CRYOGENIC CERAMIC 277 WATT Yb: YAG THIN-DISK LASER

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Subject terms: laser materials; ytterbium lasers; solid-state lasers; diode-pumped lasers; optical materials.

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1 Introduction

Cryogenic solid-state laser materials offer many improvements in thermo/optical, structural, and lasing properties; all of which are noted in the following short literature review. Our first and foremost reference to the literature, and perhaps one of the most relevant, pertains to the lasing properties of a relatively thick disk of 20-at.%-doped Yb:YAG crystal. The maximum fiber coupled pump power was 12 W at a wavelength of 970 nm with a spot size of 0.4 mm. Kasamatsu, et al. studied various lasing characteristics over a temperature range of 80 to 300 K and noted a decrease in threshold and an increase in slope efficiency as the temperature decreases. Further, they demonstrated a maximum output lasing power at a temperature near 150 K. The next study2 was a cryogenic Yb:YAG oscillator with a 15-mm, 5-at.%-doped Yb:YAG crystal sandwiched between two undoped YAG caps maintained near 100 K. A 165 W output for an input of 215 W gave a slope efficiency of 85%, optical-to-optical efficiency of 60%, and a threshold near 22 W. Later the same institution3 returned to a similar configuration and measured a threshold of 50 W, an optical-to-optical efficiency of 60%, and a pump power of 500 W that gave an output of around 300 W. In the intervening years4 a cryogenic sapphire-Yb:YAG sandwich yielded a 75 W output at a pump power of 106 W with an optical-to-optical efficiency of 70% and a slope efficiency of 80%. In a study initiated in 2007 and continued into subsequent years,5,6 multiple cryogenic sapphire-sandwiched Yb:YAG disks eventually gave an optical-to-optical efficiency of 42.2% with an output of 963 W for an input of 1871 W.7 The introduction of a total-reflection active mirror (TRAM) Yb:YAG disk sandwiched between a trapezoidal YAG prism and a liquid nitrogen (LN2) bath gave 275 W, an optical-to-optical efficiency of 65%, and a slope efficiency of 72% at an absorbed power of 450 W.8–10 Finally, we note a paper on cryogenic Yb:YAG 10-at.%-doped crystal at high peak and average power.11 They obtained 0.5 J at a 1-kHz rate and in doing so measured the small signal gain and showed that it rolls over at a pump power of 200 to 300 W for a spot size of 6 mm.

In the following sections we report the improvement in extraction, threshold, and distortion of a cryogenic ceramic Yb:YAG thin-disk laser operating at initial temperatures near 80 K, as compared to the same system operating near 280 K.
These correspond to LN$_2$ and R134a coolants. We show that for the same resonator the threshold drops dramatically by a factor of 50 while the slope efficiency changes from 54% (280 K) to 63% (80 K). These changes are all, for the most part, manifestations of the differences between quasi three-level and four-level Yb:YAG.

In contrast to the study by Kasamatsu, et al.,\textsuperscript{,1} we are interested in the scalability of cryogenic thin-disk lasers, and so we use a disk 3.5 times thinner with a pump beam size 17.5 times larger. This design addresses thermal and structural issues as well as disk and bonding failures attendant with small spot sizes. Additionally, high-power pumping at low temperatures becomes more difficult due to absorption spectral narrowing, and so we have shifted to a pump wavelength from 969 nm to 940 nm.

\section{Experimental}

Figure 1 shows a 2-D cross section of our setup. The evaporative spray cooling (ESC) nozzle unit (S) directs coolant into the interior of the cooling cap (C). Yb:YAG is attached to the cap using a void-free thin layer of indium solder. The pump is introduced into the cavity through the lens train (L). The pump is actually out of plane and does not propagate through the mirror (M2). It reflects off the one-piece parabolic mirror (M3) and corner mirrors (M1 and M2) are shown, and two additional corner mirrors are perpendicular to the plane of the paper, onto the disk surface to image the light for 16 passes of absorption\textsuperscript{,11}) The laser cavity is defined on one end by the ion-beam sputtered HR (\textgreater99.9\%) coatings applied to the Yb:YAG and the output coupler (OC R = 95 to 99\%) on the other. The resonator is approximately 35 cm long.

To manage the high heat loads, an advanced two-phase ESC technique developed by the RN\textsuperscript{i} Corporation (Oviedo, FL) is utilized. Two versions were constructed. R134a refrigerant is one coolant. This system can maintain the disk temperature from 268 to 290 K and is capable of removing heat fluxes of up to 200 W/cm$^2$ with a convective heat transfer coefficient of 125,000 W/m$^2$K. The convective heat transfer coefficient $h$, is defined by $h = q''/\Delta T_m$. $q''$ is the heat flux ($q'' = Q/A$ where $Q$ is the total heat load and $A$ is the heat transfer area) and $\Delta T_m$ is the difference between the spray surface temperature and the coolant saturation temperature. A second system uses LN$_2$, allowing laser operation near 77 K. Traditional LN$_2$ cryostats are limited to heat fluxes on the order of 20 – 25 W/cm$^2$. The two-phase ESC technique consistently removes heat fluxes of 100 W/cm$^2$ and has even removed heat fluxes as high as 160 W/cm$^2$ with a convective heat transfer coefficient of 80,000 W/m$^2$K. In order to put these heat transfer coefficient values in perspective, other commonly utilized cooling technologies can be considered for their respective performance. A typical natural convection with gases provides only 25 W/m$^2$K, while forced convection with liquids, such as jet impingement, could achieve 20,000 to 50,000 W/m$^2$K. Besides its high heat transfer performance, ESC allows very efficient coolant usage and results in lightweight and compact thermal management schemes. The ESC system in this work features an open-cycle design comprised of a nozzle located at the bottom of a stainless steel reservoir mounted on a base. The reservoir is designed to vent excess bubbles generated by an ambient heat leak, thus enabling a consistent spray. Heat generated by the thin-disk is removed by the liquid droplets impinging on the inner surface of the interface cap and vaporizing that cools the crystal through a two-phase heat transfer process. The generated vapor and any remaining liquid exit an exhaust fitting on the base section of the reservoir/nozzle assembly.

In the laboratory experiments, the difference between the two systems is the cavity environment. For cryogenic operation, the entire thin-disk laser is placed within a vacuum chamber so as to eliminate the ambient humidity. Air in the chamber is initially evacuated, and then the dry gas overflow of LN$_2$ is introduced to create a slight overpressure. This setup is sufficient to prevent condensation on optical surfaces and to cool the optics. The pump light is coupled through a port, necessitating a different set of optics to effectively collimate and launch the pump light onto the thin-disk resonator. This resulted in a modest degradation in the cumulative beam that is incident on the disk surface. In either case, the laser pump power is increased slowly so as to allow a thermopile detector time to settle. For room temperature operation (R134a), a thermal camera is imaged onto the disk surface to track the surface temperature increase. The camera was not calibrated at cryogenic temperatures and not used when the disk was LN$_2$-cooled.

The thin disk is void-free indium soldered to the mounting cap. A fluxless indium soldering process was developed by Precision Photonics (Boulder, CO) and Enerdyne Solutions (North Bend, WA). The mounting worked well for both cryogenic and room temperature operation. Both capped and uncapped ceramic Yb:YAG thin-disk gain material were tested. The uncapped gain media is 200 μm thick, 9.8% Yb-doped, 14 mm in diameter provided by Konoshima Corporation (Japan). The gain media is coated with a high reflective coating on one end. Capped thin-disks have the gain media bonded to undoped and antireflection coated for 940/1030 nm YAG. Precision Photonics bonds them using the chemically activated direct bonding (CADB) technique.
Figure 2 shows the lasing power versus the incident pump power at 15°C and 80 K. The inset spectral peak in Fig. 2 is the 80 K laser output and is centered near 1029.1 nm. At room temperatures this peak is near 1030.2 nm. The resonator uses a 98% reflective output coupler with a 2-m radius of curvature. The pump spot size is approximately 7 mm at room temperature with a nearly top hat profile. In the cryogenic case the profile is not an ideal flat top and has a spot size slightly greater than 7 mm. The room-temperature lowest lasing threshold is 155 W, at which point the disk surface temperature is 308 K. The slope efficiency is 54%, which remained linear throughout the pump range. At 510 W pump, the power reached 184 W.

In stark contrast to the R134a results, Fig. 2 shows that the cryogenic laser threshold plummets to near 10 W. This is a clear demonstration of the superiority of a four-level laser over the quasi three-level system. It also shows the improvement in emission and absorption cross sections at cold temperatures. For the cryogenic operation, the initial slope efficiency was only 43%. Above 200 W pump, the efficiency increased to 63% and remained linear beyond this point. At 520 W pump, the maximum power achieved was 277 W. We believe that this increase in slope efficiency can be traced to two factors. The first factor is the pump wavelength, which is just above 932 nm at low powers. The diode junction heats as the driving current increases, causing a wavelength red-shift of 0.0076 nm/W. At cryogenic temperatures the absorption cross-section increases substantially as the wavelength shifts from 932 to 935 nm. Thus the absorbed pump increases super-linearly with respect to the incident pump. Attempts to red-shift the light via the cooling water temperature were unsuccessful. In contrast, at room temperature the absorption profile increases slowly at these wavelengths. A second reason for the improvement is the lack of the ideal top hat shape of the pump spot. At powers very near threshold, the laser operates with few transverse modes oscillating. Yb at the center of the disk is inverted prior to the periphery. As the pump power increases, more modes reach the threshold condition and augment the total power. This effect is quite small but observable. Plotting the cryogenic laser power versus the absorbed pump would likely result in a linear slope with an efficiency superior to 63%. However, the resonator design prevented us from measuring the residual pump, and so the absorbed pump values would be an undesirable computed estimate.

Figure 3 shows the on-axis spontaneous emission when the thin disk is pumped with a low-power, 25-W, 915-nm pump source. No output coupler optic is in place. The material is ceramic Yb:YAG rather than crystal; the latter has been reported previously. The initial spectral sweep is taken at 15°C using a fiber probe positioned near the surface of the disk. The disk is subsequently cooled to 80 K and the measurement repeated. The integrated emission curves show that the emission at cryogenic temperatures is 1.8 times greater than the room temperature emission. Interestingly, there is a strong dip in the cryogenic emission spectra around the 968-nm zero-phonon line, which is due to the reduction in homogeneous broadening. The width of this dip is 8.2 nm on both sides of the line, and for ceramic thin disk the dip is not symmetrical. Note that the low temperature behavior around the zero phonon line is the major difference between ceramic and crystal emission spectra. Specifically, at low temperatures the crystal emission is symmetric and narrower than the asymmetric ceramic emission.

In the lasing situation, amplified spontaneous emission (ASE) emitted at a high angle (~80 deg) off of the optical axis of the disk was measured. This spectrum is not shown because it is very similar to the emission shown in Fig. 3. The emission peaks are amplified and narrowed. Instead of comparing ASE spectra between room and cryogenic cases, which is difficult due to different experimental arrangements and power levels, we examine the relative differences in the spectra in Fig. 3. For example, the collection angle of the fiber and its position is similar between the room and cryogenic experiments, but not identical due to a difference in the setup. Light is coupled into a 50-μm multimode fiber attached to an Agilent 86146B Optical Spectrum Analyzer. The pump laser power is 180 W, giving laser outputs of $P_{98K} = 72$ W, $P_{288K} = 7$ W. The predominant emission is scattered lasing light at 1030 nm (1029 nm at cryogenic). At 80 K there is a strong reduction in emission away from the laser line. The zero-phonon 969 nm line is 17.1 dB less than the 1029 nm peak while the peak at 1048.8 nm is 20.23 dB lower. In contrast, at room temperature the 969 nm and 1048.8 nm peaks are only 4.4 dB and 8.2 dB less than the laser peak intensity. As suggested in Fig. 2, there is a strong reduction in spurious ASE generated when Yb:YAG behaves as a four-level laser. Great effort has been spent on suppressing the ASE in the radial direction (undoped caps or beveled edges, for example). Broad band ASE is suppressed in the four-level case due to a sharpening of the emission lines. The large gain cross-section on the lasing line enhances stimulated emission into the lasing mode.

The 98% output coupler optic and 0.2-mm disk thickness are optimal for R134a or water cooling methods. It is not likely to be so for the cryogenic case. They were used here for direct comparison. Brown computed pump absorption optimization at room and cryogenic temperatures as a function of pump wavelength and the optical thickness (doping density x penetration distance), which can be applied here experimentally. Furthermore, Contag computed that the optical efficiency approaches ~85% at low temperatures regardless of the number of pump passes on the disk surface. Hence another parameter to consider is the number of pump passes. It is simpler and easier to experimentally align a few-pass resonator than the 8-pass one used here. These options make the cryogenic thin-disk laser an appealing system to simultaneously pursue very high power and efficient lasers.

We modeled our cryogenic disk/mount structure using the Finite Element Model (FEM) software COMSOL. Figures 4 and 5 show the temperatures and deformation of the disk/
mount for R134a (Fig. 4) and LN₂ (Fig. 5) cooling temperatures. Both figures show the outline of the undeformed disk (black lines): in Fig. 4 the topmost rectangle is the anti-ASE cap, while in Fig. 5 the topmost rectangle is the Yb gain medium. These figures are a cross-section of the rotationally symmetric disk/mount shown in Fig. 6. Further, the disk in Fig. 4 has a YAG anti-ASE cap while the disk in Fig. 5 is without a cap. The cap is a 1-mm YAG disk that is index-matched to the active 200 μm thick Yb: YAG. This reduces the trapped ASE rays in the doped region and hence decreases heating. These two disk configurations were determined by availability of disks and experimental constraints.

The thermal/structural constants of Yb: YAG at these two temperatures are: R134a, \( K = 5 \text{ W/m/K} \), \( E = 280 \times 10^9 \text{ Pa} \), \( \alpha = 7.6 \times 10^{-6} / \text{°K} \); LN₂, \( K = 40 \text{ W/m/K} \), \( E = 311 \times 10^9 \text{ Pa} \), \( \alpha = 2, 1 \times 10^{-6} / \text{°K} \). \( K \) is the thermal conductivity, \( E \) is Young’s constant, and \( \alpha \) is the thermal expansion coefficient.

For the Yb cap \( K = 10 \text{ W/m/K} \). A 33-μm layer of indium solder layer lies between the Yb: YAG and cap. Indium material parameters are determined by COMSOL default values.

For these experiments the maximum pump power is 520 W, which gives a pump intensity of 0.13 kW/cm² for a spot size of 7 mm. The photon efficiency is about 10% and for an absorption coefficient of 1000/m, the heat load is 1.3 kW/cm³. The dependence of this and ASE on pump power is discussed in Ref. 14. The absorbed heat load is modeled as a ninth order super-Gaussian profile. Returning to Figs. 4 and 5, note that in the figures the deformation is not to scale, but is generated as COMSOL’s default value. The coloring in these figures shows the temperature distribution of the cross section. The R134a disk shows a maximum temperature at the center of 309 K with a decrease of 16 K at the edge of the spot size and a drop of 28 K across the entire disk. For the LN₂ case, the center temperature is 93 K to 84 K at the end of the pump spot size and a 15 K drop across the entire disk to 78 K. Thus, for the same heat load the cryogenic temperature variation across the disk surface is about half of the R134a temperatures. The deformation in Figs. 4 and 5 shows that the side of the mount caves in as the disk bends convexly. A more detailed view of the deformation is shown in Figs. 7 and 8 for the top and bottom of the 200-μm Yb: YAG gain layer. These surfaces are important because the bottom is the HR side and presents a curvature to the cavity. The difference between the top and bottom contribute to the \( dn/dT \) and stress phases. The major difference between Figs. 7 and 8 is that the entire disk is pulled down by 24 μm due to cooling of the mount at LN₂ temperatures.

**Fig. 4** Thermal and structural plot for R134a. The topmost rectangle is the anti-ASE cap. Note the outline in black line of the undeformed disk/mount.

**Fig. 5** Thermal and structural plot for LN₂ cooling. The topmost rectangle is the Yb: YAG gain medium. Note the outline in black line of the undeformed disk/mount.

**Fig. 6** CuW mount for the Yb: YAG disk.
at R134a temperatures this does not occur. However, they show a sag of 0.2 μm for the R134a case and a 2-μm sag for the LN2 case, both across the pump beam spot size. Finally, we briefly note the stresses. For the uncollapsed disk at cryogenic temperatures, both the radial and azimuthal stress is compressive and uniform, except at the edge. However, for the capped disk at near-room temperatures, the azimuthal stress is compressive up to the beam edge then tensile to the edge caused by the radial tensile stress in the unpumped gain volume. This nonuniformity can lead to disk and bonding failure.

Our thin-disk laser, as well as most others, is inherently multimode. An estimate for the beam quality $M^2$ is calculated by assuming that $M$ is proportional to the radius of the pump spot $r_p$, divided by the lowest loss Gaussian beam radius on the thin disk surface $w_m$, thus $M^2 = (ar_p/w_m)^2$. For a thin-disk laser, Giesen$^{15}$ sets $a = 0.85$, which gives a $M^2$ value near 35. The beam quality can be further degraded by disk bowing or refractive index changes, $dn/dT$. Experiments were performed using a probe beam reflected from the surface of the disk and into a Shack-Hartmann wavefront analyzer. These revealed that in the lasing situation the disk can slightly bow with increasing laser power but that no thermal lensing occurs.

3 Summary

In summary, we demonstrated that cooling a 0.2-mm thick, 14-mm diameter, 9%-Yb-doped ceramic Yb:YAG thin-disk from 15°C (288 K) down to 80 K results in a laser threshold drop from 155 W to near 10 W. In concert, the cryogenic slope efficiency reached 63%, and the laser generated 277 W from a 520-W pump. A novel two-phase spray cooling method mitigates the heat produced within the Yb:YAG. Two systems, one using R134a and LN2 for the other, were built. Yb:YAG disks were indium-mounted to CuW caps that are interchangeable between the two systems. Hence the same disk was tested with cooling at 288 K (R134a) and also at 80 K (LN2). Material damage due to cycling the temperature from room to 80 K was not observed. This is ostensibly due to the soft and thick indium layer that buffers the tensile strain of the cap on the disk. These initial cryogenic results can be readily improved with wavelength stabilized pump diodes, minor refinements to the pump coupling, and optimization of the material characteristics.

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References


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