Rock, Part I, Stress-Strain Behavior During Regions I and II," Int. J. Rock Mech. Min. Sci., V. 6, No. 2, 211-225.
7. Brady, B. T. (1969b), "The Nonlinear Mechanical Behavior of Brittle Rock, Part II, Stress-Strain Behavior During Regions III and IV, Int. J. Rock Mech. Min. Sci., V. 6, No. 3, 301-310.
8. Clough, R. W., (1960), "The Finite Element Method in Plane Stress Analysis," Proc. 2nd ASCE Conf. on Electronic Computation, Pittsburgh, Pa.
9. Clough, R. W. (1965), "The Finite Element Method in Structural Mechanics," Chapter 7, STRESS ANALYSIS, ed. O.C. Zienkiewicz and G.S. Holister, London: Wiley.
10. De Arantes e Oliveira, Eduardo R. (1968), "Theoretical Foundations of the Finite Element Method," Int. J. Solids Structures, V. 4, 929-952.
11. Drucker, D. C. (1951), "A More Fundamental Approach to Plastic Stress-Strain Relation," Proc. First U.S. National Congress of Applied Mechanics, ASME, 487-491.
12. Drucker, D.C. and Prager, W. (1952), "Soil Mechanics and Plastic Analysis or Limit Design," Q. App. Math, v. 10, 157-165.

13. Drucker, D. C., Gibson, R. E. and Henkel, D. J. (1955), "Soil Mechanics and Work-Hardening Theories of Plasticity," Paper No. 2864, Trans., ASCE.
14. Duncan, J.M. and Goodman, R.E. (1968), "Finite Element Analysis of Slopes in Jointed Rock," Report No. TE-68-1, U.S. Army Engineer Waterways Exp. Station, Corps of Engineers.
15. Einstein, H.H. , Bruhn, R.W. and Hirschfeld, R.C. (1970), "Mechanics of Jointed Rock, Experimental and Theoretical Studies," Soil Mech. Publication No. 268 M. I. T.
16. Felippa, C. (1966), "Refined Finite Element Analysis of Linear and Nonlinear Two-Dimensional Structures," Ph.D. Thesis, University of California, Berkeley.
17. Felippa, C.A. and Clough, R.W., "The Finite Element Method in Solid Mechanics," V.2, SIAM-AMS Proc., Numerical Solution of Field Problems in Continuum Physics, Am. Math. Soc., Providence, R.I.
18. Goodman, R.E., Taylor, R.L. and Brekke, T.L. (1968), "A Model for the Mechanics of Jointed Rock," ASCE, SM3, 637-658.
19. Green, A.E. and Naghdi, P.M. (1965), "A General Theory of an Elastic-Plastic Continuum," Arch. Rational Mech. Anal, V. 18, 251-281.
20. Hill, R. (1950), The Mathematical Theory of Plasticity, Oxford University Press.
21. Holand and Bell (1970), "Finite Element Methods in Stress Analysis," TAPIR, The Technical University of Norway, Trondheim, Norway, 2nd print.
22. Koiter, W.T. (1953), "Stress-Strain Relations, Uniqueness and Variational Theorems for Elastic-Plastic Materials with a Singular Yield Surface," Quar. Appl. Math. Vol. 11, 350-354.
23. Malina, H. (1970), "The Numerical Determination of Stresses and Deformations in Rock Taking into Account Discontinuities," Rock Mechanics 2, 1-16.
24. Melosh, R. J. (1963), "Basis for Derivation of Matrices for the Direct Stiffness Method," AIAA, V. 1, No. 7.

25. Naghdi, P.M. (1960), "Stress-Strain Relations in Plasticity and Thermo-plasticity," *Plasticity: Proceedings of the 2nd Symposium on Naval Structural Mechanics*, Pergman Press, 121-169.
26. Oden, J.T. (1969), "A General Theory of Finite Elements, I. Topological Considerations; II. Applications, *Int. J. Num. Math. Eng.*, V. 1, 205-221, 247-259.
27. Pian, T.H.H., (1968), "Variational Principles and Their Application to Finite Element Methods," *Lecture Notes on "Finite Element Methods in Solid Mechanics,"* 24-38, M.I.T.
28. Prager, W. (1949), "Recent Developments in the Mathematical Theory of Plasticity," *J. App. Phys.*, V. 20, p. 235.
29. Proc. Conf. Matrix Methods in Structural Mechanics, Wright-Patterson Air Force Base, Ohio, 1965.
30. Proc. 2nd Conf. on Matrix Methods in Structural Mechanics, AFFDL-TR-68-150, Wright Patterson AFB, Ohio, 1968.
31. Reyes, S.F. (1965), "Elasto-Plastic Analysis of Underground Openings by the Finite Element Method," Ph.D. Thesis, University of Illinois.
32. Reyes, S.F. and Deere, D.V. (1966), "Elastic-Plastic Analysis of Underground Openings by Finite Element Method," *Proc., First Congress International Society of Rock-Mechanics, Lisbon.*
33. Sandhu, R.S. and Wilson, E.L. (1970), "Finite Element Analysis of Stresses in Mass Concrete Structures," Preprint, Symposium on Impact of Computers on the Practice of Structural Engineering in Concrete, ACI, St. Louis, Missouri.
34. Swanson, S.R. (1970), "Development of Constitutive Equations for Rocks," Ph.D. Thesis, University of Utah.
35. Walsh, J. B. (1965a), "The Effect of Cracks on the Compressibility of Rock," *J. Geophys. Res.*, V. 70, No. 2, 381-389.
36. Walsh, J.B. (1965b), "The Effect of Cracks on the Uniaxial Elastic Compression of Rocks," *J. Geophys. Res.*, V. 70, No. 2, 399-411.

37. Walsh, J. B. (1965c), "The Effect of Cracks in Rocks on Poisson's Ratio," J. Geophysical Res., V. 70, N. 20, 5249-57.
38. Wilson, E.L. (1963), "Finite Element Method in Two-Dimensional Problems," D. Engg. Thesis, University of California, Berkeley.
39. Wilson, E.L. (1965), "Structural Analysis of Axisymmetric Solids," J. AIAA, V. 3.
40. Yamada, Y., Yoshimura, N., and Sakurai, T. (1968), "Plastic Stress-Strain Matrix and its Application for Solution of Elastic-Plastic Problems by the Finite Element Method," Int. J. Mech. Sci., Pergamon Press, V. 10, 343-354.
41. Zienkiewicz, O.C. and Cheung, Y.K., (1967), The Finite Element Method in Structural and Continuum Mechanics, McGraw-Hill.
42. Zienkiewicz, O.C., Valliappan and King, I.P. (1969), "Elasto-Plastic Solutions of Engineering Problems, Initial Stress, Finite Element Approach," International Journal of Numerical Methods in Engg., V. 1, 75-100.
43. Zienkiewicz, O.C., Valliappan, S., and King, I.P. (1968), "Stress Analysis of Rock as a 'No Tension' Material," Geotechnique 18, No. 1, 56-66.
44. Zlamal, M. (1968), "On the Finite Element Method," Numer. Math. 12, 394-409.